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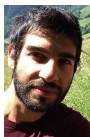
Metastable dynamics of stochastic Allen-Cahn PDEs on the torus

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March 26, 2021
(Online)

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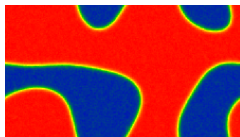
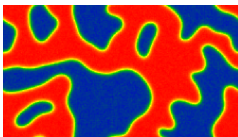
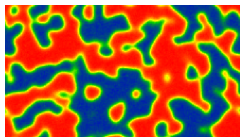
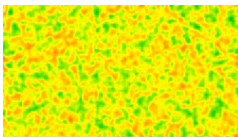
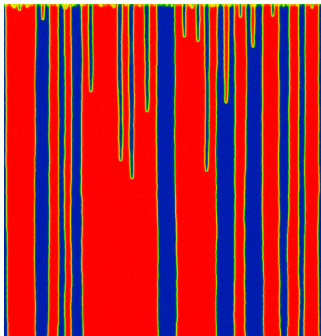
Deterministic Allen–Cahn PDE

[Chafee & Infante 74, Allen & Cahn 75]

$$\partial_t u = \Delta u + u - u^3$$

▷ $u = u(t, x) \in \mathbb{R}$, $t \geq 0$, $x \in \mathbb{T}_L^d = (\mathbb{R}/L\mathbb{Z})^d$, $L > 0$

Phase separation ($d = 1$: [Carr & Pego 89, Chen 04])



Remark: ϕ^4 model from QFT: $\partial_t u = \Delta u - u^3$

Deterministic Allen–Cahn PDE

$$\partial_t u = \Delta u + u - u^3$$

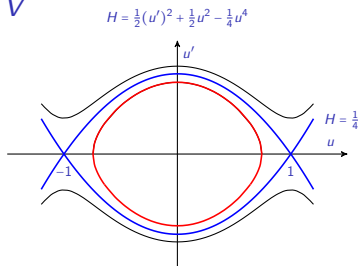
Energy function:

$$V[u] = \int_{\mathbb{T}_L^d} \left[\frac{1}{2} |\nabla u(x)|^2 - \frac{1}{2} u(x)^2 + \frac{1}{4} u(x)^4 \right] dx \quad \Rightarrow \quad \nabla_v V[u] = -\langle \partial_t u, v \rangle$$

Stationary solutions ($d = 1$):

$$u_0''(x) = -u_0(x) + u_0(x)^3 \quad \text{critical points of } V$$

- ▷ $u_{\pm}(x) \equiv \pm 1$
- ▷ $u_0(x) \equiv 0$
- ▷ Nonconstant solutions if $L > 2\pi$
(expressible in terms of Jacobi elliptic fcts)



Stability of stationary solutions

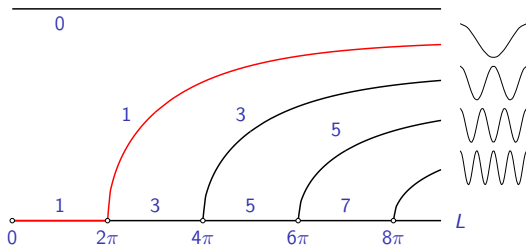
Linearisation at u_0 :

$$\partial_t v_t(x) = (\mathcal{L}v)(x) := v_t''(x) + [1 - 3u_0(x)^2]v_t(x)$$

Stability: Find eigenvalues of \mathcal{L} \longrightarrow Sturm–Liouville problem

- ▷ $u_{\pm}(x) \equiv \pm 1$: stable
- ▷ $u_0(x) \equiv 0$: unstable
- ▷ Nonconstant solutions: unstable

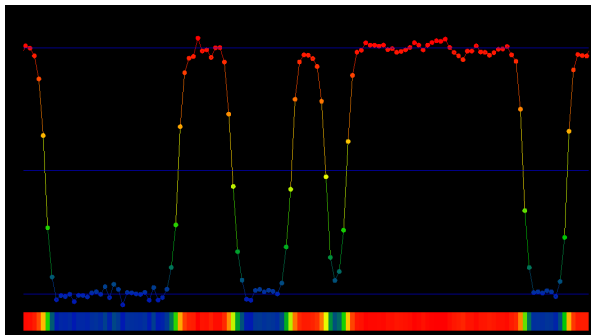
Number of positive
eigenvalues
(= unstable directions)
Transition state



