Does noise create or suppress mixed-mode oscillations?

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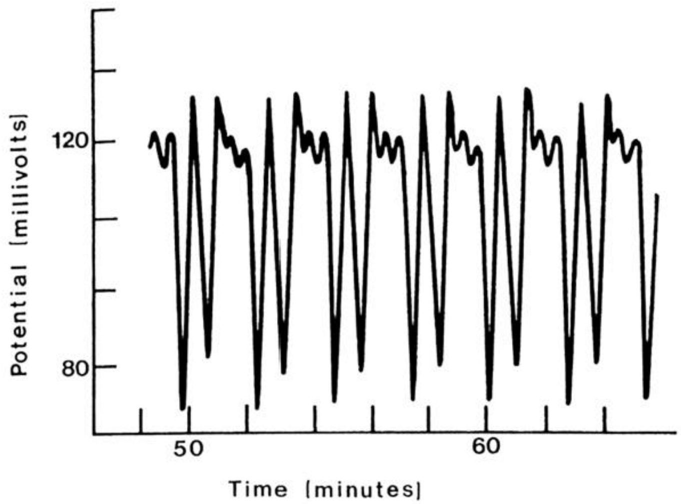
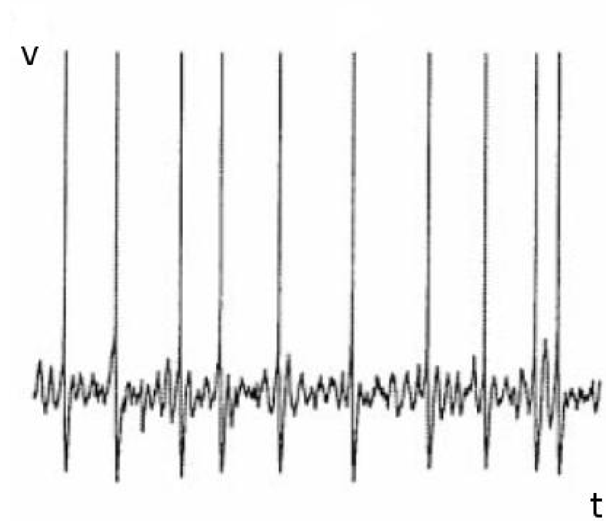
ANR project MANDy, Mathematical Analysis of Neuronal Dynamics

Stochastic Dynamics in Mathematics, Physics and Engineering
ZiF, Bielefeld, November 2, 2011

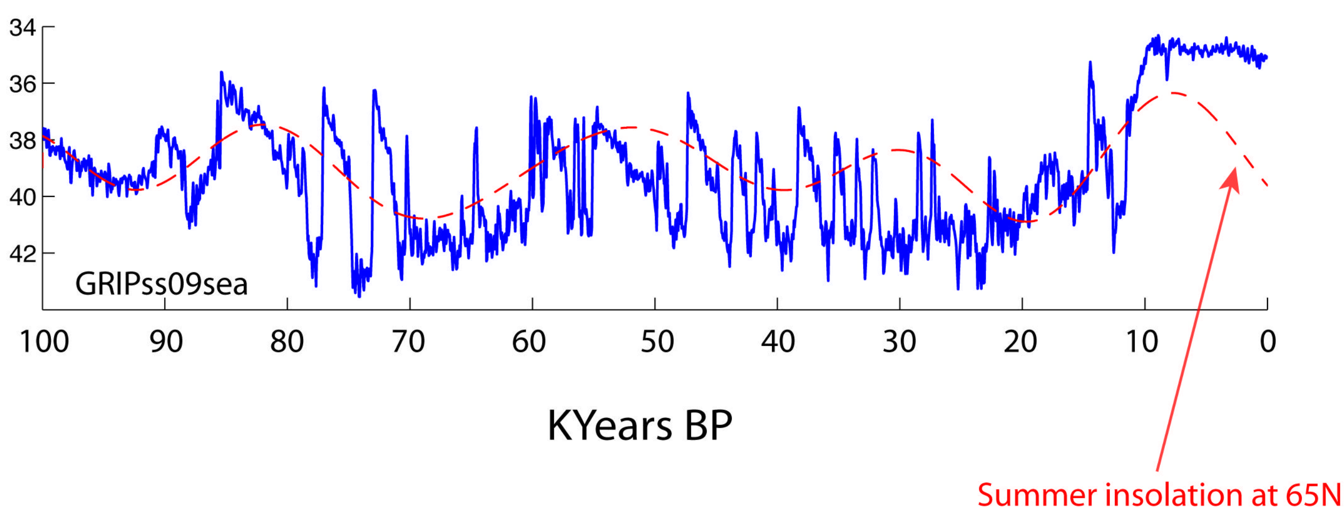
Mixed-mode oscillations (MMOs)



Belousov-Zhabotinsky reaction [Hudson 79]

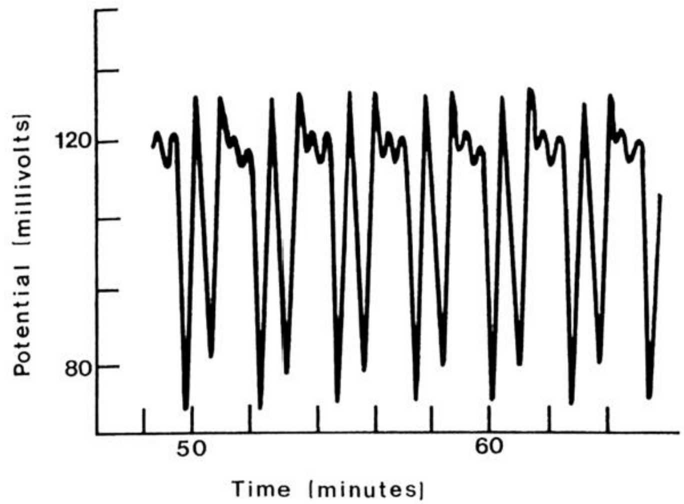


Stellate cells [Dickson 00]



Mean temperature based on ice core measurements [Johnson et al 01]

Mixed-mode oscillations (MMOs)



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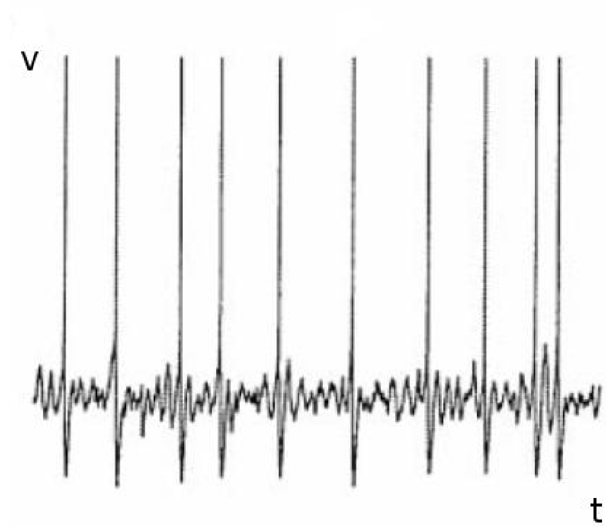

Stellate cells [Dickson 00]

▷ Deterministic models reproducing these oscillations exist and have been abundantly studied

They often involve singular perturbation theory

▷ We want to understand the effect of noise on oscillatory patterns

Noise may also induce oscillations not present in deterministic case

Part I

Where noise creates MMOs

Deterministic FitzHugh–Nagumo (FHN) equations

Consider the FHN equations in the form

$$\varepsilon \dot{x} = x - x^3 + y$$
$$\dot{y} = a - x$$

- ▷ x $\propto$ membrane potential of neuron
- ▷ y $\propto$ proportion of open ion channels (recovery variable)
- ▷ $\varepsilon \ll 1$ $\Rightarrow$ fast–slow system

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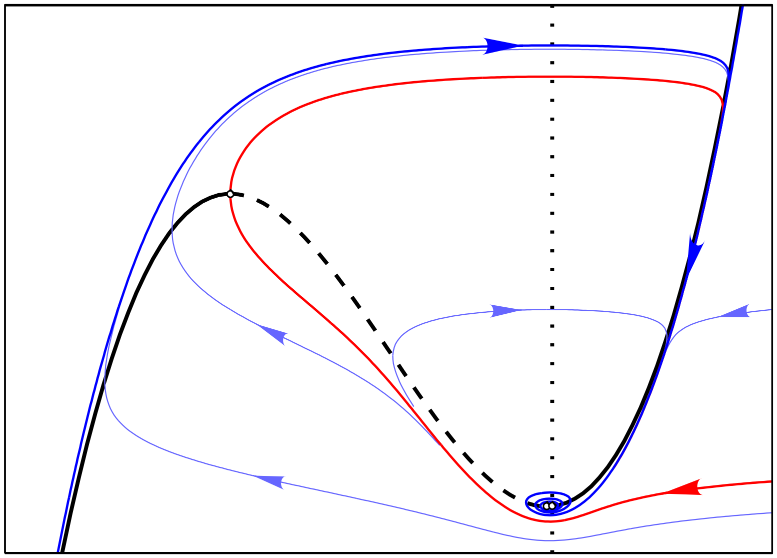
Stationary point $P = (a, a^3 - a)$
Linearisation has eigenvalues $\frac{-\delta \pm \sqrt{\delta^2 - \varepsilon}}{\varepsilon}$ where $\delta = \frac{3a^2-1}{2}$

