

15th AIMS Conference, Athens

Using Feynman diagrams in renormalisation, some new results

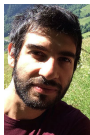
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Hendrik Weber (Münster), Tom Klose (London) and Nikolas Tapia (Berlin)



Motivation: Allen–Cahn SPDE

$$\partial_t \phi(t, x) = \Delta \phi(t, x) + f(\phi(t, x)) + \sqrt{2\varepsilon} \xi(t, x) \quad f(\phi) = \phi - \phi^3$$

▷ $x \in \mathbb{T}_L^d = (\mathbb{R}/L\mathbb{Z})^d$, L bifurcation parameter

▷ ξ : space-time white noise

▷ $\varepsilon = 0$: Gradient system with potential/energy function V

$$V[\phi] = \int_0^L \left[\frac{1}{2} \|\nabla \phi\|^2 - \frac{1}{2} \phi^2 + \frac{1}{4} \phi^4 \right] dx \quad \Rightarrow \quad \partial_t \langle \phi, \psi \rangle = -\nabla_\psi V(\phi)$$

▷ Stationary solutions for $\varepsilon = 0$, $d = 1$:

- $\phi_\pm(x) \equiv \pm 1$: stable
- $\phi_0(x) \equiv 0$: unstable
- $L > 2\pi$ ($L > \pi$ for Neumann b.c.): $\exists$ nonconstant stationary solutions

▷ Transition state: ϕ_{ts} lowest saddle point of V between ϕ_- and ϕ_+

- $L \leq 2\pi$ ($L \leq \pi$ for Neumann b.c.): $\phi_{ts} = \phi_0$
- $L > 2\pi$ ($L > \pi$): ϕ_{ts} is nonconstant in space

▷ **Question:** When $\varepsilon > 0$, what is the mean transition time between ϕ_- and ϕ_+ ?

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Eyring–Kramers law

$$dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$$

- ▷ $V : \mathbb{R}^n \rightarrow \mathbb{R}$ confining potential with minima x, y and saddle z
- ▷ $\tau_y^x = \inf\{t > 0 : x_t \in \mathcal{B}_\varepsilon(y)\}$ first-hitting time of $\mathcal{B}_\varepsilon(y)$, starting in x
- ▷ Eyring–Kramers law (1935, 1940):

Eigenvalues of Hessian of V at minimum x : $0 < \nu_1 \leq \nu_2 \leq \dots \leq \nu_n$

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$$\mathbb{E}[\tau_y^x] = 2\pi \sqrt{\frac{\lambda_2 \dots \lambda_n}{|\lambda_1| \nu_1 \dots \nu_n}} e^{[V(z) - V(x)]/\varepsilon} [1 + \mathcal{O}_\varepsilon(1)]$$

- ▷ Potential-theoretic proof:

Theorem: $A, B \subset \mathbb{R}^n$ disjoint. $\exists$ proba measure ν_{AB} on ∂A s.t.

$$\int_{\partial A} \mathbb{E}^x[\tau_B] \nu_{AB}(dx) = \frac{1}{\text{cap}(A, B)} \int_{B^c} e^{-V(y)/\varepsilon} h_{AB}(y) dy$$

◊ $h_{AB}(y) = \mathbb{P}^y\{\tau_A < \tau_B\}$ commitor, $\text{cap}(A, B)$ capacity

◊ Dirichlet principle: $\text{cap}(A, B) = \inf_{h|_A=1, h|_B=0} \varepsilon \int_{\mathbb{R}^d} e^{-V(x)/\varepsilon} |\nabla h(x)|^2 dx$

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Eyring–Kramers law for 1D SPDEs: main result

- ▷ Hessian of V at $\phi_0 = 0$: $-\Delta - 1$, eigenvalues $\mu_k = \left(\frac{2k\pi}{L}\right)^2 - 1$
- ▷ Hessian of V at $\phi_- = -1$: $-\Delta + 2$, eigenvalues $\nu_k = \left(\frac{2k\pi}{L}\right)^2 + 2$

