IECN Nancy Hopf algebras, operads, deformations for singular dynamics

Perturbation theory for the Φ_3^4 measure, revisited with Hopf algebras

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Joint works with Tom Klose (Berlin) and Yvain Bruned (Nancy)









https://www.idpoisson.fr/berglund/

Nils Berglund

The Φ_d^4 model

▷ Lattice system: $\Lambda_N = (\mathbb{Z}/N\mathbb{Z})^d$, $y \in \mathbb{R}^{\Lambda_N}$

$$V_{N,\varepsilon}(y) = \frac{1}{2}N^2 \sum_{\substack{i,j \in \Lambda \\ \|i-j\|=1}} (y_i - y_j)^2 + \sum_{i \in \Lambda} U_{\varepsilon}(y_i)$$

$$U_{\varepsilon}(\xi) = \frac{1}{2}\xi^2 + \frac{\varepsilon}{4}\xi^4$$

Gibbs measure
$$\mu_{N,\varepsilon}(dy) = \frac{1}{Z_{N,\varepsilon}} e^{-V_{N,\varepsilon}(y)} dy$$

▷ Continuum limit:
$$y_i = \phi(i/N), N \to \infty,$$

 $V_{\varepsilon}(\phi) = \int_{\Lambda} \left(\frac{1}{2} \|\nabla \phi(x)\|^2 + \frac{1}{2} \phi(x)^2 + \frac{\varepsilon}{4} \phi(x)^4\right) dx$

where $\Lambda = (\mathbb{R}/\mathbb{Z})^d =: \mathbb{T}^d$

Definition of Gibbs measure?

$$\mu_{\varepsilon}(\mathsf{d}\phi) = \frac{1}{Z_{\varepsilon}} e^{-V_{\varepsilon}(\phi)} d\phi$$

 \triangleright Alternative: Spectral Galerkin approx. (Fourier modes with $|k| \leq N$)

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The case d = 1 $\triangleright \varepsilon = 0$: $V_0(\phi) = \int_{\Lambda} \left(\frac{1}{2} \|\nabla \phi(x)\|^2 + \frac{1}{2} \phi(x)^2\right) dx = \frac{1}{2} \langle \phi, (-\Delta + 1)\phi \rangle$ μ_0 is Gaussian free field with covariance $(-\Delta + 1)^{-1}$ (well-defined since $(-\Delta + 1)^{-1}$ trace class: $\lambda_k = (2\pi k)^2$, $\sum_{k \in \mathbb{Z}} \frac{1}{\lambda_k + 1} < \infty$)

 $\triangleright \varepsilon > 0$

$$\frac{\mathrm{d}\mu_{\varepsilon}}{\mathrm{d}\mu_{0}} = \frac{Z_{0}}{Z_{\varepsilon}} \,\mathrm{e}^{-\left[V_{\varepsilon} - V_{0}\right]} = \frac{Z_{0}}{Z_{\varepsilon}} \,\mathrm{e}^{-\frac{\varepsilon}{4}\int_{\Lambda}\phi(x)^{4}\,\mathrm{d}x}$$

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$$\frac{Z_{\varepsilon}}{Z_{0}} = \mathbb{E}^{\mu_{0}} \left[e^{-\frac{\varepsilon}{4} \int_{\Lambda} \phi(x)^{4} dx} \right] = \frac{1}{Z_{0}} \int e^{-V_{0}(\phi) - \frac{\varepsilon}{4} \int_{\Lambda} \phi(x)^{4} dx} d\phi$$

Fourier representation:

$$\phi_{\mathsf{GFF}}(x) = \sum_{k \in \mathbb{Z}} \frac{Z_k}{\sqrt{\lambda_k + 1}} e_k(x), \quad Z_k \sim \mathcal{N}(0, 1) \text{ iid}$$

$$\Rightarrow \mathbb{E} \Big[\int_{\Lambda} \phi_{\mathsf{GFF}}(x)^{2n} dx \Big] \lesssim \Big(\sum_{k \in \mathbb{Z}} \frac{1}{\lambda_k + 1} \Big)^n < C^n$$