Stochastic Allen–Cahn PDE

$$\partial_t u = \Delta u + u - u^3 + \sqrt{2\varepsilon}\xi$$

- ▷ ξ space-time white noise: centered, Gaussian, $\mathbb{E}[\xi(t, x)\xi(s, y)] = \delta(t - s)\delta(x - y)$
- ▷ ξ distribution, $\langle \xi, \varphi \rangle \sim \mathcal{N}(0, \|\varphi\|_{L^2}^2)$, $\mathbb{E}[\langle \xi, \varphi_1 \rangle \langle \xi, \varphi_2 \rangle] = \langle \varphi_1, \varphi_2 \rangle_{L^2}$



Question: Speed of convergence to equilibrium measure?

Reversible diffusion in a double-well

$$dx_t = -\nabla V(x_t) dt + \sqrt{2\varepsilon} dW_t$$

$V : \mathbb{R}^d \rightarrow \mathbb{R}$ confining 2-well potential

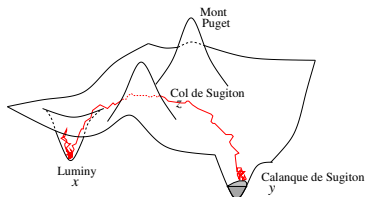
▷ **Generator:**

$$\mathcal{L} = \varepsilon \Delta - \nabla V \cdot \nabla = \varepsilon e^{V/\varepsilon} \nabla \cdot e^{-V/\varepsilon} \nabla$$

▷ **Invariant probability: Gibbs measure**

$$\pi(dx) = \frac{1}{Z} e^{-V(x)/\varepsilon} dx \quad \Rightarrow \quad \mathcal{L}^\dagger \pi = 0$$

▷ **Reversible:** $\langle f, \mathcal{L}g \rangle_\pi = \langle \mathcal{L}f, g \rangle_\pi$ for $\langle f, g \rangle_\pi = \int_{\mathbb{R}^d} f(x)g(x)\pi(dx)$



How to characterise convergence to equilibrium?

▷ **Exponential ergodicity:** $|\mathbb{E}^x[f(x_t)] - \langle \pi, f \rangle| \leq C(x, f) e^{-\beta t}$

Possible via **Lyapunov functions** [Meyn & Tweedie], bad control of β

▷ **Mean hitting time:** Estimate $\mathbb{E}^x[\tau_y]$ where $\tau_y = \inf\{t > 0 : x_t \in \mathcal{B}_\varepsilon(y)\}$

▷ **Spectral gap** of \mathcal{L}

Mean first hitting time

- ▷ Arrhenius' law (1889): $\mathbb{E}^x[\tau_y] \simeq e^{[V(z)-V(x)]/\varepsilon}$

Proved by [Freidlin, Wentzell, 1979] using large deviations

- ▷ Eyring–Kramers law (1935, 1940):

Eigenvalues of Hessian of V at minimum x : $0 < \nu_1 \leq \nu_2 \leq \dots \leq \nu_d$

Eigenvalues of Hessian of V at saddle z : $\lambda_1 < 0 < \lambda_2 \leq \dots \leq \lambda_d$

$$\mathbb{E}^x[\tau_y] = 2\pi \sqrt{\frac{\lambda_2 \dots \lambda_d}{|\lambda_1| \nu_1 \dots \nu_d}} e^{[V(z)-V(x)]/\varepsilon} [1 + \mathcal{O}_\varepsilon(1)]$$

Spectral gap of $\mathcal{L} = \mathbb{E}^x[\tau_y]^{-1} [1 + \mathcal{O}_\varepsilon(1)]$

Proved by [Bovier, Eckhoff, Gayraud, Klein, 2004] using potential theory,
by [Helffer, Klein, Nier, 2004] using Witten Laplacian, ...

