

Convergence to equilibrium in some singular parabolic SPDEs

Based on: NB & B. Gentz, EJP (2013), NB, G. Di Gesu & H. Weber, EJP (2017)
See also arXiv: 1901.07420

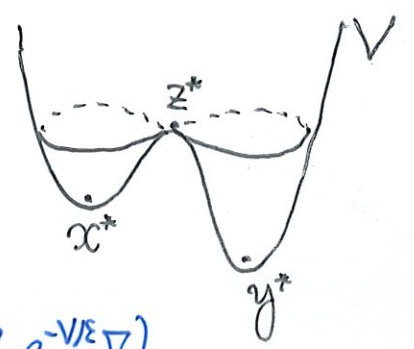
Allen-Cahn: $\partial_t \phi - \Delta \phi = \phi - \phi^3 + \sqrt{\varepsilon} \xi^\delta + C_{\delta, \varepsilon} \phi$ $\phi(t, x), t \geq 0, x \in (\mathbb{R}/\mathbb{Z})^d$

Φ_d^4 model: $\partial_t \phi - \Delta \phi = -m^2 \phi - \phi^3 + \sqrt{\varepsilon} \xi^\delta$ $m^2 \geq 0$

Existence/uniqueness of sol. (local/global)	Invariant measure	Large deviation principle	Eyring-Kramers law
d=1: Faris & Jona-Lasinio 82	Da Prato & Zabczyk 96	Faris & Jona-Lasinio 82	Berglund & Gentz 13 Barret 15
d=2: Da Prato & Debussche 03 Tsatsoulis & Weber 18	Röckner, Zhu & Zhu 17 Tsatsoulis & Weber 18 Hairer & Mattingly 18	Jona-Lasinio & Mitter 90 Hairer & Weber 15	B, Di Gesu & Weber 17 Tsatsoulis & Weber 18
d=3: Hairer 14, Catellier & Chauk 18, Kupiainen 18		Hairer & Weber 15	?

1. Reversible SDE

$$dx_t = -\nabla V(x_t) dt + \sqrt{\varepsilon} dW_t \quad x \in \mathbb{R}^n$$



Inv. measure: $\pi(x) dx = \frac{1}{Z} e^{-V(x)} dx$

reversible since $\langle f, \mathcal{L}g \rangle_\pi = \langle \mathcal{L}f, g \rangle_\pi$ ($\mathcal{L} = \varepsilon e^{V/\varepsilon} \nabla \cdot e^{-V/\varepsilon} \nabla$)

Convergence to π :

(a) [Meyn & Tweedie 93] $(\mathcal{L}\tilde{V})(x) \leq -c\tilde{V}(x) + d$ $c > 0$ $d \geq 0$
 $\Rightarrow \sup_{|f| \leq V+1} |E^x(F(x_t)) - \langle \pi, f \rangle| \leq C[V(x)+1] e^{-\beta t} \quad \forall x$
 $\beta, C > 0$

(b) Expected transition time: $E^{x^*}[\tau_{B_r}(y^*)] = E = ?$ } link with (a):
 [Day 83] $P^{x^*}(\tau_{B_r}(y^*) > tE) \xrightarrow{\varepsilon \rightarrow 0} e^{-t}$ } $f = \mathbb{1}_{B_r}(y^*)$

Behaviour of E : -Arrhenius [Fredlin & Wentzell 70] $\lim_{\varepsilon \rightarrow 0} \varepsilon \log E = V(z^*) - V(x^*)$
 via LDP

- Eyring-Kramers [Bovier et al 2004] $E = \frac{2\pi}{|\lambda_{-}(z^*)|} \sqrt{\frac{|\det \text{Hess } V(z^*)|}{|\det \text{Hess } V(x^*)|}} e^{[V(z^*) - V(x^*)]/\varepsilon}$
 via potential theory [1+o(1)]

(c) Spectral theory: [Bovier et al 2005] (spectral gap of \mathcal{L}) = $\frac{1}{E} (1+O(\varepsilon^{-ck}))$