Theorem: Neumann b.c. [B & Gentz, 2013]

- ▷ If $L < \pi - c$ with $c > 0$, then

$$\mathbb{E}^{\phi_-}[\tau_+] = 2\pi \sqrt{\frac{1}{|\lambda_0|\nu_0} \prod_{k=1}^{\infty} \frac{\lambda_k}{\nu_k}} e^{(V[\phi_{ts}] - V[\phi_-])/\varepsilon} [1 + \mathcal{O}(\varepsilon)]$$

- ▷ If $L > \pi + c$, then same formula with extra factor $\frac{1}{2}$ (since 2 saddles) and λ'_k ev at nonconst transition state instead of λ_k

- ▷ Results also for L near π and periodic b.c.
- ▷ Proof uses spectral Galerkin approximation & passage to the limit
- ▷ Prefactor involves a Fredholm determinant:

$$\prod_{k=1}^{\infty} \frac{\lambda_k}{\nu_k} = \det[(-\Delta_{\perp} - 1)(-\Delta_{\perp} + 2)^{-1}] = \det[\mathbb{1} - 3(-\Delta_{\perp} + 2)^{-1}]$$

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2. The two-dimensional case

- ▷ Large-deviation principle: [Hairer & Weber, 2015]
- ▷ Naive computation of prefactor fails:

$$\begin{aligned} \log \prod_{k \in (\mathbb{N}^2)^*} \frac{1 - \left(\frac{L}{|k|\pi}\right)^2}{1 + 2\left(\frac{L}{|k|\pi}\right)^2} &\simeq \sum_{k \in (\mathbb{N}^2)^*} \log \left(1 - \frac{3L^2}{|k|^2 \pi^2}\right) \\ &\simeq - \sum_{k \in (\mathbb{N}^2)^*} \frac{3L^2}{|k|^2 \pi^2} \simeq -\frac{3L^2}{\pi^2} \int_1^\infty \frac{r \, dr}{r^2} = -\infty \end{aligned}$$

- ▷ In fact, the equation needs to be renormalised ($P * \xi$ is distribution!)

Theorem: [Da Prato & Debussche 2003]

Let ξ^δ be a mollification on scale δ of white noise. Then

$$\partial_t \phi = \Delta \phi + \phi - [\phi^3 - 3\epsilon C(\delta)\phi] + \sqrt{2\epsilon} \xi^\delta$$

with $C(\delta) \simeq \log(\delta^{-1})$ admits local solution converging as $\delta \rightarrow 0$

Moreover, $\phi - \sqrt{2\epsilon}(P * \xi^\delta) \in \mathcal{C}([0, T], \mathcal{B}_{p,r}^\alpha) \cap L^p([0, T], \mathcal{B}_{p,r}^s)$

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Computation of the prefactor

- ▷ Consider for simplicity $L < 2\pi \Rightarrow$ transition state in 0
- ▷ Galerkin-truncated renormalised potential

$$V_N = \frac{1}{2} \int_{\mathbb{T}^2} [\|\nabla \phi_N(x)\|^2 - \phi_N(x)^2] dx + \frac{1}{4} \int_{\mathbb{T}^2} :\phi_N(x)^4: dx$$

where $:\phi_N^4: = \phi_N^4 - 6\varepsilon C_N \phi_N^2 + 3\varepsilon^2 C_N^2$, $C_N = \sum_{|k| \leq N} \frac{1}{\lambda_k + 1} \asymp \log(N)$

- ▷ Using variational principles: $\text{cap}(A, B) \simeq \sqrt{\frac{|\lambda_0| \varepsilon}{2\pi}} \prod_{0 < |k| \leq N} \sqrt{\frac{2\pi \varepsilon}{\lambda_k}}$
- ▷ Symmetry argument:

$$\int_{B^c} h_{A,B}(z) e^{-V_N(z)/\varepsilon} dz = \frac{1}{2} \int e^{-V_N(z)/\varepsilon} dz = \frac{1}{2} \mathcal{Z}_N(\varepsilon)$$