Potential-theoretic proof of Eyring–Kramers law

$$\triangleright w_A(x) = \mathbb{E}^x[\tau_A] \quad \text{satisfies} \quad \begin{cases} (\mathcal{L} w_A)(x) = -1 & x \in A^c \\ w_A(x) = 0 & x \in A \end{cases}$$

$$\triangleright h_{AB}(x) = \mathbb{P}^x\{\tau_A < \tau_B\} \quad \text{satisfies} \quad \begin{cases} (\mathcal{L} h_{AB})(x) = 0 & x \in (A \cup B)^c \\ h_{AB}(x) = 1 & x \in A \\ h_{AB}(x) = 0 & x \in B \end{cases}$$

Theorem: $A, B \subset \mathbb{R}^d$ disjoint. \exists proba measure ν_{AB} on ∂A s.t.

$$\int_{\partial A} \mathbb{E}^x[\tau_B] \nu_{AB}(dx) = \frac{1}{\text{cap}(A, B)} \int_{B^c} e^{-V(y)/\varepsilon} h_{AB}(y) dy$$

Proof: $w_B(x) = - \int_{B^c} G_{B^c}(x, y) dy \quad h_{AB}(y) = - \int_{\partial A} G_{B^c}(y, x) e_{AB}(dx)$

$$\nu_{AB}(dx) := \frac{1}{\text{cap}(A, B)} e^{-V(x)/\varepsilon} e_{AB}(dx) \quad \text{cap}(A, B) := \int_{\partial A} e^{-V/\varepsilon} e_{AB}(dx)$$

$$\text{cap}(A, B) \int_{\partial A} w_B(x) \nu_{AB}(dx) = - \int_{\partial A} \int_{B^c} \underbrace{G_{B^c}(x, y) e^{-V(x)/\varepsilon}}_{= G_{B^c}(y, x) e^{-V(y)/\varepsilon}} e_{AB}(dx) dy$$

Potential-theoretic proof of Eyring–Kramers law

Theorem: $A, B \subset \mathbb{R}^d$ disjoint. \exists proba measure ν_{AB} on ∂A s.t.

$$\int_{\partial A} \mathbb{E}^x[\tau_B] \nu_{AB}(dx) = \frac{1}{\text{cap}(A, B)} \int_{B^c} e^{-V(y)/\varepsilon} h_{AB}(y) dy$$

Apply to A, B neighbourhoods of x, y

- ▶ Laplace asymptotics: $\int_{B^c} h_{A,B}(y) e^{-V(y)/\varepsilon} dy \simeq \sqrt{\frac{(2\pi\varepsilon)^d}{\nu_1 \dots \nu_d}} e^{-V(x)/\varepsilon}$
- ▶ Capacity: $\text{cap}(A, B) = \mathcal{E}(h_{AB})$
- ▶ Dirichlet form: $\mathcal{E}(f) = \langle f, -\mathcal{L}f \rangle = \varepsilon \int_{\mathbb{R}^d} e^{-V(x)/\varepsilon} |\nabla f(x)|^2 dx$

Theorem: Dirichlet principle

Let $\mathcal{H}_{AB} = \{h : \mathbb{R}^d \rightarrow [0, 1] : h|_A = 1, h|_B = 0\}$. Then

$$\text{cap}(A, B) = \inf_{h \in \mathcal{H}_{AB}} \mathcal{E}(h) = \mathcal{E}(h_{AB})$$

Appropriate test function yields $\text{cap}(A, B) \simeq \varepsilon \sqrt{\frac{|\lambda_1|}{2\pi\varepsilon}} \sqrt{\frac{(2\pi\varepsilon)^{d-1}}{\lambda_2 \dots \lambda_d}} e^{-V(z)/\varepsilon}$

Stochastic Allen–Cahn PDE

$$\partial_t u = \Delta u + u - u^3 + \sqrt{2\varepsilon}\xi \quad x \in \mathbb{T}_L$$