- ▷ $\delta > 0$: stable node ($\delta > \sqrt{\varepsilon}$) or focus ($0 < \delta < \sqrt{\varepsilon}$)
- ▷ $\delta = 0$: singular Hopf bifurcation [Erneux & Mandel '86]
- ▷ $\delta < 0$: unstable focus ($-\sqrt{\varepsilon} < \delta < 0$) or node ($\delta < -\sqrt{\varepsilon}$)

Deterministic FitzHugh–Nagumo (FHN) equations

$\delta > 0$:

- ▷ P is asymptotically stable
- ▷ the system is excitable
- ▷ one can define a separatrix



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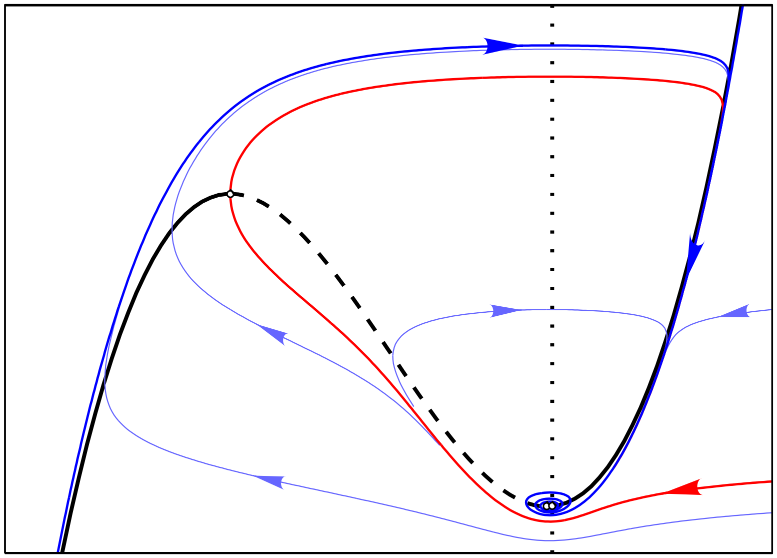
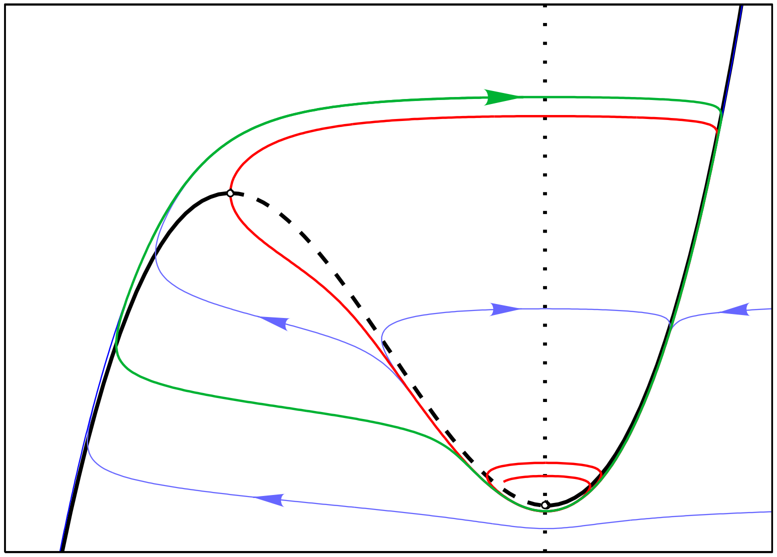


$\delta < 0$:

- ▷ P is unstable
- ▷ $\exists$ asympt. stable periodic orbit
- ▷ sensitive dependence on δ: canard (duck) phenomenon [Callot, Diener, Diener '78, Benoît '81, ...]



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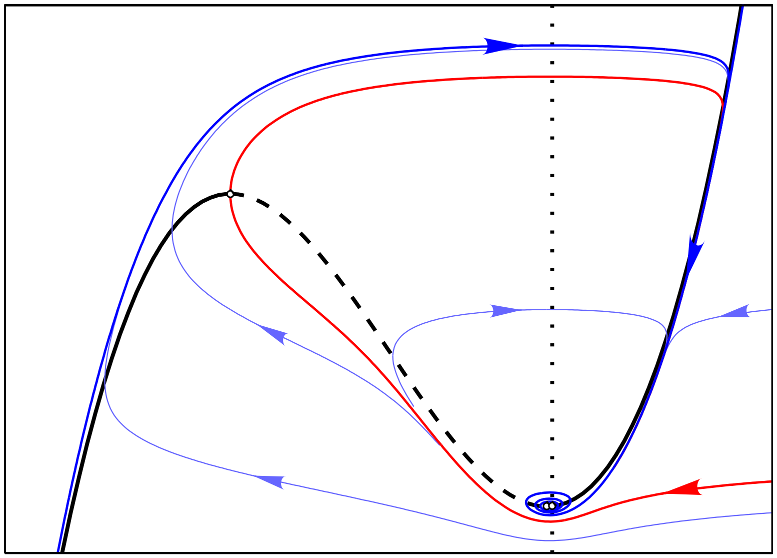

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Stochastic FHN equations

$$\mathrm{d}x_t = \frac{1}{\varepsilon}[x_t - x_t^3 + y_t]\,\mathrm{d}t + \frac{\sigma_1}{\sqrt{\varepsilon}}\,\mathrm{d}W_t^{(1)}$$
$$\mathrm{d}y_t = [a - x_t]\,\mathrm{d}t + \sigma_2\,\mathrm{d}W_t^{(2)}$$

▷ $W_t^{(1)}, W_t^{(2)}$: independent Wiener processes

▷ $0 < \sigma_1, \sigma_2 \ll 1$, $\sigma = \sqrt{\sigma_1^2 + \sigma_2^2}$

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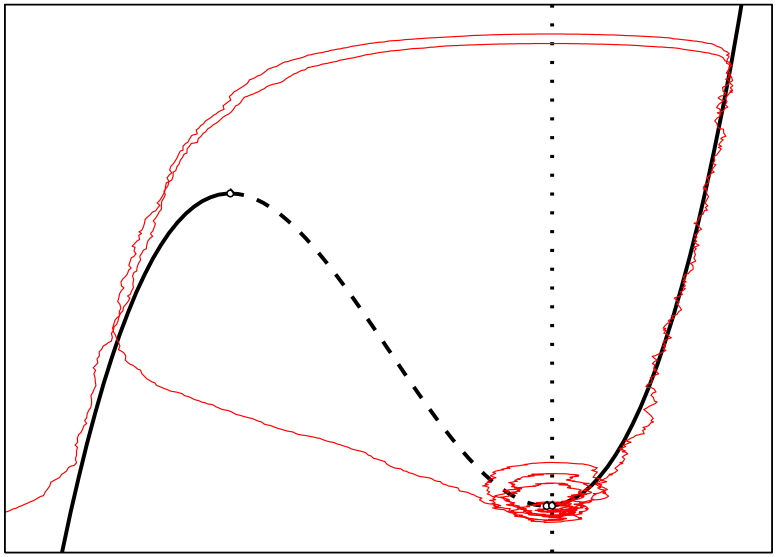
$\varepsilon = 0.1$
$\delta = 0.02$
$\sigma_1 = \sigma_2 = 0.03$



Some previous work

- ▷ Numerical: Kosmidis & Pakdaman '03, . . . , Borowski et al '11
- ▷ Moment methods: Tanabe & Pakdaman '01
- ▷ Approx. of Fokker–Planck equ: Lindner et al '99, Simpson & Kuske '11
- ▷ Large deviations: Muratov & Vanden Eijnden '05, Doss & Thieullen '09
- ▷ Sample paths near canards: Sowers '08