Renormalisation of static and dynamic Φ_d^4 models

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so that
$$\frac{Z_{\varepsilon}}{Z_0} = 1 + \mathcal{O}(\varepsilon)$$

 $\triangleright (-\Delta + 1)^{-1}$ no longer trace class, since $\lambda_k = (2\pi ||k||)^2$, $\sum_{k \in \mathbb{Z}^2} \frac{1}{\lambda_k + 1} = \infty$

▷ Truncated GFF:

$$\phi_{\mathsf{GFF},N}(x) = \sum_{k \in \mathbb{Z}^2, |k| \leq N} \frac{\mathbb{Z}_k}{\sqrt{\lambda_k + 1}} e_k(x)$$

$$C_N = \int_{\Lambda} \mathbb{E}[\phi_{\mathsf{GFF},N}(x)^2] \, \mathrm{d}x = \sum_{|k| \le N} \frac{1}{\lambda_k + 1} = \mathsf{Tr}[(-\Delta_N + 1)^{-1}] \sim \log N$$

▷ Wick calculus: $:\phi(x)^n := H_n(\phi(x); C_N)$ where H_n Hermite polynomials If (X, Y) centred jointly Gaussian rv, $\mathbb{E}[X^2] = C$, $\mathbb{E}[Y^2] = C'$ then $\mathbb{E}[H_n(X; C)H_m(Y; C')] = n!\delta_{nm}\mathbb{E}[XY]^n$

Consequence: $\sup_{N} \mathbb{E} \left[\left(\int_{\Lambda} : \phi_{\mathsf{GFF},N}(x)^{n} : \mathrm{d}x \right)^{2} \right] < \infty \qquad \forall n$

 $\triangleright~$ Gibbs measure defined as in 1d case, with

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Renormalisation of static and dynamic Φ_d^4 models

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Renormalisation of static and dynamic Φ_d^4 models

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Theorem: Potential needs exactly 4 counterterms:

$$V(\phi) = \int_{\Lambda} \left(\frac{1}{2} \| \nabla \phi(x) \|^2 + \frac{1}{2} \left[1 - \varepsilon^2 C_N^{(2)} \right] \phi(x)^2 + \frac{\varepsilon}{4} \cdot \phi(x)^4 \cdot \frac{\varepsilon}{C_N^{(1)}} + \varepsilon^2 C_N^{(3)} - \varepsilon^3 C_N^{(4)} \right) dx$$

where

$$C_{N}^{(1)} = G_{N}(0) = \operatorname{Tr}((-\Delta_{N} + 1)^{-1}) = \mathcal{O}(N)$$

$$C_{N}^{(2)} = 3! \int_{\Lambda} G_{N}(x)^{3} dx = \mathcal{O}(\log N)$$

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$$C_{N}^{(4)} = \frac{2^{3}}{3!4^{3}} {\binom{4}{2}}^{3} \int_{\Lambda} \int_{\Lambda} G_{N}(x)^{2} G_{N}(y)^{2} G_{N}(x - y)^{2} dx dy = \mathcal{O}(\log N)$$

and
$$G_N(x) = \sum_{|k| \leq N} \frac{1}{\lambda_k + 1} e_k(x)$$
 is the Green function of Δ_N

Renormalisation of static and dynamic Φ_d^4 models

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Renormalisation of static and dynamic Φ_d^4 models

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Some literature

- Glimm & Jaffe (1968, 1973), Feldman (1974): Combinatorics of Feynman diagrams
- Benfatto, Cassandro, Gallavotti, Nicolò & Olivieri (1978, 1980): Renormalisation group (integrating out scales)
- Brydges, Fröhlich & Sokal (1983):
 Generating function and skeleton inequalities
- Brydges, Dimock & Hurd (1995):
 Polymer expansions
- Connes & Kreimer (2000, 2001): Hopf algebras

▷ ...

