

Online probability seminar, Seoul (Online)

Metastable dynamics of one-dimensional Allen–Cahn-type equations

Nils Berglund

Institut Denis Poisson, University of Orléans, France



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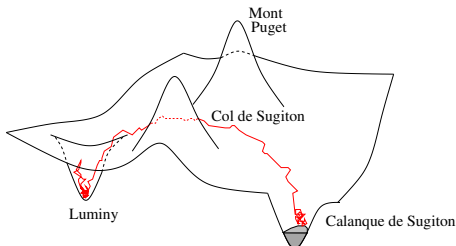
Joint works with Bastien Fernandez (Paris) and Barbara Gentz (Bielefeld)



Project
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1. Reversible SDEs:

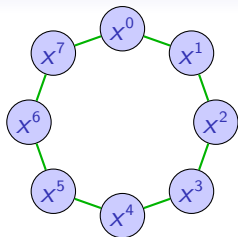
The potential-theoretic approach to metastability



A particle system

[B, Fernandez, Gentz, Nonlinearity 2007]

- ▷ N particles on a circle $\mathbb{Z}/N\mathbb{Z}$
- ▷ Bistable local dynamics
- ▷ Ferromagnetic nearest neighbour coupling
- ▷ Independent noise on each site

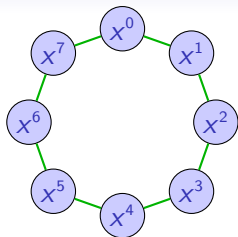


$$dx_t^i = [x_t^i - (x_t^i)^3] dt + \frac{\gamma}{2} [x_t^{i+1} - 2x_t^i + x_t^{i-1}] dt + \sqrt{2\varepsilon} dW_t^i$$

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Gradient system $dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$

$$\text{potential } V(x) = \sum_i U(x^i) + \frac{\gamma}{4} \sum_i (x^{i+1} - x^i)^2 \quad U(\xi) = \frac{1}{4}\xi^4 - \frac{1}{2}\xi^2$$

Coarsening dynamics with noise

([Link to simulation](#))

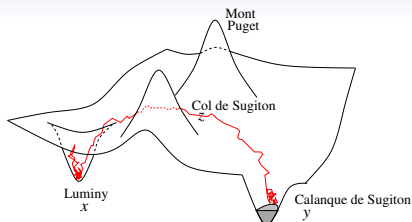
Reversible diffusion in a double-well

$$dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$$

$V : \mathbb{R}^d \rightarrow \mathbb{R}$ confining potential

$$\tau_y^x = \inf\{t > 0 : x_t \in \mathcal{B}_\varepsilon(y)\}$$

first-hitting time of small ball $\mathcal{B}_\varepsilon(y)$,
when starting in x



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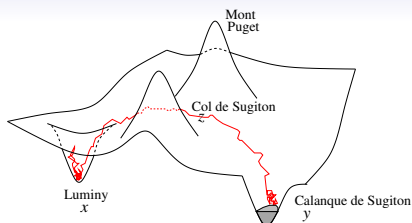
Arrhenius' law (1889): $\mathbb{E}[\tau_y^x] \simeq e^{[V(z)-V(x)]/\varepsilon}$

Eyring–Kramers law (1935, 1940):

Eigenvalues of Hessian of V at minimum x : $0 < \nu_1 \leq \nu_2 \leq \dots \leq \nu_d$

Eigenvalues of Hessian of V at saddle z : $\lambda_1 < 0 < \lambda_2 \leq \dots \leq \lambda_d$

$$\mathbb{E}[\tau_y^x] = 2\pi \sqrt{\frac{\lambda_2 \dots \lambda_d}{|\lambda_1| \nu_1 \dots \nu_d}} e^{[V(z)-V(x)]/\varepsilon} [1 + \mathcal{O}_\varepsilon(1)]$$



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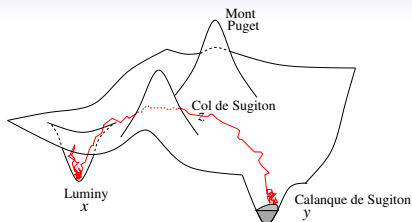
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Arrhenius' law: proved by [Freidlin, Wentzell, 1979] using large deviations

Eyring–Kramers law: [Bovier, Eckhoff, Gayard, Klein, 2004] using potential theory,
[Helffer, Klein, Nier, 2004] using Witten Laplacian, ...