- ▷ [Faris & Jona-Lasinio '82]: existence/uniqueness of solution
- ▷ [Da Prato & Zabczyk '90s]: invariant measure is **Gibbs measure** associated with

$$V[u] = \int_{\mathbb{T}_L} \left[\frac{1}{2} |\nabla u(x)|^2 - \frac{1}{2} u(x)^2 + \frac{1}{4} u(x)^4 \right] dx$$

- ▷ [Faris & Jona-Lasinio '82]: **Arrhenius law** $\mathbb{E}^{u_-} [\tau_{u_+}] \simeq e^{(V[u_{\text{tr}}] - V[u_-])/\varepsilon}$
where u_{tr} transition state, $u_{\text{tr}} = u_0$ if $L < 2\pi$

Question: is there an **Eyring–Kramers** law? With what prefactor?

Heuristics: $V[u] = \frac{1}{2} \langle u, (-\Delta - 1)u \rangle + \mathcal{O}(u^4)$

- ▷ Hessian of V at u_0 : $-\Delta - 1$, eigenvalues $\lambda_k = \left(\frac{2k\pi}{L}\right)^2 - 1$
- ▷ Hessian of V at u_- : $-\Delta + 2$, eigenvalues $\nu_k = \left(\frac{2k\pi}{L}\right)^2 + 2$

Formally, product of ratios of λ_k/ν_k converges [Maier & Stein '01]

Formal computation and Fredholm determinant

Formally (for $L < 2\pi$)

$$\mathbb{E}^{u_-}[\tau_{u_+}] = \frac{2\pi}{|\lambda_1|} \sqrt{\frac{|\det \text{Hess } V[u_0]|}{\det \text{Hess } V[u_-]}} e^{(V[u_0] - V[u_-])/\varepsilon} [1 + \mathcal{O}_\varepsilon(1)]$$

Δ_\perp Laplacian acting on mean zero functions

$$\begin{aligned} \det([- \Delta_\perp - 1][- \Delta_\perp + 2]^{-1}) &= \det([- \Delta_\perp + 2 - 3][- \Delta_\perp + 2]^{-1}) \\ &= \det(\mathbb{1} - 3[- \Delta_\perp + 2]^{-1}) \end{aligned}$$

Fredholm determinant

$$\begin{aligned} \log \det(\mathbb{1} - 3[- \Delta_\perp + 2]^{-1}) &= \text{Tr} \log(\mathbb{1} - 3[- \Delta_\perp + 2]^{-1}) \\ &= - \sum_{n \geq 1} \frac{3^n}{n} \underbrace{\text{Tr}([- \Delta_\perp + 2]^{-n})}_{\lesssim \left[\left(\frac{2\pi}{L} \right)^2 + 2 \right]^{-n}} < \infty \quad (L < 2\pi) \end{aligned}$$

General fact: $\det(\mathbb{1} + T) < \infty$ if T is trace class

Main result for $d = 1$

Theorem: [B & Gentz, Elec. J. Proba 2013]

- ▶ If $L < 2\pi - c$ with $c > 0$, then

$$\mathbb{E}^{u_-}[\tau_+] = 2\pi \sqrt{\det(\mathbb{1} - 3[-\Delta + 2]^{-1})} e^{(V[u_0] - V[u_-])/\varepsilon} [1 + \mathcal{O}_\varepsilon(1)]$$

- ▶ Similar explicit expressions for $L > 2\pi - c$ and $L \simeq 2\pi$
(with different u_{tr} and slightly different ε -dependence due to 0 eigenvalue)

Remarks:

- ▶ Proof relies on **spectral Galerkin approximation**
- ▶ Error more precise than $\mathcal{O}_\varepsilon(1)$
- ▶ If $u_{tr} \neq u_0$, Fredholm determinant computed with techniques from **path integrals** [Maier & Stein]
- ▶ Similar results for **Neumann b.c.**
- ▶ Similar results for other nonlinearities than $-u^3$