2. Allen-Cahn in dim $d=1$

$$\partial_t \phi - \Delta \phi = \phi - \phi^3 + \sqrt{2\varepsilon} \xi$$

$$V(\phi) = \int_0^L \left[\frac{1}{2} \phi'(x)^2 + \frac{1}{4} \phi(x)^4 - \frac{1}{2} \phi(x)^2 \right] dx$$

$$\Rightarrow \frac{\partial}{\partial \lambda} V(\phi + \lambda \psi) \Big|_{\lambda=0} = - \langle \Delta \phi + \phi - \phi^3, \psi \rangle_{L^2}$$

\Rightarrow Gradient,
Inv. measure $\sim \frac{e^{-V/\varepsilon}}{Z}$
(abs. cont. w.r.t. GFF)

- Stationary sol:
- $\phi \equiv \pm 1$ (loc. min. of V , stable)
 - $\phi \equiv 0$ (transition state if $L < 2\pi$)
 - other, non-const. stat. states if $L > 2\pi$ \otimes LDP

Hessian: $V(\phi) = \frac{1}{2} \langle \phi, (-\Delta - 1)\phi \rangle + O(\phi^4) \Rightarrow HV(0) = -\Delta - 1 \quad \text{ev} \left(\frac{2\pi k}{L} \right)^2 - 1$
 $HV(\pm 1) = -\Delta + 2 \quad \text{ev} \left(\frac{2\pi k}{L} \right)^2 + 2$

In E-k prefactor: $\det \left(\begin{matrix} \downarrow \\ \Delta \text{ acting on mean zero fct} \\ [-\Delta_L - 1] [-\Delta_L + 2]^{-1} \end{matrix} \right) = \det \left([-\Delta_L + 2 - 3] [-\Delta_L + 2]^{-1} \right)$
 $= \det \left(1 - 3[-\Delta_L + 2]^{-1} \right)$ Fredholm det
 $\log \det \left([-\Delta_L - 1] [-\Delta_L + 2]^{-1} \right) = \text{Tr} \log \left(1 - 3[-\Delta_L + 2]^{-1} \right)$
 $= - \sum_{n \geq 1} \frac{3^n}{n} \text{Tr} \left([-\Delta_L + 2]^{-n} \right) < \infty$
 $= \sum_{k \neq 0} \frac{1}{\left[\left(\frac{2\pi k}{L} \right)^2 + 2 \right]^n} \leq \frac{\text{const}}{\left[\left(\frac{2\pi}{L} \right)^2 + 2 \right]^n}$

Thm: [B, Gentz 13] EK law holds true
(result also exists for $L \geq 2\pi$, more general equ.)

Proof uses spectral Galerkin approx.

\otimes LDP [F&JL'82]:

LDP with cube fct $\mathcal{I}_{[0,T]}(\gamma) = \int_0^T \int_0^L \left[\frac{\partial}{\partial t} \gamma(t,x) - \Delta \gamma(t,x) - (\phi - \phi^3)(t,x) \right]^2 dx dt$
 $+ \infty$

\Rightarrow Arrhenius law holds $\lim_{\varepsilon \rightarrow 0} \varepsilon \log \mathbb{E}^{-1} [\tau_{\mathbb{R}^{(+1)}}] = V(0) - V(-1) = \frac{L}{4}$

$\mathbb{P}^{-1} \{ \tau_{\mathbb{R}^{(+1)}} < T \} = \inf \{ \mathcal{I}_{[0,T]}(\gamma) : \gamma(0) = -1, \gamma(T) \in \mathbb{R}^{(+1)} \}$
 $\mathcal{I}_{[0,T]}(\gamma) = \frac{1}{2} \int_0^T \int_0^L \left[\frac{\partial}{\partial t} \gamma + \Delta \gamma + (\phi - \phi^3) \right]^2 dx dt + 2 \int_0^T \int_0^L \left[-\Delta \gamma - (\phi - \phi^3) \right] \frac{\partial}{\partial t} \gamma dx dt$
 $\geq 2 \int_0^T \int_0^L \left[\frac{\partial}{\partial t} \gamma \frac{\partial}{\partial t} \gamma - (\phi - \phi^3) \frac{\partial}{\partial t} \gamma \right] dx dt$
 $= 2 [V(\gamma_T) - V(\gamma_0)]$

3. Allen-Cahn in dim $d=2$

$\partial_t \phi - \Delta \phi = \phi - \phi^3 + 3\varepsilon C_\delta \phi + \sqrt{2\varepsilon} \xi^S \Rightarrow$ Stat. sol in $\pm \sqrt{1+3\varepsilon C_\delta}$??