Potential-theoretic proof

$$dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$$

- ▷ **Generator:** $\mathcal{L} = \varepsilon \Delta - \nabla V \cdot \nabla = \varepsilon e^{V/\varepsilon} \nabla \cdot e^{-V/\varepsilon} \nabla$
- ▷ **Invariant probability:** $\pi(dx) = \frac{1}{Z} e^{-V(x)/\varepsilon} dx \Rightarrow \mathcal{L}^\dagger \pi = 0$

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- ▷ **Invariant probability:** $\pi(dx) = \frac{1}{Z} e^{-V(x)/\varepsilon} dx \Rightarrow \mathcal{L}^\dagger \pi = 0$
- ▷ **Reversible:** $\langle f, \mathcal{L}g \rangle = \langle \mathcal{L}f, g \rangle$ for $\langle f, g \rangle = \int_{\mathbb{R}^d} e^{-V(x)/\varepsilon} f(x)g(x) dx$
- ▷ **Dirichlet form:** $\mathcal{E}(f) = \langle f, -\mathcal{L}f \rangle = \varepsilon \int_{\mathbb{R}^d} e^{-V(x)/\varepsilon} |\nabla f(x)|^2 dx$
 $\mathcal{E}(f, g) = \langle f, -\mathcal{L}g \rangle$

Expected hitting time

▷ Expected hitting time:

$$w_A(x) = \mathbb{E}^x[\tau_A] \quad \text{satisfies} \quad \begin{cases} (\mathcal{L}w_A)(x) = -1 & x \in A^c \\ w_A(x) = 0 & x \in A \end{cases}$$

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- ▷ Green function:

$$\begin{cases} (\mathcal{L}G_A)(x) = \delta(x-y) & x \in A^c \\ G_A(x,y) = 0 & x \in A \end{cases}$$

$$\Rightarrow \quad w_A(x) = - \int_{A^c} G_A(x,y)(dy)$$

Committer

▷ Committor:

$$h_{AB}(x) = \mathbb{P}^x\{\tau_A < \tau_B\} \quad \text{satisfies} \quad \begin{cases} (\mathcal{L}h_{AB})(x) = 0 & x \in (A \cup B)^c \\ h_{AB}(x) = 1 & x \in A \\ h_{AB}(x) = 0 & x \in B \end{cases}$$

▷ Equilibrium measure: $e_{AB}(dx) = (-\mathcal{L}h_{AB})(dx)$ measure on $x \in \partial A$

$$\Rightarrow \quad h_{AB}(x) = - \int_A G_B(x, y) e_{AB}(dy)$$

Capacity

Capacity: $\text{cap}(A, B) = \int_{\partial A} e^{-V(x)/\varepsilon} e_{AB}(dx)$

$\Rightarrow \nu_{AB}(dx) = \frac{1}{\text{cap}(A, B)} e^{-V(x)/\varepsilon} e_{AB}(dx)$ is a probability measure on ∂A

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Theorem (“Magic” formula):

$$\mathbb{E}^{\nu_{AB}}[\tau_B] := \int_{\partial A} \mathbb{E}^x[\tau_B] \nu_{AB}(dx) = \frac{1}{\text{cap}(A, B)} \int_{B^c} e^{-V(x)/\varepsilon} h_{AB}(x) dx$$

Dirichlet principle

Theorem: Dirichlet principle

Let $\mathcal{H}_{AB} = \{h : \mathbb{R}^d \rightarrow [0, 1] : h|_A = 1, h|_B = 0\}$. Then

$$\text{cap}(A, B) = \inf_{h \in \mathcal{H}_{AB}} \mathcal{E}(h) = \mathcal{E}(h_{AB})$$

Thomson principle

Theorem: Thomson principle [Landim, Mariani, Seo 2018]