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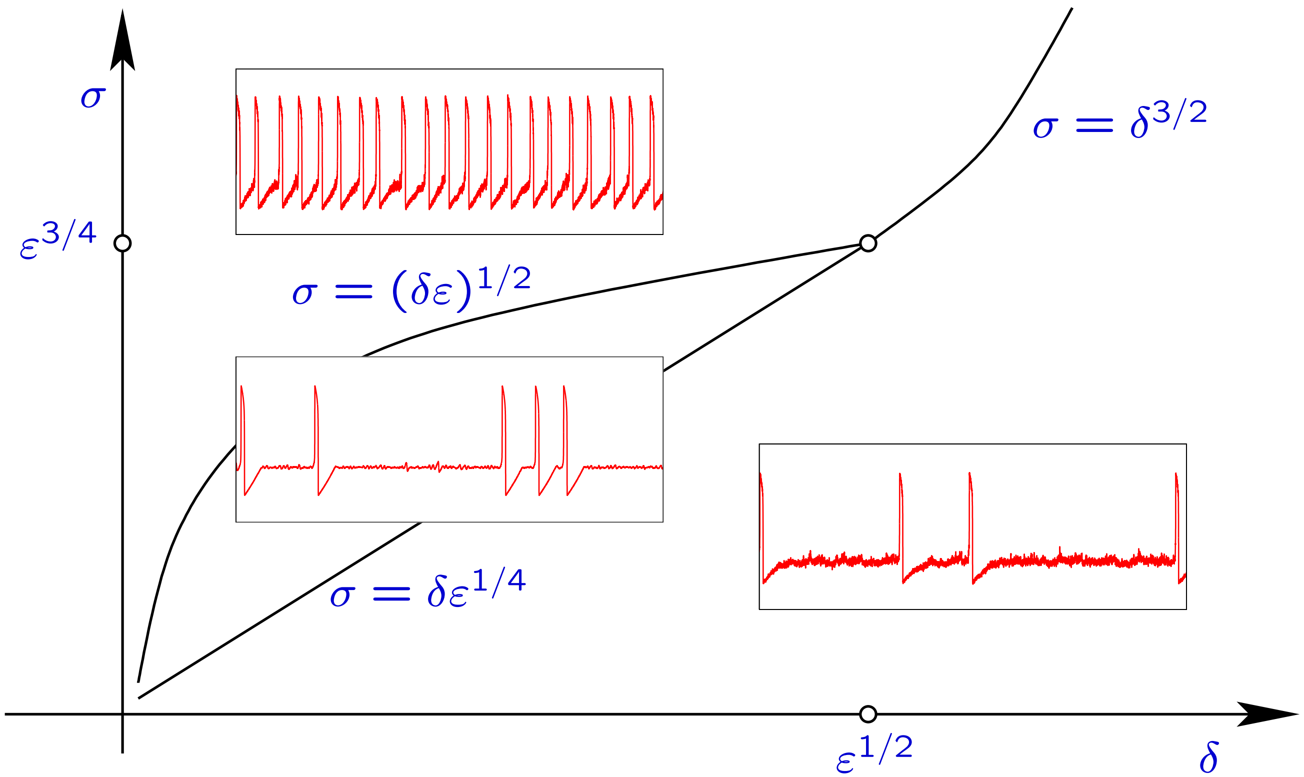
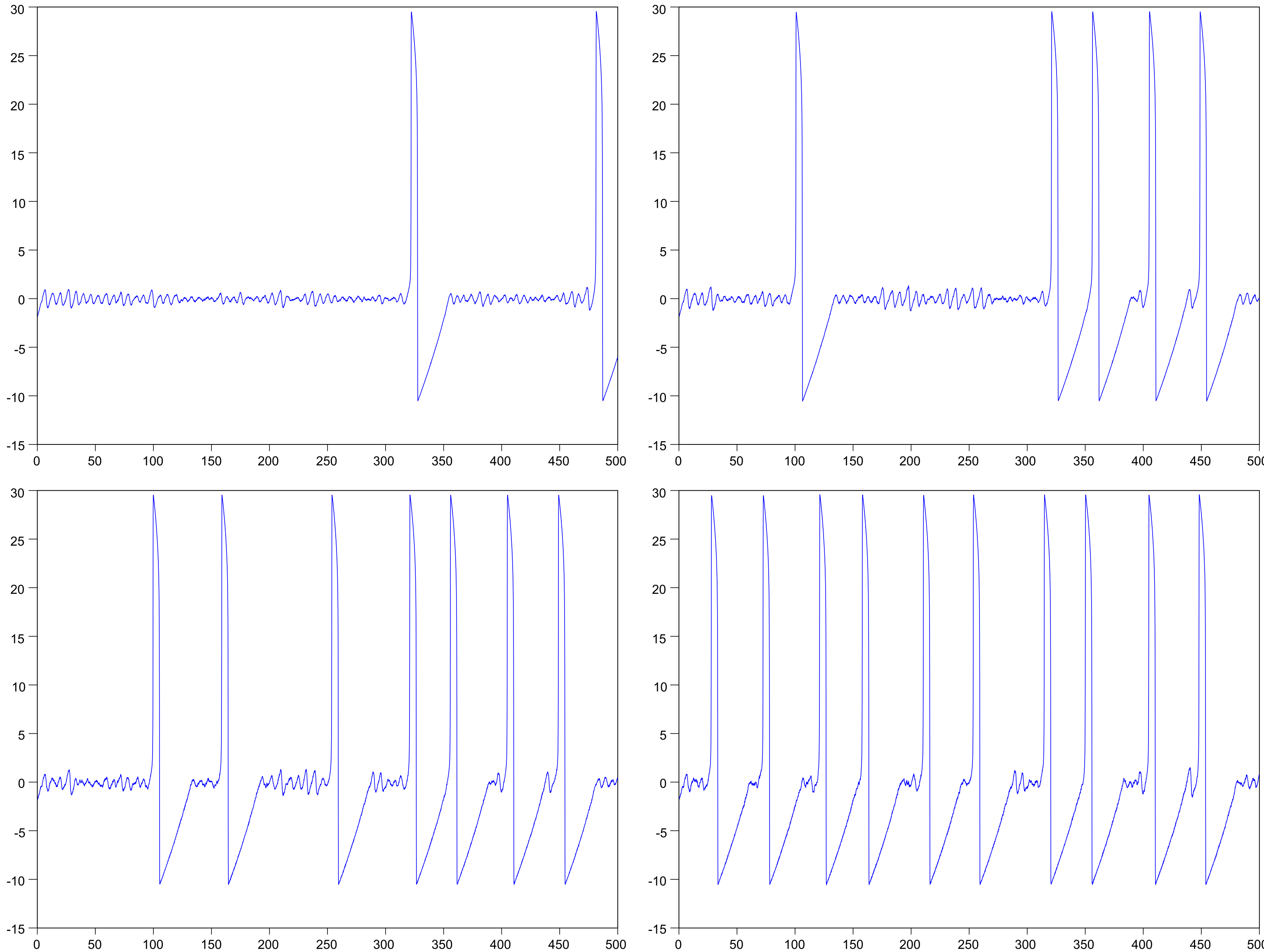
Proposed "phase diagram" [Muratov & Vanden Eijnden '08]

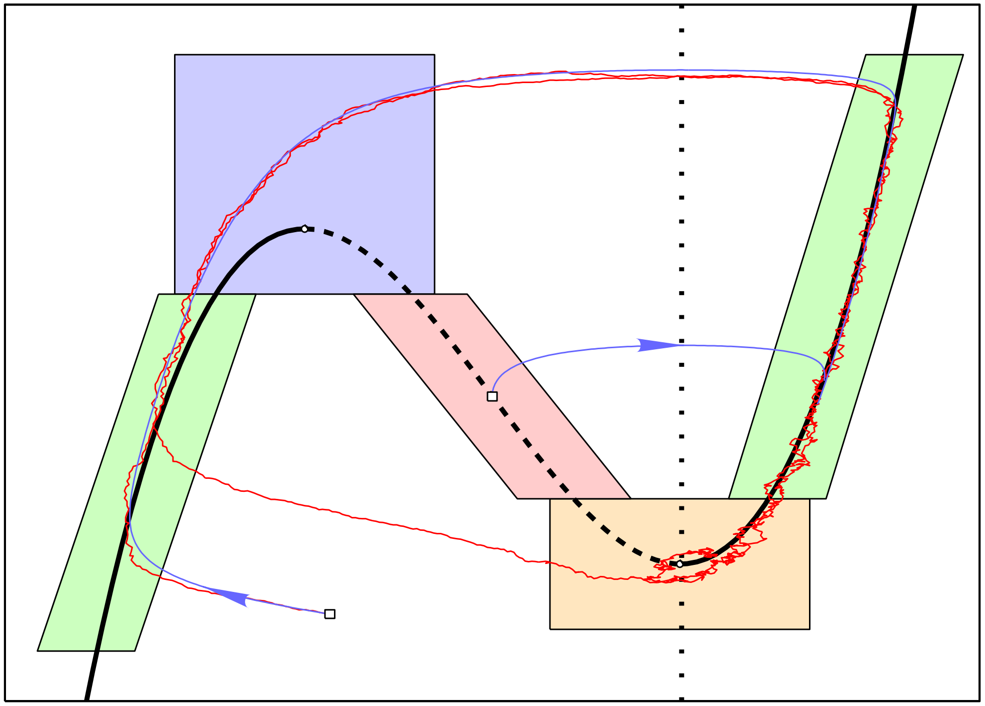


Intermediate regime: mixed-mode oscillations (MMOs)



Time series $t \mapsto -x_t$ for $\varepsilon = 0.01$, $\delta = 3 \cdot 10^{-3}$, $\sigma = 1.46 \cdot 10^{-4}, \dots, 3.65 \cdot 10^{-4}$