Allen–Cahn SPDE for $d = 2$

- ▷ Arrhenius law holds via large-deviation principle [Hairer & Weber, 2015]
- ▷ Naive computation of prefactor fails:

$$\begin{aligned}\mathrm{Tr} \log(\mathbb{1} - 3[-\Delta_{\perp} + 2]^{-1}) &\simeq \sum_{k \in (\mathbb{Z}^2)^*} \log\left(1 - \frac{3L^2}{|k|^2 \pi^2}\right) \\ &\simeq - \sum_{k \in (\mathbb{Z}^2)^*} \frac{3L^2}{|k|^2 \pi^2} \simeq -\frac{3L^2}{\pi^2} \int_1^{\infty} \frac{r \, dr}{r^2} = -\infty\end{aligned}$$

- ▷ In fact, the equation needs to be **renormalised**

Theorem: [Da Prato & Debussche 2003]

Let ξ^{δ} be a mollification on scale δ of white noise. Then

$$\partial_t u = \Delta u + [1 + 3\varepsilon C(\delta)]u - u^3 + \sqrt{2\varepsilon} \xi^{\delta}$$

with $C(\delta) \simeq \log(\delta^{-1})$ admits local solution converging as $\delta \rightarrow 0$

- ▷ $C(\delta) \sim$ variance of mollified **Gaussian free field (GFF)**
- ▷ Naively, one could expect $u_{\pm} = \pm \sqrt{1 + 3\varepsilon C(\delta)}$ but this is **not** the case

Main result in dimension 2

- ▶ Use **spectral Galerkin approximation** with cut-off N instead of mollification, $L^2 C_N = \text{Tr}(P_N[-\Delta + 2]^{-1}) \sim \log(N)$
- ▶ $V_N[u_0] - V_N[u_-] = \frac{L^2}{4} + \frac{3}{2}L^2 \varepsilon C_N$
- ▶ (Prefactor) $^2 = \det(\mathbb{1} - 3P_N[-\Delta_\perp + 2]^{-1}) e^{3\text{Tr}(P_N[-\Delta_\perp + 2]^{-1})}$
 $\det_2(\mathbb{1} - T) = \det(\mathbb{1} - T) e^{\text{Tr } T}$ **Carleman–Fredholm determinant**

Theorem: [B, Di Gesù, Weber, Elec. J. Proba 2017]

For $L < 2\pi$, appropriate $A \ni u_-$, $B \ni u_+$, $\exists \mu_N$ probability measures on ∂A :

$$\limsup_{N \rightarrow \infty} \mathbb{E}^{\mu_N}[\tau_B] \leq 2\pi \sqrt{\det_2(\mathbb{1} - 3[-\Delta + 2]^{-1})} e^{(V[u_0] - V[u_-])/\varepsilon} [1 + c_+ \sqrt{\varepsilon}]$$

$$\liminf_{N \rightarrow \infty} \mathbb{E}^{\mu_N}[\tau_B] \geq 2\pi \sqrt{\det_2(\mathbb{1} - 3[-\Delta + 2]^{-1})} e^{(V[u_0] - V[u_-])/\varepsilon} [1 - c_- \varepsilon]$$

- ▶ [Tsatsoulis & Weber, PTRF 2018]: Same result for $\mathbb{E}^{u_-}[\tau_B]$
- ▶ \det_2 defined whenever T is only **Hilbert–Schmidt** (true for $d \leq 3$)

References

- ▷ N. B. & Barbara Gentz, *Sharp estimates for metastable lifetimes in parabolic SPDEs: Kramers' law and beyond*, *Electronic J. Probability* **18**, (24):1–58 (2013)
- ▷ N. B., Giacomo Di Gesù & Hendrik Weber, *An Eyring–Kramers law for the stochastic Allen–Cahn equation in dimension two*, *Electronic J. Probability* **22**, 1–27 (2017)
- ▷ N. B., *An introduction to singular stochastic PDEs: Allen–Cahn equations, metastability and regularity structures*, Lecture notes, Sarajevo Stochastic Analysis Winter School, January 2019 [hal-02004985](#)
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Thanks for your attention!