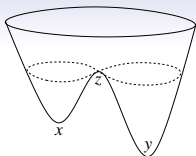
Let $\mathcal{U}_{AB} = \{f: \nabla \cdot f|_{(A \cup B)^c} = 0, \int_{\partial A} f(x) \cdot n_A(x) \sigma(dx) = 1\}$. Then

$$\text{cap}(A, B) = \sup_{f \in \mathcal{U}_{AB}} \frac{1}{\mathcal{D}(f)} = \frac{1}{\mathcal{D}(f_{AB})} \quad \mathcal{D}(f) = \frac{1}{\varepsilon} \int e^{V(x)/\varepsilon} |f(x)|^2 dx$$

Proof of Eyring–Kramers law

$$dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$$

▷ A, B small balls around x, y



Potential landscape for the particle system

$$V(x) = \sum_{i \in \mathbb{Z}/N\mathbb{Z}} U(x^i) + \frac{\gamma}{4} \sum_{i \in \mathbb{Z}/N\mathbb{Z}} (x^{i+1} - x^i)^2 \quad U(\xi) = \frac{1}{4}\xi^4 - \frac{1}{2}\xi^2$$

$$\gamma = 0: \mathcal{X} = \{-1, 0, 1\}^N, \mathcal{X}_0 = \{-1, 1\}^N, \mathcal{X}_1 = \{x \in \mathcal{X} : \text{one } x^i = 0\}$$

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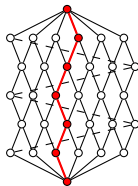
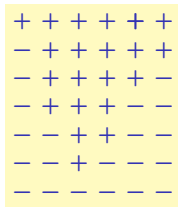
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Theorem [BFG, Nonlinearity 2007]

No bifurcation for $0 \leq \gamma \leq \gamma^*(N)$

where $\gamma^*(N) > \frac{1}{4} \quad \forall N \geq 2$

$V_\gamma(z_\gamma^*) = V_0(z_0^*) + \gamma(\# \text{ interfaces}) + \dots$
Ising-like dynamics



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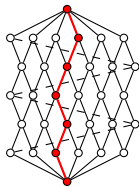
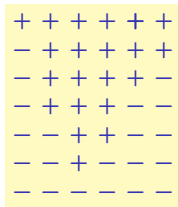
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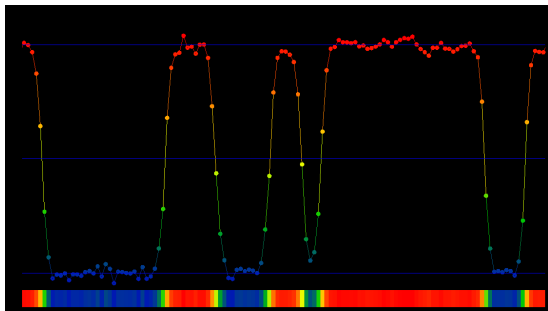
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Theorem [BFG, Nonlinearity 2007]

$\gamma > \frac{1}{2 \sin^2(\pi/N)} \Leftrightarrow \mathcal{X}_0 = \{\pm(1, \dots, 1)\}, \mathcal{X}_1 = \{0\} \Leftrightarrow \text{Synchronization}$

2. The stochastic Allen–Cahn PDE on the 1d torus



Deterministic Allen–Cahn PDE

[Chafee & Infante 74, Allen & Cahn 75]

$$\partial_t u(t, x) = \Delta u(t, x) + f(u(t, x))$$

- ▷ $x \in [0, L]$, L : bifurcation parameter
- ▷ $u(t, x) \in \mathbb{R}$
- ▷ Either periodic or zero-flux Neumann boundary conditions
- ▷ In this talk: $f(u) = u - u^3$ (results more general)

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Energy function:

$$V[u] = \int_0^L \left[\frac{1}{2} u'(x)^2 - \frac{1}{2} u(x)^2 + \frac{1}{4} u(x)^4 \right] dx \quad \rightarrow \min$$

Scaling limit of particle system with $\gamma = 2 \frac{N^2}{L^2}$

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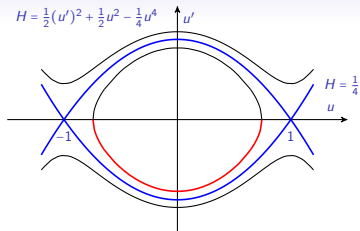
Stationary solutions: $u_0''(x) = -u_0(x) + u_0(x)^3$ critical points of V