Precise analysis of sample paths



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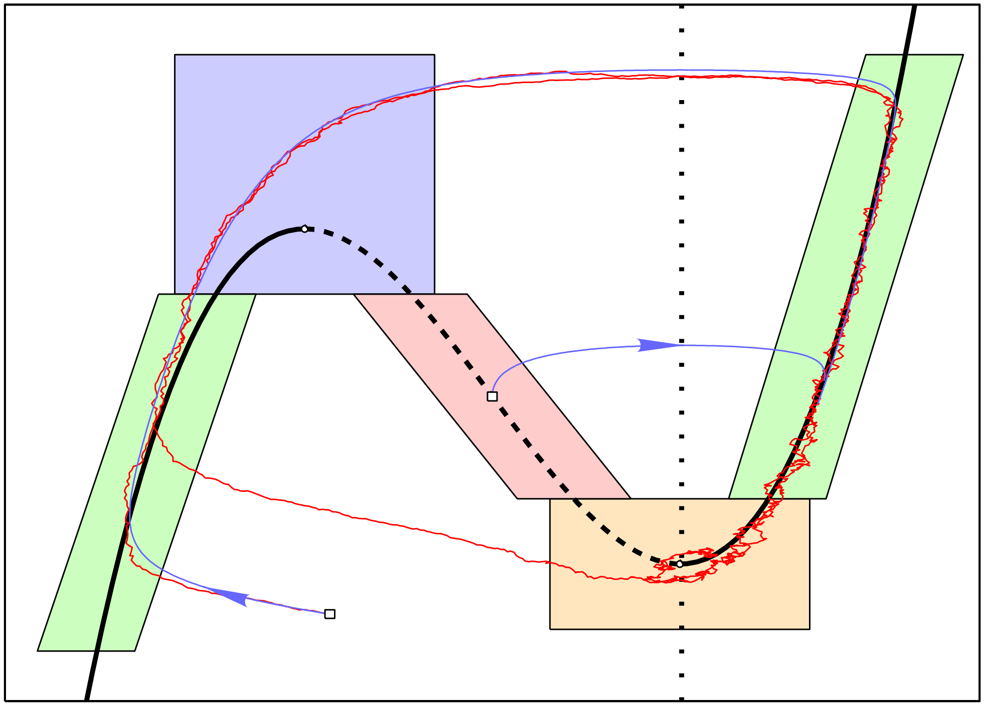
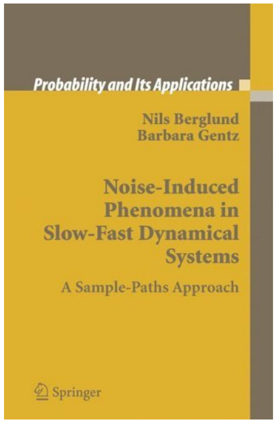


▷ Dynamics near stable branch, unstable branch and saddle–node bifurcation: already done in [B & Gentz '05]



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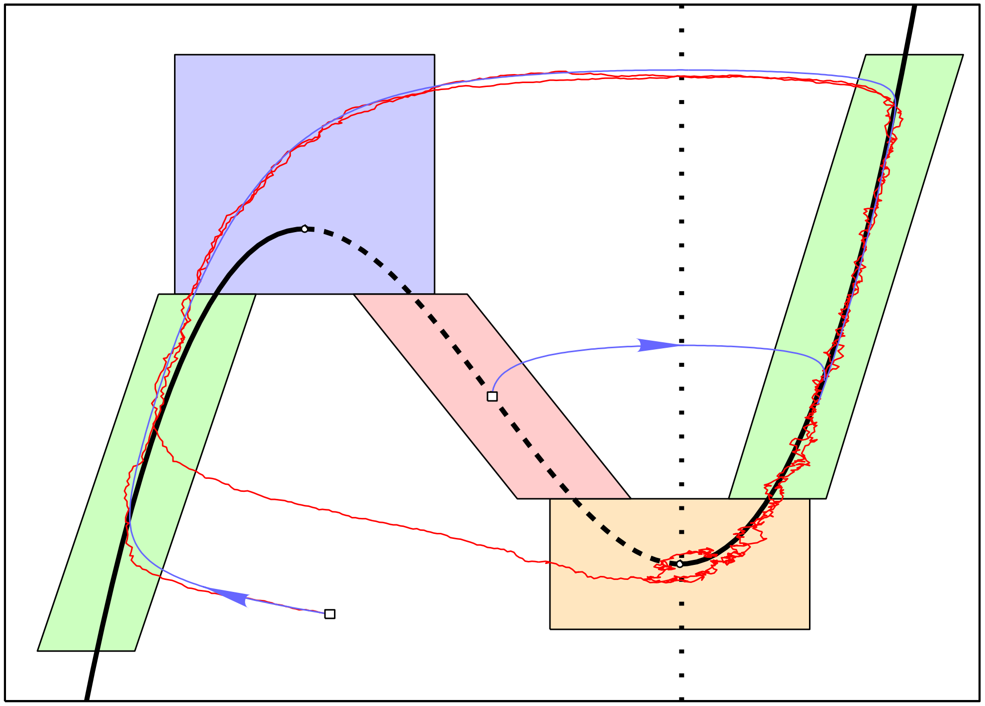

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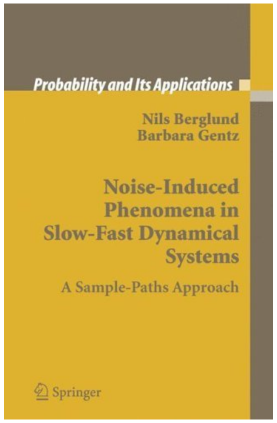
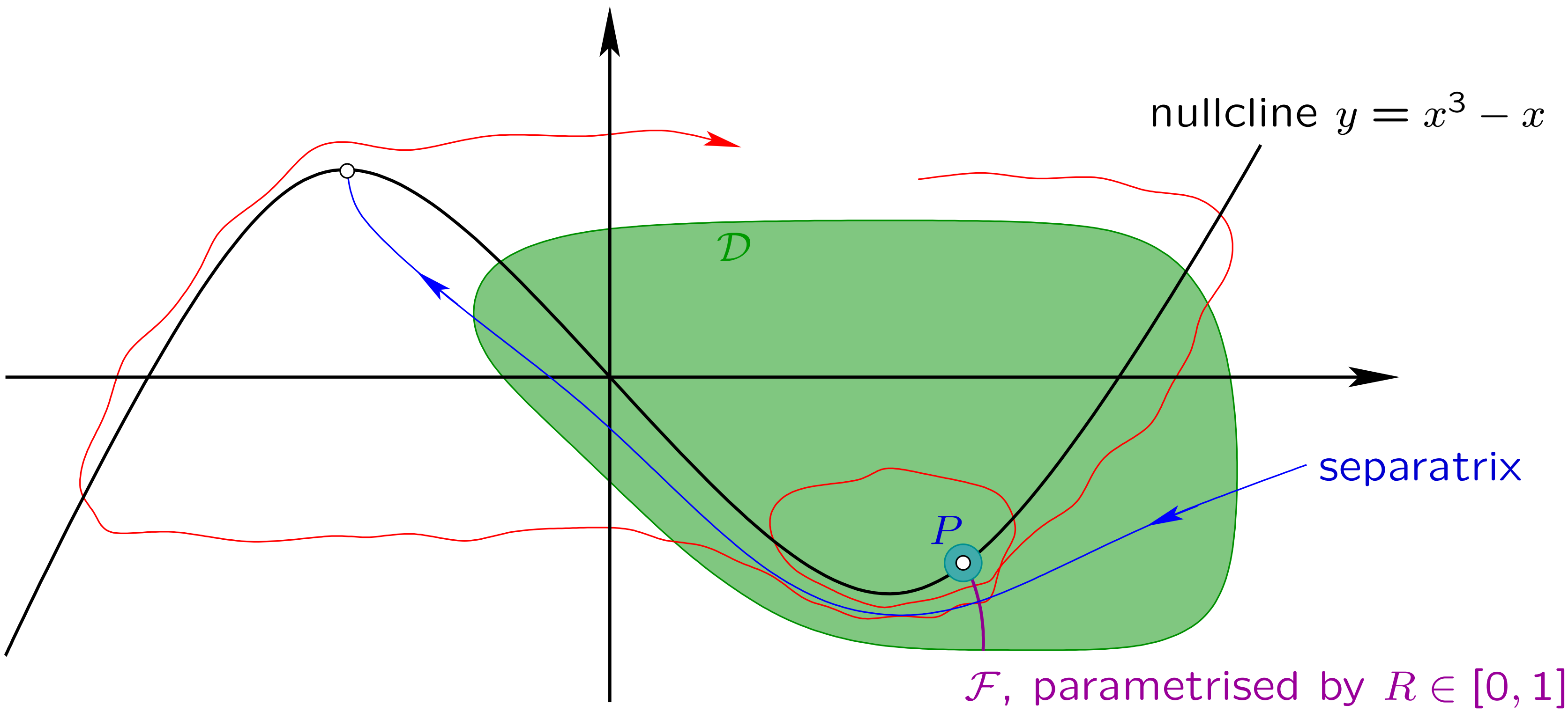
▷ Dynamics near singular Hopf bifurcation: To do



Small-amplitude oscillations (SAOs)

Definition of random number of SAOs N:



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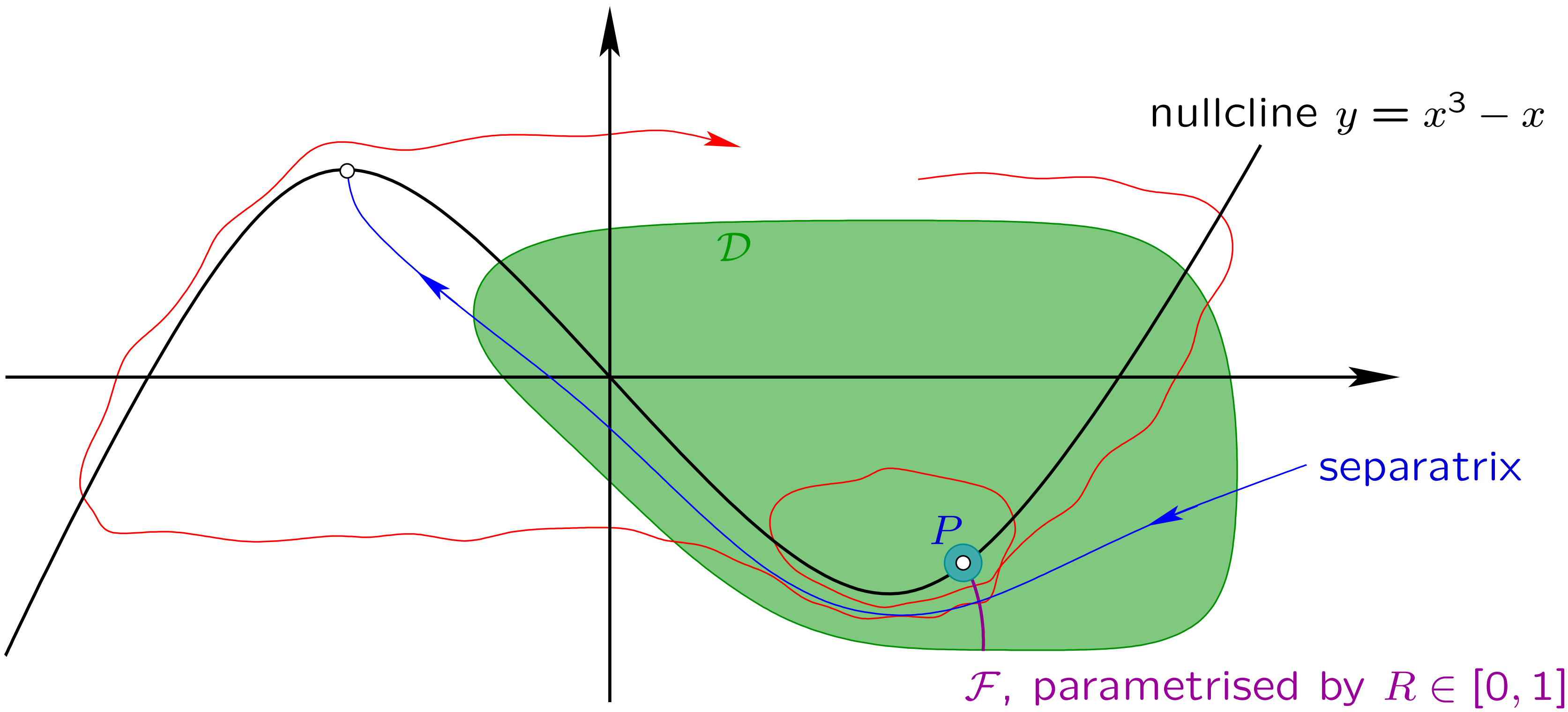

$(R_0, R_1, \dots, R_{N-1})$ substochastic Markov chain with kernel

$$K(R_0, A) = \mathbb{P}^{R_0}\{R_\tau \in A\}$$

$R \in \mathcal{F}$, $A \subset \mathcal{F}$, τ = first-hitting time of $\mathcal{F}$ (after turning around P)

N = number of turns around P until leaving $\mathcal{D}$

General results on distribution of SAOs

General theory of continuous-space Markov chains: [Orey '71, Nummelin '84]

Principal eigenvalue: eigenvalue λ_0 of K of largest module. $\lambda_0 \in \mathbb{R}$

Quasistationary distribution: prob. measure π_0 s.t. $\pi_0 K = \lambda_0 \pi_0$

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Theorem 1: [B & Landon, 2011] Assume $\sigma_1, \sigma_2 > 0$

- $\lambda_0 < 1$
- K admits quasistationary distribution π_0
- N is almost surely finite
- N is asymptotically geometric:

$$\lim_{n\to\infty} \mathbb{P}\{N = n+1 | N > n\} = 1 - \lambda_0$$

- $\mathbb{E}[r^N] < \infty$ for $r < 1/\lambda_0$, so all moments of N are finite

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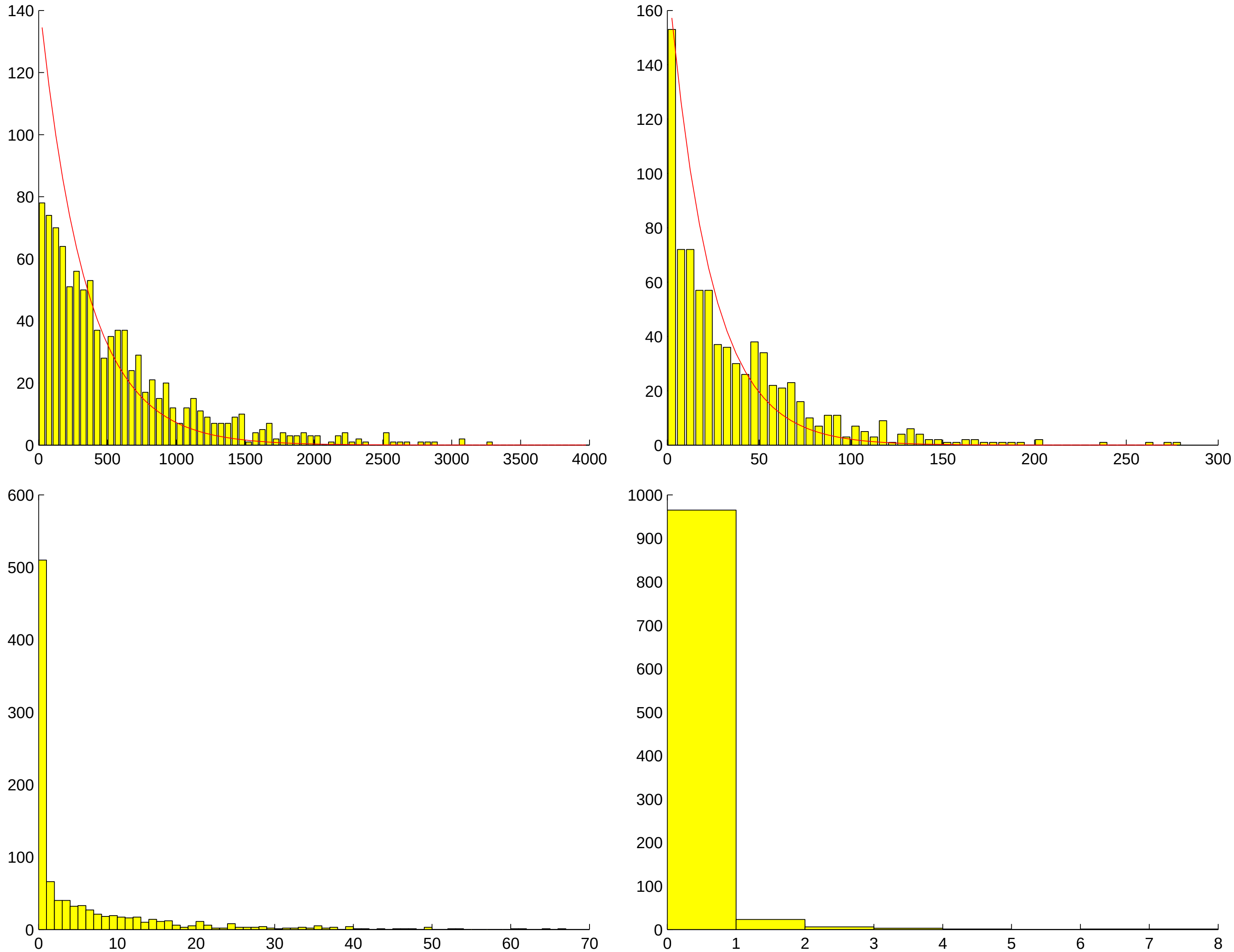
Proof uses Frobenius–Perron–Jentzsch–Krein–Rutman–Birkhoff theorem and uniform positivity of K, which implies spectral gap

Histograms of distribution of SAO number N (1000 spikes)



The weak-noise regime

Theorem 2: [B & Landon 2011]
Assume ε and $\delta/\sqrt{\varepsilon}$ sufficiently small
There exists $\kappa > 0$ s.t. for $\sigma^2 \leqslant (\varepsilon^{1/4}\delta)^2/\log(\sqrt{\varepsilon}/\delta)$

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$$\mathbb{E}^{\mu_0}[N] \geqslant C(\mu_0)\exp\left\{\kappa\frac{(\varepsilon^{1/4}\delta)^2}{\sigma^2}\right\}$$

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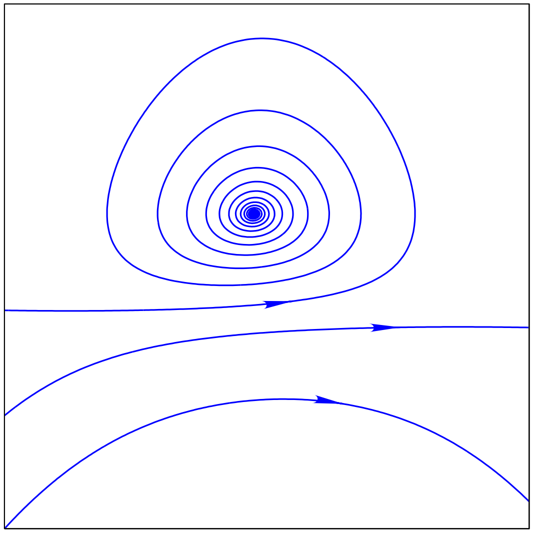
Proof:

▷ Construct $A \subset \mathcal{F}$ such that $K(x, A)$ exponentially close to 1 for all $x \in A$

▷ Use two different sets of coordinates to approximate K:
Near separatrix, and during SAO

Dynamics near the separatrix



Change of variables:

- ▷ Translate to Hopf bif. point
- ▷ Scale space and time
- ▷ Straighten nullcline $\dot{x} = 0$

⇒ variables (ξ, z) where nullcline: $\{z = \frac{1}{2}\}$

$$\mathrm{d}\xi_t = \left(\frac{1}{2} - z_t - \frac{\sqrt{\varepsilon}}{3}\xi_t^3\right)\mathrm{d}t + \tilde{\sigma}_1\,\mathrm{d}W_t^{(1)}$$

$$\mathrm{d}z_t = \left(\tilde{\mu} + 2\xi_t z_t + \frac{2\sqrt{\varepsilon}}{3}\xi_t^4\right)\mathrm{d}t - 2\tilde{\sigma}_1\xi_t\,\mathrm{d}W_t^{(1)} + \tilde{\sigma}_2\,\mathrm{d}W_t^{(2)}$$

where

$$\tilde{\mu} = \frac{\delta}{\sqrt{\varepsilon}} - \tilde{\sigma}_1^2 \qquad \tilde{\sigma}_1 = -\sqrt{3}\frac{\sigma_1}{\varepsilon^{3/4}} \qquad \tilde{\sigma}_2 = \sqrt{3}\frac{\sigma_2}{\varepsilon^{3/4}}$$

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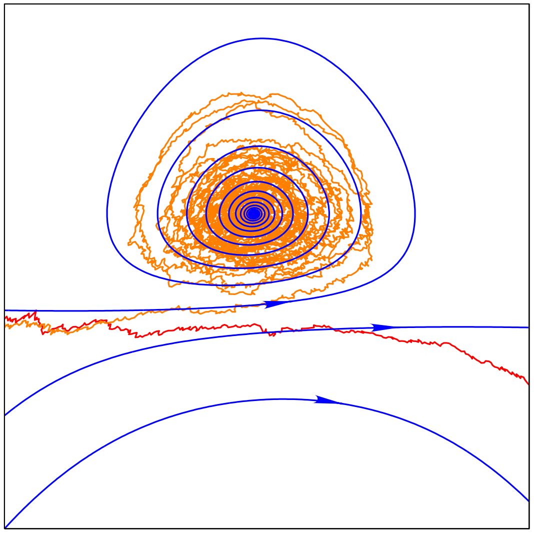

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$$d\xi_t = \left(\frac{1}{2} - z_t - \frac{\sqrt{\varepsilon}}{3}\xi_t^3\right) dt + \tilde{\sigma}_1 \, dW_t^{(1)}$$

$$dz_t = \left(\tilde{\mu} + 2\xi_t z_t + \frac{2\sqrt{\varepsilon}}{3}\xi_t^4\right) dt - 2\tilde{\sigma}_1 \xi_t \, dW_t^{(1)} + \tilde{\sigma}_2 \, dW_t^{(2)}$$

where

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Upward drift dominates if $\tilde{\mu}^2 \gg \tilde{\sigma}_1^2 + \tilde{\sigma}_2^2 \Rightarrow (\varepsilon^{1/4}\delta)^2 \gg \sigma_1^2 + \sigma_2^2$

Rotation around P: use that $2z \, e^{-2z-2\xi^2+1}$ is constant for $\tilde{\mu} = \varepsilon = 0$

Transition from weak to strong noise

Linear approximation:

$$\mathrm{d}z_t^0 = \left(\tilde{\mu} + t z_t^0\right) \mathrm{d}t - \tilde{\sigma}_1 t \,\mathrm{d}W_t^{(1)} + \tilde{\sigma}_2 \,\mathrm{d}W_t^{(2)}$$

$$\Rightarrow \quad \mathbb{P}\{\text{no SAO}\} \simeq \Phi\left(-\pi^{1/4}\frac{\tilde{\mu}}{\sqrt{\tilde{\sigma}_1^2+\tilde{\sigma}_2^2}}\right) \qquad \Phi(x) = \int_{-\infty}^{x} \frac{\mathrm{e}^{-y^2/2}}{\sqrt{2\pi}}\,\mathrm{d}y$$

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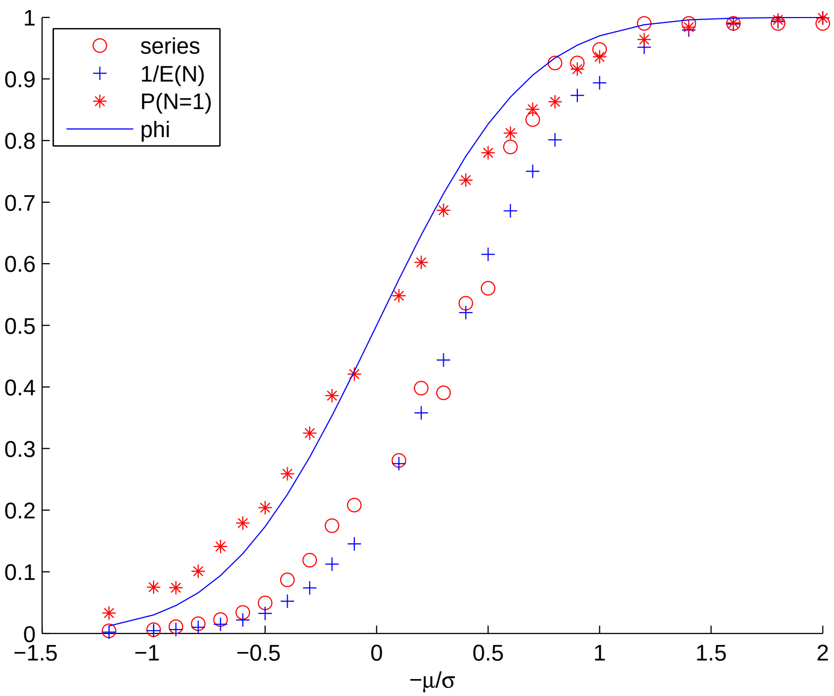


$*$: $\mathbb{P}\{\text{no SAO}\}$

$+$: $1/\mathbb{E}[N]$

$\circ$: $1 - \lambda_0$

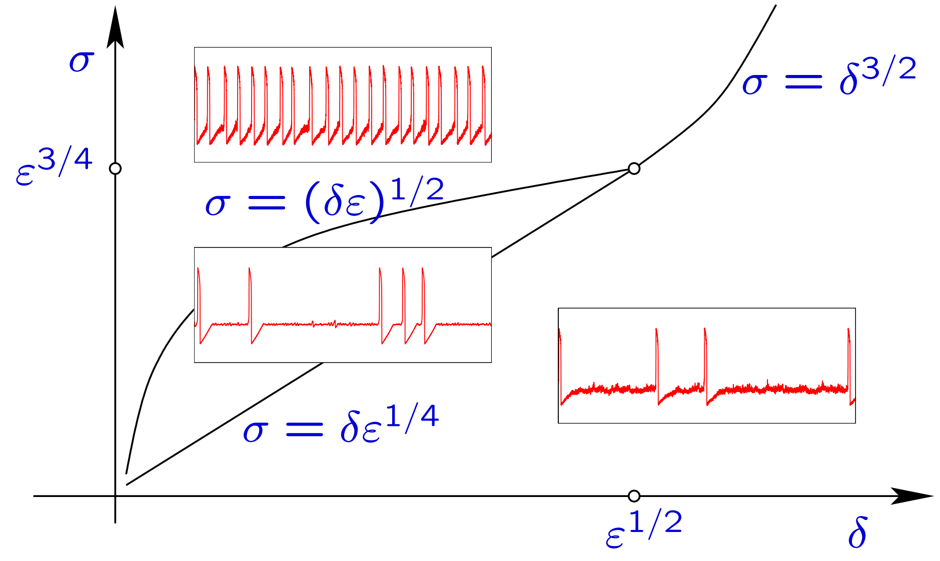
curve: $x \mapsto \Phi(\pi^{1/4} x)$

$$x = \frac{\tilde{\mu}}{\sqrt{\tilde{\sigma}_1^2 + \tilde{\sigma}_2^2}} = \frac{\varepsilon^{1/4}(\delta - \sigma_1^2/\varepsilon)}{\sqrt{\sigma_1^2 + \sigma_2^2}}$$

Conclusions

Three regimes for $\delta < \sqrt{\varepsilon}$:

- ▷ $\sigma \ll \varepsilon^{1/4}\delta$: rare isolated spikes
  interval $\simeq \mathcal{E}xp(\sqrt{\varepsilon}\,\mathrm{e}^{-(\varepsilon^{1/4}\delta)^2/\sigma^2})$
- ▷ $\varepsilon^{1/4}\delta \ll \sigma \ll \varepsilon^{3/4}$: transition
  geometric number of SAOs
  $\sigma = (\delta\varepsilon)^{1/2}$: geometric(1/2)
- ▷ $\sigma \gg \varepsilon^{3/4}$: repeated spikes



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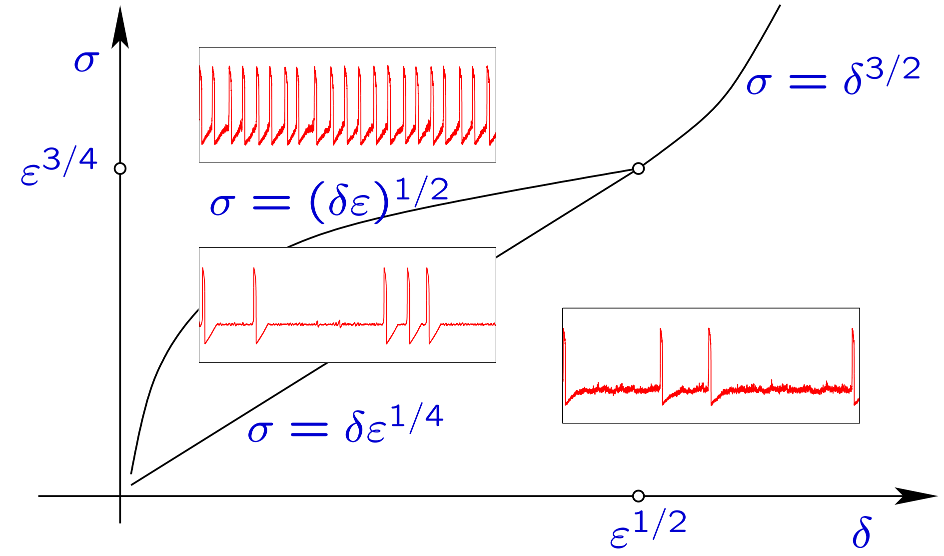

Outlook

▷ sharper bounds on λ_0 (and π_0)

▷ relation between $\mathbb{P}\{$no SAO$\}$, $1/\mathbb{E}[N]$ and $1 - \lambda_0$

▷ consequences of postspike distribution $\mu_0 \neq \pi_0$

▷ interspike interval distribution $\simeq$ periodically modulated exponential – how is it modulated?

Part II

Where noise modifies or suppresses MMOs

Folded node singularity

Normal form [Benoît, Lobry '82, Szmolyan, Wechselberger '01]:

$$
\begin{aligned}
\epsilon \dot{x} &= y - x^2 \\
\dot{y} &= -(\mu + 1)x - z \qquad\qquad (+ \text{ higher-order terms}) \\
\dot{z} &= \frac{\mu}{2}
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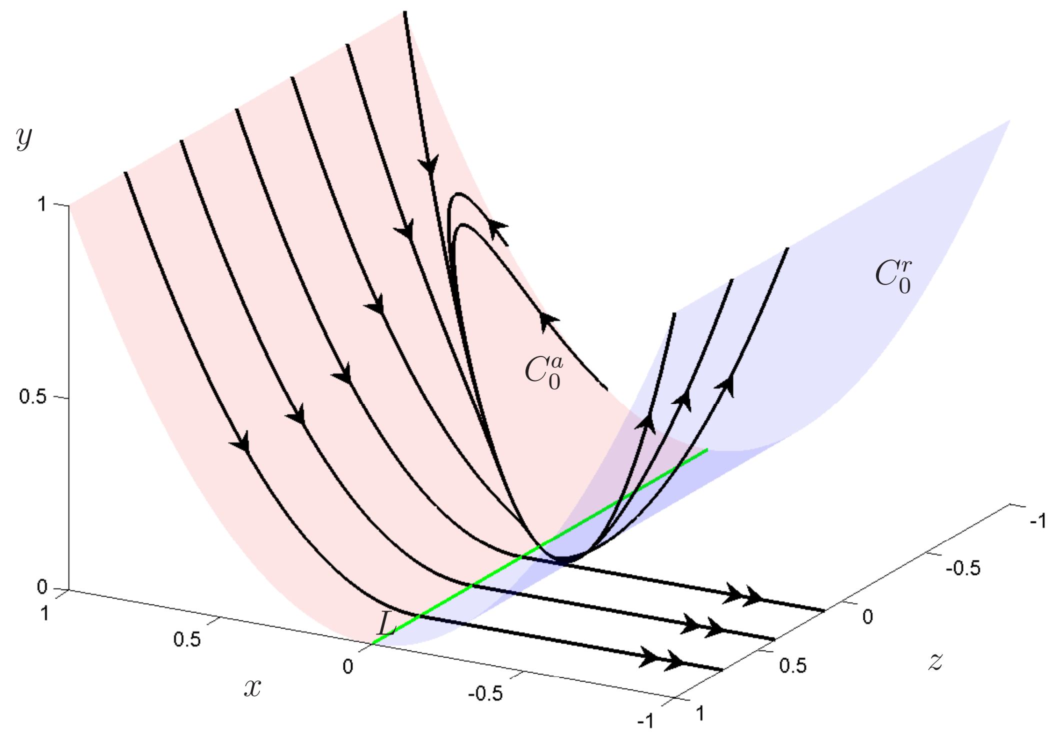
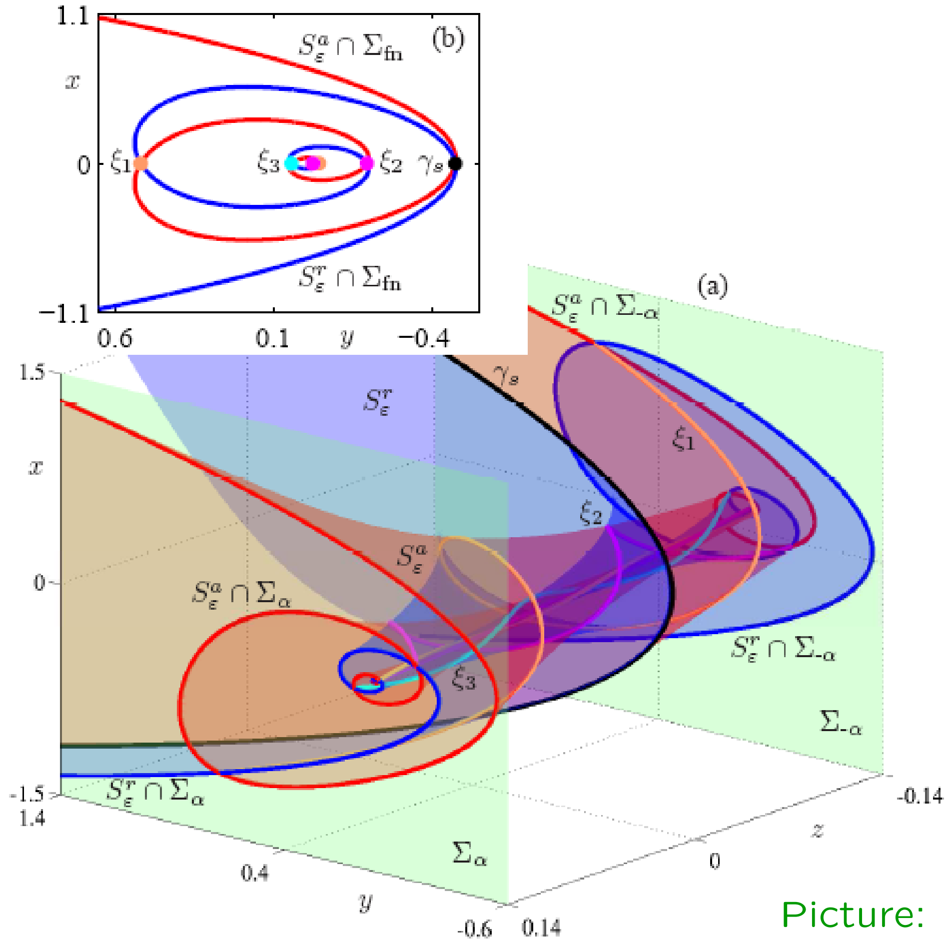
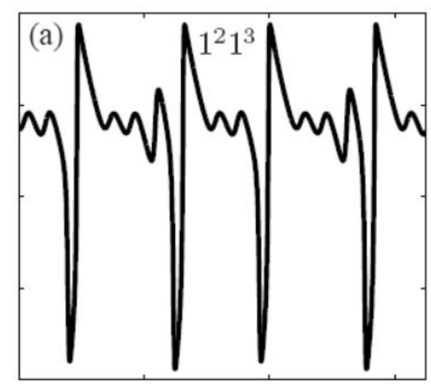


Folded node singularity

Theorem [Benoît, Lobry '82, Szmolyan, Wechselberger '01]:
For $2k+1 < \mu^{-1} < 2k+3$, the system admits k canard solutions
The j^{th} canard makes $(2j+1)/2$ oscillations