Barashkov & Gubinelli (2020):
 Boué–Dupuis formula

Singular stochastic PDEs

$$\partial_t \phi(t,x) = \Delta \phi(t,x) - \phi(t,x)^3 + \xi(t,x)$$

space-time white noise

- Parisi & Wu (1981):
 Stochastic quantization
- Faris & Jona-Lasinio (1982), ...:
 1d case: Well-posed, large-deviation principle
- ⊳ Da Prato & Debussche (2003):

2d case: Besov spaces, fixed-point argument for difference between ϕ and stochastic convolution

▷ Hairer (2014):

3d case: regularity structures, Banach spaces of modeled distributions Ad-hoc renormalisation for Φ_3^4 and PAM (parabolic Anderson model)

 Bruned, Chandra, Chevyrev, Hairer, Zambotti (2016+):
 Solution theory and renormalisation for general locally subcritical (superrenormalisable) parabolic SPDEs, using BPHZ renormalisation

Graphical notations

▷ Wick powers: $X = \sum_{N=0}^{\infty} = \int_{\Lambda} :\phi(x)^4 : dx, Y = ---- = \int_{\Lambda} :\phi(x)^2 : dx$ ▷ Parameters: $\alpha = \frac{\varepsilon}{4}, \beta = \frac{1}{2}\varepsilon^2 C_N^{(2)}, \gamma = \varepsilon^2 C_N^{(3)} - \varepsilon^3 C_N^{(4)}$ Then $\frac{Z_{N,\varepsilon}}{Z_{N,0}} = \mathbb{E}^{\mu_0} \left[e^{-\alpha X - \beta Y - \gamma} \right] = e^{-\gamma} \mathbb{E}^{\mu_0} \left[e^{-\alpha X - \beta Y} \right]$

▷ Let $\Gamma = (\mathscr{V}, \mathscr{E})$ be a multigraph, $\mathscr{G} = \operatorname{span}{\Gamma}$. Its valuation is

$$\Pi_N(\Gamma) = \int_{\Lambda^{\mathscr{V}}} \prod_{e \in \mathscr{E}} G_N(x_{e_+} - x_{e_-}) \, \mathrm{d}x$$

For instance

$$C_{N}^{(1)} = \Pi_{N} \bigcirc$$

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Renormalisation of static and dynamic Φ_d^4 models

Cumulant expansion

Examples:

$$\mu_{2} = \alpha^{2} 4! \Pi_{N} \bigoplus + \beta^{2} 2! \Pi_{N} \bigoplus$$
$$\mu_{3} = -\alpha^{3} {\binom{4}{2}}^{3} 2^{3} \Pi_{N} \bigoplus - 3\alpha^{2} \beta (4^{2} \cdot 2 \cdot 3!) \Pi_{N} \bigoplus$$
$$-3\alpha\beta^{2} 4! \Pi_{N} \bigoplus -8\beta^{3} \Pi_{N} \bigwedge$$

Cumulant expansion: (Leonov & Shiraev)

$$-\log \mathbb{E}[e^{-\alpha X - \beta Y - \gamma}] = \gamma - \sum_{n=2}^{\infty} \frac{\kappa_n}{n!} \qquad \kappa_n = \mu_n - \sum_{m=2}^{n-2} \binom{n-1}{m} \kappa_m \mu_{n-m}$$

Linked Cluster Theorem: κ_n projection of μ_n on connected graphs
 Proof: for instance Peccati & Taqqu (2011)

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Cumulant expansion

$$\models \mu_n = (-1)^n \mathbb{E}^{\mu_0} \Big[\Big(\alpha \longrightarrow + \beta \longrightarrow \Big)^n \Big] = (-1)^n \sum_{m=0}^n \binom{n}{m} \alpha^m \beta^{n-m} A_{nm}$$

where $A_{nm} = \mathbb{E}^{\mu_0} \Big[\longrightarrow \stackrel{m \longrightarrow n-m}{\longrightarrow} \Big]$

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▷ Degree of Γ: deg(Γ) = 3(|𝒴/ − 1) − |𝔅|. Γ divergent if deg(Γ) ≤ 0.
 ▷ Examples:



Theorem: [Dyson]

If deg $\overline{\Gamma} > 0$ for all subgraphs $\overline{\Gamma} \subset \Gamma$, then $\Pi_N(\Gamma)$ is bounded unif in N

Renormalisation of static and dynamic Φ_d^4 models

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Renormalisation of static and dynamic Φ_d^4 models

▷ Degree of Γ : deg $(\Gamma) = 3(|\mathcal{V}| - 1) - |\mathcal{E}|$. Γ divergent if deg $(\Gamma) \leq 0$.