- ▷ $\mathcal{Z}_N(\varepsilon) \simeq 2 \prod_{|k| \leq N} \sqrt{\frac{2\pi \varepsilon}{\nu_k}} e^{-V_N(L,0)/\varepsilon}$ where $-V_N(L,0) = \frac{1}{4} L^2 + \frac{3}{2} L^2 C_N \varepsilon$
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$$\prod_{0 < |k| \leq N} \frac{\lambda_k}{\lambda_k + 3} e^{3/\lambda_k} \quad \text{converges since} \quad \log \left[\frac{\lambda_k}{\lambda_k + 3} e^{3/\lambda_k} \right] = \mathcal{O} \left(\frac{1}{|k|^4} \right)$$

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Main result in dimension 2

Theorem: [B, Di Gesù, Weber, 2017]

For $L < 2\pi$, $A \ni \phi_-$, $B \ni \phi_+$ balls in $\|\cdot\|_{H^s}$, $s < 0$, $\exists \mu_N$ prob measures on ∂A :

$$\limsup_{N \rightarrow \infty} \mathbb{E}^{\mu_N}[\tau_B] \leq \frac{2\pi}{|\lambda_0|} \sqrt{\prod_{k \in \mathbb{Z}^2} \frac{|\lambda_k|}{\nu_k} e^{\frac{\nu_k - \lambda_k}{|\lambda_k|}} e^{(V[\phi_{ts}] - V[\phi_-])/\varepsilon} [1 + c_+ \sqrt{\varepsilon}]}$$

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▷ Inverse of prefactor involves Carleman–Fredholm determinant:

$$\det_2(\mathbb{1} + T) = \det(\mathbb{1} + T) e^{-\text{Tr } T}$$

with $T = 3(-\Delta_{\perp} - 1)^{-1}$

$\det_2$ defined whenever T is only Hilbert–Schmidt (true for $d \leq 3$)

▷ [Tsatsoulis & Weber 2018]: Same result for $\mathbb{E}^{\phi_0}[\tau_B]$

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Main result in dimension 2

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3. Allen–Cahn equation in dimension 3

- ▷ Equation needs two counterterms:

$$\partial_t \phi = \Delta \phi + \phi - [\phi^3 - (3\varepsilon C_1(\delta) - 9\varepsilon^2 C_2(\delta))\phi] + \sqrt{2\varepsilon} \xi^\delta$$

with $C_1(\delta) \sim \delta^{-1}$, $C_2(\delta) \sim \log(\delta^{-1})$

- ▷ Solution theories:
 - ◊ Regularity structures [Hairer 2014]
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- ▷ **Metastability**: one expects equivalent result (T is still Hilbert–Schmidt)
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Φ_3^4 measure

- ▷ Potential (before renormalisation):

$$V_\alpha(\phi) = \int_\Lambda \left(\frac{1}{2} \|\nabla \phi(x)\|^2 + \frac{1}{2} \phi(x)^2 + \alpha \phi(x)^4 \right) dx$$

where $\Lambda = (\mathbb{R}/\mathbb{Z})^3 =: \mathbb{T}^3$, $\alpha \geq 0$

- ▷ Definition of Gibbs measure?

$$\nu_\alpha(d\phi) \stackrel{“=”}{=} \frac{1}{Z_\alpha} e^{-V_\alpha(\phi)} \text{ “}d\phi\text{”}$$

- ▷ ν_0 is Gaussian (it's a GFF). For random variable F

$$\mathbb{E}^{\nu_\alpha}[F] = \frac{Z_0}{Z_\alpha} \mathbb{E}^{\nu_0}[F e^{-\alpha \int_\Lambda \phi^4(x) dx}]$$