Stability: Sturm–Liouville problem $\partial_t v_t(x) = v_t''(x) + [1 - 3u_0(x)^2]v_t(x)$

Stationary solutions

$$u_0''(x) = -f(u_0(x)) = -u_0(x) + u_0(x)^3$$

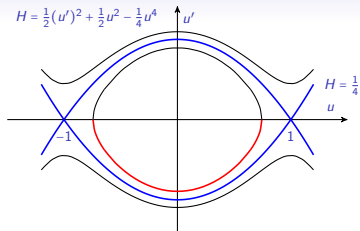
- ▷ $u_{\pm}(x) \equiv \pm 1$
- ▷ $u_0(x) \equiv 0$
- ▷ Nonconstant solutions satisfying b.c.
(expressible in terms of Jacobi elliptic fcts)



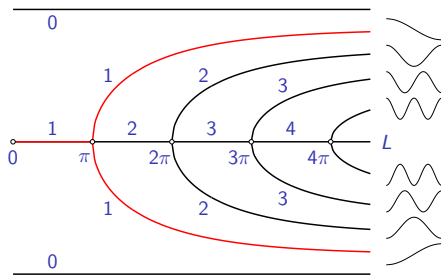
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- ▷ Neumann b.c: $2k$ nonconstant solutions when $L > k\pi$



Number of positive
eigenvalues
(= unstable directions)
Transition state



- ▷ Periodic b.c: k families when $L > 2k\pi$

Eyring–Kramers law for 1D SPDEs: heuristics

$$\partial_t u(t, x) = \Delta u(t, x) + f(u(t, x)) + \sqrt{2\varepsilon} \xi(t, x) \quad (f(u) = u - u^3)$$

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Initial condition: u_{in} near $u_- \equiv -1$ with eigenvalues $\nu_k = \left(\frac{\beta k \pi}{L}\right)^2 + 2$

Target: $u_+ \equiv 1$, $\tau_+ = \inf\{t > 0: \|u_t - u_+\|_{L^\infty} < \rho\}$

Transition state: ($\beta = 1$ for Neumann b.c., $\beta = 2$ for periodic b.c.)

$$u_{\text{ts}}(x) = \begin{cases} u_0(x) \equiv 0 & \text{if } L \leq \beta\pi \quad \text{with ev } \lambda_k = \left(\frac{\beta k \pi}{L}\right)^2 - 1 \\ u_1(x) \text{ } \beta\text{-kink stationary sol.} & \text{if } L > \beta\pi \quad \text{with ev } \lambda'_k \end{cases}$$

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[Faris & Jona-Lasinio 82]: large-deviation principle

\Rightarrow Arrhenius law: $\mathbb{E}^{u_{\text{in}}}[\tau_+] \simeq e^{(V[u_{\text{ts}}] - V[u_-])/\varepsilon}$

[Maier & Stein 01]: formal computation; for Neumann b.c.

$$\Rightarrow \mathbb{E}^{u_{\text{in}}}[\tau_+] \simeq 2\pi \sqrt{\frac{1}{|\lambda_0| \nu_0} \prod_{k=1}^{\infty} \frac{\lambda_k}{\nu_k}} e^{(V[u_{\text{ts}}] - V[u_-])/\varepsilon}$$

Eyring–Kramers law for 1D SPDEs: heuristics

Eyring–Kramers law for 1D SPDEs: main result

Theorem: Neumann b.c. [B & Gentz, 2013]

▷ If $L < \pi - c$ with $c > 0$, then

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▷ If $L > \pi + c$, then same formula with extra factor $\frac{1}{2}$ (since 2 saddles) and λ'_k instead of λ_k

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- ▶ If $L > \pi + c$, then same formula with extra factor $\frac{1}{2}$ (since 2 saddles) and λ'_k instead of λ_k

- ▶ Results also for L near π and periodic b.c.