▷ Examples:



Theorem: [Dyson]

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Hopf algebras and renormalisation

- ▷ Character: linear form $g : \mathscr{G} \to \mathbb{R}$ such that $\langle g, \Gamma_1 \cdot \Gamma_2 \rangle = \langle g, \Gamma_1 \rangle \langle g, \Gamma_2 \rangle$ Renormalisation map: $M^g : \mathscr{G} \to \mathscr{G}, M^g(\Gamma) := (g \otimes \mathrm{id}) \Delta \Gamma$ Property: If $\langle f \circ g, \Gamma \rangle = \langle f \otimes g, \Delta \Gamma \rangle$ and $\langle \mathscr{A}^*(f), \Gamma \rangle = \langle f, \mathscr{A}(\Gamma) \rangle$ then $M^{g \circ h} = M^g M^h$ and $(M^g)^{-1} = M^{\mathscr{A}^*(g)} \Rightarrow$ group structure

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Renormalisation of static and dynamic Φ_d^4 models

- ▷ BPHZ character: $\langle g^{\text{BPHZ}}, \Gamma \rangle = \prod_N \mathscr{A}(\Gamma) \mathbf{1}_{\text{deg } \Gamma \leqslant 0}$
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 $\Pi_{N}^{\mathsf{BPHZ}}(\Gamma) = \Pi_{N} \mathcal{M}^{\mathcal{G}^{\mathsf{BPHZ}}}(\Gamma)$ $= (g^{\mathsf{BPHZ}} \otimes \Pi_{N}) \Delta \Gamma$ $= (\Pi_{N} \tilde{\mathscr{A}} \otimes \Pi_{N}) \Delta \Gamma$ $\tilde{\mathscr{A}}(\Gamma) = \mathscr{A}(\Gamma) 1_{\mathsf{deg} \Gamma \leqslant 0}$



Theorem: [Bogolyubov, Parasiuk, Hepp, Zimmermann] If deg $\Gamma > 0$ then $\Pi^{BPHZ}(\Gamma)$ hdd uniformly in M

Theorem: [B & Klose]

Write $\kappa_n = (-1)^n \sum_{m=0}^n {n \choose m} \alpha^m \beta^{n-m} \sum_k b_{nm}^{(k)} \Pi_N(\Gamma_{nm}^{(k)})$ Then $\sum_{n=2}^\infty \frac{\kappa_n}{n!} = -\sum_{p=2}^\infty \frac{1}{p!} (-\alpha)^p \sum_k b_{pp}^{(k)} \Pi_N^{\mathsf{BPHZ}}(\Gamma_{pp}^{(k)}) \quad \deg \Gamma_{pp}^{(k)} = p-3$

Consequence: all terms in cumulant expansion bounded uniformly in N

Renormalisation of static and dynamic Φ_d^4 models

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Commutative diagram



 $\triangleright \mathcal{P} = \prod_{\text{connected}} \left(\sum_{\text{pairings}} \right)$

 $\triangleright \ e^{-\alpha X}, e^{-\alpha X - \beta Y} \in H = \operatorname{span} \left\{ \mathsf{X}^{\boldsymbol{n}} : \ \boldsymbol{n} \in \mathbb{N}^2 \right\} \qquad \mathsf{X}^{\boldsymbol{n}} \coloneqq \mathsf{X}^{\boldsymbol{n}_1} Y^{\boldsymbol{n}_2}$