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The static Φ_3^4 model

Theorem: Potential needs exactly 4 counterterms:

$$V_\alpha(\phi) = \int_\Lambda \left(\frac{1}{2} \|\nabla \phi(x)\|^2 + \frac{1}{2} [1 - \alpha^2 C_N^{(2)}] \phi(x)^2 + \alpha : \phi(x)^4 :_{C_N^{(1)}} + \alpha^2 C_N^{(3)} - \alpha^3 C_N^{(4)} \right) dx$$

where N ($\sim \delta^{-1}$) UV cut-off and

$$C_N^{(1)} = G_N(0) = \text{Tr}((- \Delta_N + 1)^{-1}) = \mathcal{O}(N)$$

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and $G_N(x) = \sum_{|k| \leq N} \frac{1}{\lambda_k + 1} e_k(x)$ is the Green function of Δ_N

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Perturbative computation of partition function

- ▷ Wick powers: $X = \text{diagram of two lines crossing at a point} = \int_{\Lambda} : \phi(x)^4 : dx$, $Y = \text{diagram of two lines connected by a dot} = \int_{\Lambda} : \phi(x)^2 : dx$
- ▷ Potential: $V_{\alpha} = V_0 + \alpha X + \beta Y + \gamma$, $\beta = \frac{1}{2} \alpha^2 C_N^{(2)}$, $\gamma = \alpha^2 C_N^{(3)} - \alpha^3 C_N^{(4)}$
- Then $\frac{Z_{N,\alpha}}{Z_{N,0}} = \mathbb{E}^{\nu_0} [e^{-\alpha X - \beta Y - \gamma}] = e^{-\gamma} \mathbb{E}^{\nu_0} [e^{-\alpha X - \beta Y}]$
- ▷ Let $\Gamma = (\mathcal{V}, \mathcal{E})$ be a multigraph. Its valuation is

$$\Pi_N(\Gamma) = \int_{\Lambda^{\mathcal{V}}} \prod_{e \in \mathcal{E}} G_N(x_{e_+} - x_{e_-}) dx$$

For instance

$$C_N^{(1)} = \Pi_N \text{diagram of a tadpole}$$

“tadpole”

$$C_N^{(2)} = 4^2 3! \Pi_N \text{diagram of two vertices connected by two lines}$$

“sunset”

$$C_N^{(3)} = 4^2 \frac{4!}{2!4^2} \Pi_N \text{diagram of two vertices connected by three lines}$$

“water melon”

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“PSG”

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Cumulant expansion

$$\triangleright \mathbb{E}^{\nu_0} [e^{-\alpha X - \beta Y}] = \sum_{n \geq 0} \frac{1}{n!} \mu_n$$

$$\mu_n = (-1)^n \mathbb{E}^{\nu_0} \left[\left(\alpha \begin{array}{c} \diagup \quad \diagdown \\ \bullet \end{array} + \beta \text{---} \bullet \text{---} \right)^n \right] = (-1)^n \sum_{m=0}^n \binom{n}{m} \alpha^m \beta^{n-m} A_{nm}$$

where $A_{nm} = \mathbb{E}^{\nu_0} \left[\begin{array}{c} \diagup \quad \diagdown \\ \bullet \end{array}^m \text{---} \bullet \text{---}^{n-m} \right]$ computed using Isserlis' thm

Examples: $\mu_2 = \alpha^2 4! \Pi_N \begin{array}{c} \diagup \quad \diagdown \\ \bullet \end{array} + \beta^2 2! \Pi_N \text{---} \bullet \text{---}$

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$$-\log \mathbb{E}[e^{-\alpha X - \beta Y - \gamma}] = \gamma - \sum_{n=2}^{\infty} \frac{\kappa_n}{n!} \quad \kappa_n = \mu_n - \sum_{m=2}^{n-2} \binom{n-1}{m} \kappa_m \mu_{n-m}$$

$\triangleright$ Linked Cluster Theorem: κ_n projection of μ_n on **connected** graphs

Proof: for instance Peccati & Taqqu (2011)

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Proof: for instance Peccati & Taqqu (2011)

Divergences and subdivergences

▷ **Degree** of Γ : $\deg(\Gamma) = 3(|\mathcal{V}| - 1) - |\mathcal{E}|$. Γ **divergent** if $\deg(\Gamma) \leq 0$.