$V(\phi) = \int_{\mathbb{T}^2} \left[\frac{1}{2} |\nabla \phi|^2 + \frac{1}{4} \phi^4 - \frac{1}{2} \phi^2 - \frac{3}{2} \varepsilon C_\delta \phi^2 \right] dx \Rightarrow$ inv. meas. $\sim \frac{e^{-V/\varepsilon}}{\varepsilon}$
(can be written using ϕ^4, ϕ^3)

[Hairer & Weber 15] LDP with $I_{\text{opt}}(\gamma) = \frac{1}{2} \int_0^T \int_{\mathbb{T}^2} \left[\frac{\partial \gamma}{\partial t} - \Delta \gamma - (\phi - \phi^3) \right]^2 dx dt$

\Rightarrow does not depend on $C_\delta \rightsquigarrow$ Stat sol in ± 1 ?

Arrhenius law $\lim_{\varepsilon \rightarrow 0} \varepsilon \log \mathbb{E}^{-1}[\tau_{\mathbb{R}^{(1)}}] = V(0) - V(-1) \Big|_{\varepsilon=0} = \frac{L^2}{4}$

Eyring-Kramers prefactor: $\text{Tr} [(-\Delta_1 + 2)^{-n}] \cong \sum_{k \neq (0,0)} \frac{1}{\|k\|^{2n}} \sim \int_1^\infty \frac{r dr}{r^{2n}} \begin{cases} = +\infty & n=1 \\ < +\infty & n \geq 2 \end{cases}$

$\Rightarrow (-\Delta_1 + 2)^{-1}$ is not trace class, but Hilbert-Schmidt

Spectral Galerkin: $\partial_t \phi_N - \Delta \phi_N = \phi_N - P_N \phi_N^3 + 3\varepsilon C_N \phi_N + \sqrt{2\varepsilon} \xi_N$

$\phi_N^{\text{GFF}}(x) = \sum_{|k| \leq N} \frac{z_k}{\sqrt{|k|+1}} e_k(x) \Rightarrow \mathbb{E}[\phi_N^{\text{GFF}}(x) \phi_N^{\text{GFF}}(y)] = \sum_{|k| \leq N} e_k(x) (-\Delta + 1)^{-1} e_k(y)$

$L^2 C_N = \mathbb{E}[\|\phi_N^{\text{GFF}}\|_{L^2}^2] = \text{Tr} [(-\Delta_N + 1)^{-1}] \sim \log N$

Eyring-Kramers: $V(0) - V(-1) = \frac{L^2}{4} + \frac{3}{2} L^2 \varepsilon C_N$

$\Rightarrow \sqrt{\det [(-\Delta_N^1 + 1)^{-1} (-\Delta_N^1 + 2)]} e^{(V(0)-V(-1))/\varepsilon} = \left[\det (1 - 3(-\Delta_N^1 + 2)^{-1}) e^{3 \text{Tr} [(-\Delta_N^1 + 2)^{-1}] + \Theta} \right]^{1/2} e^{L^2/4\varepsilon}$
Carleman-Fredholm det

$\log[\dots] = \text{Tr} \left[\log (1 - 3(-\Delta_N^1 + 2)^{-1}) + 3(-\Delta_N^1 + 2) \right] + \Theta$
 $= \sum_{0 \neq |k| < N} \left(\log \left(1 - \frac{3}{(\frac{2\pi|k|}{L})^2 + 2} \right) + \frac{3}{(\frac{2\pi|k|}{L})^2 + 2} \right) + \Theta < +\infty$
shift in reference GFF

[B, Di Gesù, Weber 17]: uniform in N bounds on EK law (starting in eq. measure)
[Tzatsoulis & Weber 18]: full proof of EK law (limit $N \rightarrow \infty$ & arbitrary initial cond)