- ▶ Prefactor involves a **Fredholm determinant**:

Δ_{\perp} Laplacian acting on mean zero functions

$$\prod_{k=1}^{\infty} \frac{\lambda_k}{\nu_k} = \det[(-\Delta_{\perp} - 1)(-\Delta_{\perp} + 2)^{-1}] = \det[\mathbf{1} - 3(-\Delta_{\perp} + 2)^{-1}]$$

converges because $\log \det = \text{Tr} \log$ and $(-\Delta_{\perp} + 2)^{-1}$ is **trace class**

(limit = $\frac{\sqrt{2} \sin(L)}{\sinh(\sqrt{2}L)}$)

Ideas of the proof ($L < \pi$)

- ▷ Spectral Galerkin approximation: $u(t, x) = \sum_{|k| \leq N} z_k(t) e_k(x)$ (Fourier)

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▷ Change of variables $u(x) = u_0 + \sqrt{\varepsilon} u_\perp(x)$ where $\int_0^L u_\perp(x) dx = 0$

$$\Rightarrow \frac{1}{\varepsilon} V[u_0 + \sqrt{\varepsilon} u_\perp] = \frac{1}{\varepsilon} \underbrace{\left(\frac{1}{4} u_0^4 - \frac{1}{2} u_0^2 \right)}_{V_0(u_0)} + Q_{u_0}[u_\perp] + \sqrt{\varepsilon} R_{u_0}[u_\perp]$$

where

$$\begin{cases} Q_{u_0}[u_\perp] = \frac{1}{2} \int_0^L [u_\perp'(x)^2 - (1 - 3u_0^2) u_\perp(x)^2] dx = \frac{1}{2} \langle u_\perp, [-\Delta - (1 - 3u_0^2)] u_\perp \rangle \\ R_{u_0}[u_\perp] = u_0 \int_0^L u_\perp(x)^3 dx + \sqrt{\varepsilon} \int_0^L u_\perp(x)^4 dx \quad (\text{remainder}) \end{cases}$$

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$$\begin{cases} Q_{u_0}[u_\perp] = \frac{1}{2} \int_0^L [u_\perp'(x)^2 - (1 - 3u_0^2) u_\perp(x)^2] dx = \frac{1}{2} \langle u_\perp, [-\Delta - (1 - 3u_0^2)] u_\perp \rangle \\ R_{u_0}[u_\perp] = u_0 \int_0^L u_\perp(x)^3 dx + \sqrt{\varepsilon} \int_0^L u_\perp(x)^4 dx \quad (\text{remainder}) \end{cases}$$

▷ Dirichlet principle with $h = h(u_0)$ s.t. $h'(u_0) = -\frac{1}{c} e^{V_0(u_0)/\varepsilon}$, $c \simeq \sqrt{\frac{2\pi\varepsilon}{|\lambda_0|}}$

$$\begin{aligned} \text{cap}(A, B) \leq \mathcal{E}(h) &= \frac{\varepsilon^{1+\frac{N}{2}}}{c^2} \int_{-1}^1 e^{V_0(u_0)/\varepsilon} \underbrace{\int e^{-Q_{u_0}[u_\perp]} e^{-\sqrt{\varepsilon} R_{u_0}[u_\perp]} du_\perp}_{=} du_0 \\ &= \sqrt{\frac{(2\pi)^N}{\det[-\Delta_\perp - (1 - 3u_0^2)]}} \mathbb{E} \gamma [e^{-\sqrt{\varepsilon} R_{u_0}}] \end{aligned}$$

Ideas of the proof ($L < \pi$)

- ▷ Thomson principle with divergence-free unit flow $f = K^{-1} e^{-Q_0[u_\perp]} e_{u_0}$

Normalisation
$$K = \varepsilon^{\frac{N}{2}} \int e^{-Q_0[u_\perp]} du_\perp = \sqrt{\frac{(2\pi\varepsilon)^N}{\det[-\Delta_\perp - 1]}}$$

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- ▷ Conclusion: $\text{cap}(A, B) = \varepsilon \sqrt{\frac{|\lambda_0|}{2\pi\varepsilon}} \sqrt{\frac{(2\pi\varepsilon)^N}{\det[-\Delta_\perp - 1]}} [1 + \mathcal{O}(\varepsilon)]$

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Other elements of the proof:

- ▶ A priori bounds on h_{AB} : large deviations (or symmetry argument)
- ▶ Convergence of hitting times as $N \rightarrow \infty$: a priori estimate for $\mathbb{E}[\tau_B^2]$
- ▶ Coupling argument for start in u_{in} [Martinelli, Olivieri & Scoppola]
- ▶ Bifurcation at $L = \beta\pi$

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Thanks for your attention!

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