▷ Examples:

$$\deg(\text{bubble}) = 0$$

$$\Pi_N(\text{bubble}) = \mathcal{O}(\log N)$$

$$\deg(\text{figure-eight}) = -1$$

$$\Pi_N(\text{figure-eight}) = \mathcal{O}(N)$$

$$\deg(\text{triangle}) = 0$$

$$\Pi_N(\text{triangle}) = \mathcal{O}(\log N)$$

▷ It looks like $\Pi_N(\Gamma) = \begin{cases} \mathcal{O}(N^{-\deg(\Gamma)}) & \text{if } \deg(\Gamma) < 0 \\ \mathcal{O}(\log N) & \text{if } \deg(\Gamma) = 0 \end{cases}$

However, $\deg(\text{figure-eight with subdivergence}) = 1$, while $\Pi_N(\text{figure-eight with subdivergence}) = \mathcal{O}(\log N)$
because it contains a **subdivergence** 

Theorem: [Bogolyubov, Parasiuk, Hepp, Zimmermann]

$\Pi_N^{\text{BPHZ}}(\Gamma) := -\Pi_N \mathcal{A}(\Gamma) 1_{\deg \Gamma > 0}$ is bdd uniformly in N for any Γ

where $\mathcal{A}$ defined recursively by $\mathcal{A}(\Gamma) = -\Gamma - \sum_{\substack{\bar{\Gamma} \subsetneq \Gamma, \deg(\bar{\Gamma}) \leq 0}} \mathcal{A}(\bar{\Gamma}) \cdot (\Gamma/\bar{\Gamma})$

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Main result: Commutative diagram

$$\begin{array}{ccccc}
 e^{-\alpha X} & \xrightarrow{\mathcal{P}} & \sum_p \frac{(-\alpha)^p}{p!} \mathcal{P}(X^p) & & \\
 \downarrow W & & \downarrow M^{\text{BPHZ}+\Theta} & \searrow \Pi_N^{\text{BPHZ}} + \Pi_N \Theta & \\
 e^{-\alpha X - \beta Y} & \xrightarrow{\mathcal{P}} & \sum_{n,m} \frac{(-\alpha)^m (-\beta)^{n-m}}{m!(n-m)!} \mathcal{P}(X^m Y^{n-m}) & \xrightarrow{\Pi_N} & \log \mathbb{E}[e^{-\alpha X - \beta Y}]
 \end{array}$$

- ▷ $\mathcal{P} = \Pi_{\text{connected}}(\sum_{\text{pairings}})$, satisfies $\deg \mathcal{P}(X^n) = n - 3$
- ▷ W : Wick map, $W(X^n) = H_n(X; \beta Y)$ is Hermite polynomial (for more general models: Bell polynomial)
- ▷ Θ associated with energy renormalisation, $\Pi_N \Theta(e^{-\alpha X}) = \gamma$

Theorem: [B, Klose, Tapia 25]

$$\log \frac{Z_{N,\alpha}}{Z_{N,0}} = \log \mathbb{E}[e^{-\alpha X - \beta Y}] - \gamma \asymp - \sum_{n \geq 4} \frac{(-\alpha)^n}{n!} \Pi_N \mathcal{A}(\mathcal{P}(X^n))$$

as asymptotic expansion with terms uniformly bounded in N

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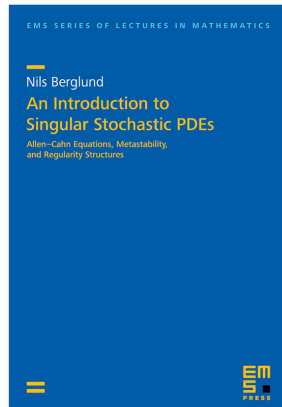
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Thanks for your attention!

Slides available at https://www.idpoisson.fr/berglund/Athens26_Feynman.pdf