 \triangleright Construction of $\hat{\chi}$ inspired by Ebrahimi-Fard et al

Lemma: [B & Klose]

$$\triangleright \hat{\chi}(e^{-\alpha X}) = e^{-\alpha X - \beta Y}$$

$$\triangleright \mathscr{P} \circ \hat{\chi} = M^{g^{\mathsf{BPHZ}}} \circ \mathscr{P}$$

Proof of commutativity based on Zimmermann's forest formula for A

Commutative diagram



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Proof of commutativity based on Zimmermann's forest formula for \mathscr{A}

Ebrahimi-Fard et al type construction



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▷ Zimmermann forest formula: $\mathscr{A}(\Gamma) = -\sum_{\mathscr{X}} (-1)^{|\mathscr{F}|} \mathscr{C}_{\mathscr{F}}\Gamma$

where sum ranges over all forests \mathscr{F} (set of subgraphs, pairwise vertex-disjoint or included) and $\mathscr{C}_{\mathscr{F}}$ extracts all subgraphs in \mathscr{F}

▷ Our case:
$$\mathscr{A}(\Gamma) = -\sum_{S \subset \{1,...,g\}} (-1)^{|S|} \bigoplus^{|S|} \mathscr{C}_S \Gamma$$

where \mathscr{C}_S extracts all "bubbles" \longleftrightarrow labelled by element of S , g is number of bubbles

Since $\Pi_N(\bigcirc) = \frac{\beta}{3\varepsilon^2} = \frac{\beta}{48\alpha^2}$ we have

$$\Pi_{N}^{\mathsf{BPHZ}}\left(\Gamma_{pp}^{(k)}\right) = -\sum_{S \subset \{1, \dots, g\}} \left(-\frac{\beta}{48\alpha^{2}}\right)^{|S|} \Pi_{N}\left(\mathscr{C}_{S}\Gamma_{pp}^{(k)}\right)$$

Lemma: [B & Klose] The diagram commutes, that is, $\mathcal{P} \circ \mathcal{M} \circ \chi_{\eta} = (\prod_{N} \tilde{\mathcal{A}} \otimes id) \Delta \circ \mathcal{P}$

Renormalisation of static and dynamic Φ_d^4 models

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- $\bigvee V(\phi) = \frac{1}{2}\phi^2 + \frac{\varepsilon}{4}\phi^4$ $Z(\varepsilon) = \int_{-\infty}^{\infty} e^{-V(\phi)} d\phi = \int_{-\infty}^{\infty} e^{-\phi^2/2} e^{-\varepsilon\phi^4/4} d\phi$ $Z(\varepsilon) \asymp \sqrt{2\pi} \sum_{n \ge 0} \left(-\frac{\varepsilon}{4}\right)^n \frac{(4n-1)!!}{n!} = \sum_{n \ge 0} a_n \varepsilon^n, \qquad a_n \sim n!$
- ▷ Borel transform:

$$Z(\varepsilon) \asymp \sum_{n \ge 0} a_n \varepsilon^n \frac{\Gamma(n+1)}{n!} = \sum_{n \ge 0} \frac{a_n \varepsilon^n}{n!} \int_0^\infty t^n e^{-t} dt$$
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Theorem (Watson 1912, Sokal 1980) $D_R = \{\varepsilon: \operatorname{Re} \varepsilon^{-1} > R^{-1}\}$ If Z analytic in D_R and $Z(\varepsilon) = \sum_{k=0}^n a_k \varepsilon^k + R_n(\varepsilon)$ with $|R_n(\varepsilon)| \leq Cr^n n! |\varepsilon|^n$ unif in n and ε , then $\mathcal{B}Z(t)$ cv for $|t| < \frac{1}{r}$ and $Z(\varepsilon) = Z_{\operatorname{Borel}}(\varepsilon)$ in D_R

Renormalisation of static and dynamic Φ_d^4 models

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- Moment bound
- ♦ Sharp estimates on $(\Pi_N^{\text{BPHZ}} \circ \mathscr{P})\Gamma_{PP}^{(k)}$ (Hairer 2018, B & Bruned 2019)

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Control remainder by using

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Thanks for your attention!

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