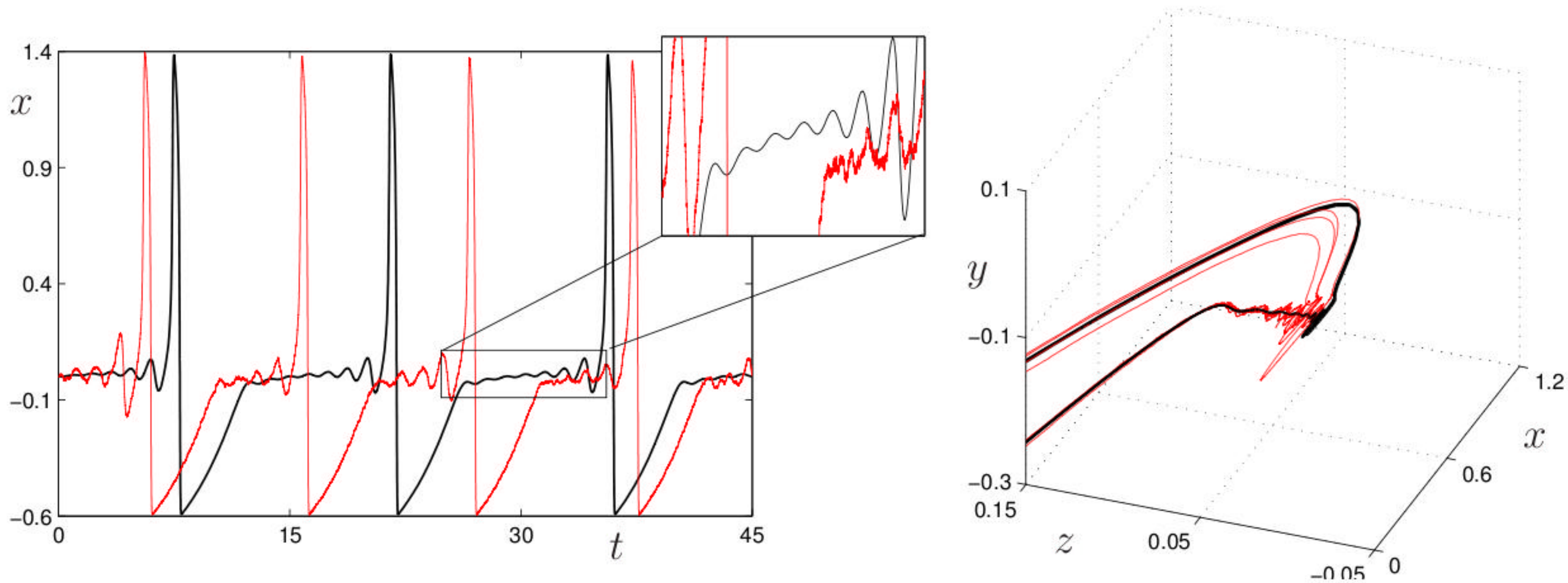


Mixed-mode oscillations (MMOs)

Picture: Mathieu Desroches

Effect of noise

$$
\begin{aligned}
\mathrm{d}x_t &= \frac{1}{\varepsilon}(y_t - x_t^2)\,\mathrm{d}t + \frac{\sigma}{\sqrt{\varepsilon}}\,\mathrm{d}W_t^{(1)} \\
\mathrm{d}y_t &= [-(\mu+1)x_t - z_t]\,\mathrm{d}t + \sigma\,\mathrm{d}W_t^{(2)} \\
\mathrm{d}z_t &= \frac{\mu}{2}\,\mathrm{d}t
\end{aligned}
$$



- Noise smears out small amplitude oscillations
- Early transitions modify the mixed-mode pattern

Covariance tubes

Linearized stochastic equation around a canard $(x_t^{\text{det}}, y_t^{\text{det}}, z_t^{\text{det}})$

$$\mathrm{d}\zeta_t = A(t)\zeta_t \,\mathrm{d}t + \sigma \,\mathrm{d}W_t \qquad A(t) = \begin{pmatrix} -2x_t^{\text{det}} & 1 \\ -(1+\mu) & 0 \end{pmatrix}$$

$$\zeta_t = U(t)\zeta_0 + \sigma \int_0^t U(t,s)\,\mathrm{d}W_s$$

($U(t,s)$: principal solution of $\dot{U} = AU$)

Gaussian process with covariance matrix

$$\mathrm{Cov}(\zeta_t) = \sigma^2 V(t) \qquad V(t) = U(t)V(0)U(t)^{-1} + \int_0^t U(t,s)U(t,s)^T \,\mathrm{d}s$$

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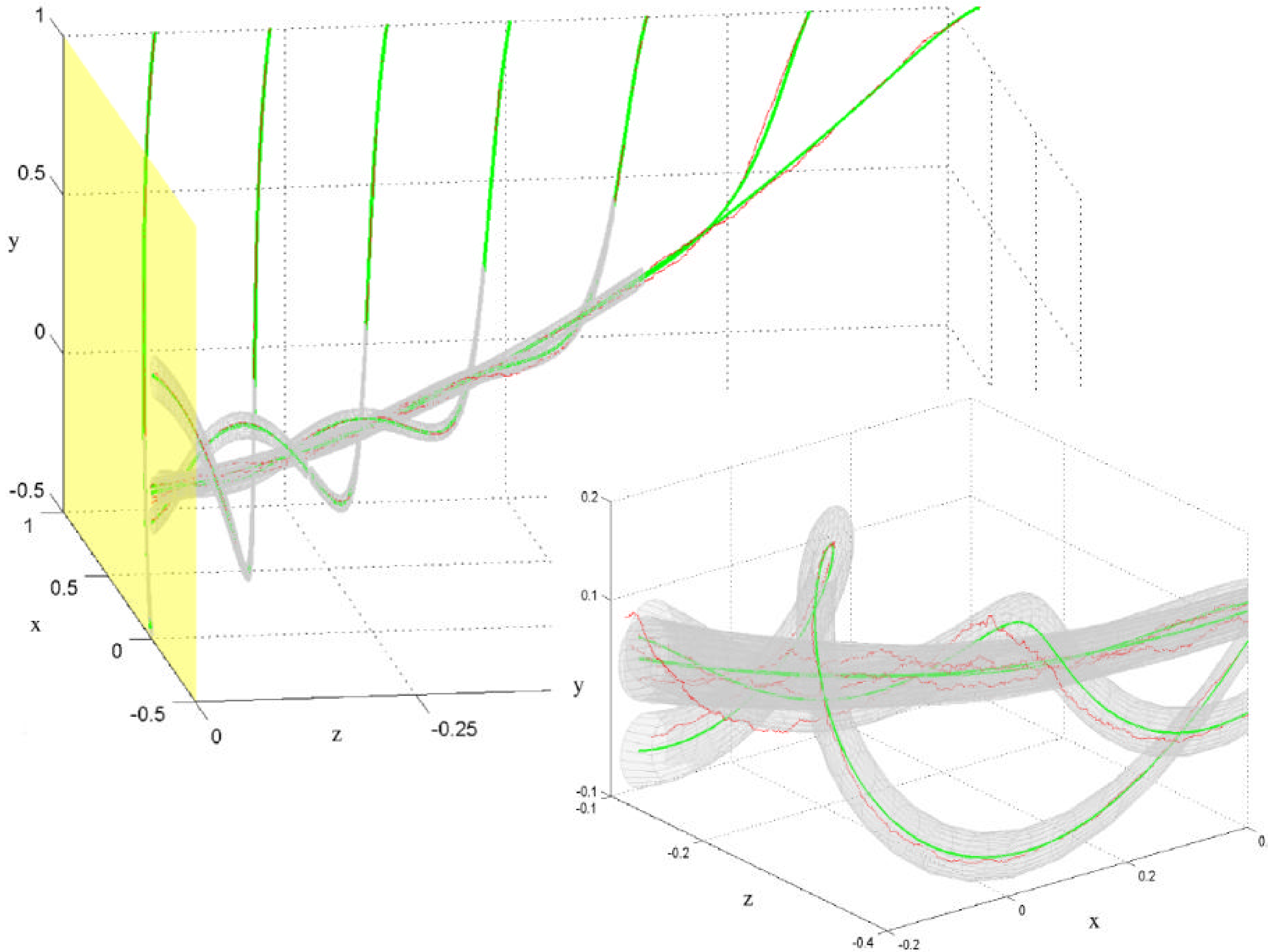
Covariance tube :

$$\mathcal{B}(h) = \Big\{ \langle (x,y) - (x_t^{\text{det}}, y_t^{\text{det}}), V(t)^{-1}[(x,y) - (x_t^{\text{det}}, y_t^{\text{det}})] \rangle < h^2 \Big\}$$

Theorem 3: [B, Gentz, Kuehn 2010]
Probability of leaving covariance tube before time t (with $z_t \leqslant 0$) :

$$\mathbb{P}\Big\{ \tau_{\mathcal{B}(h)} < t \Big\} \leqslant C(t)\,\mathrm{e}^{-\kappa h^2/2\sigma^2}$$



17-b

Small-amplitude oscillations and noise

One shows that for $z = 0$

- ▷ The distance between the k^{th} and $k+1^{\text{st}}$ canard has order $\mathrm{e}^{-(2k+1)^2\mu}$
- ▷ The section of $\mathcal{B}(h)$ is close to circular with radius $\mu^{-1/4}h$

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Corollary:
Let
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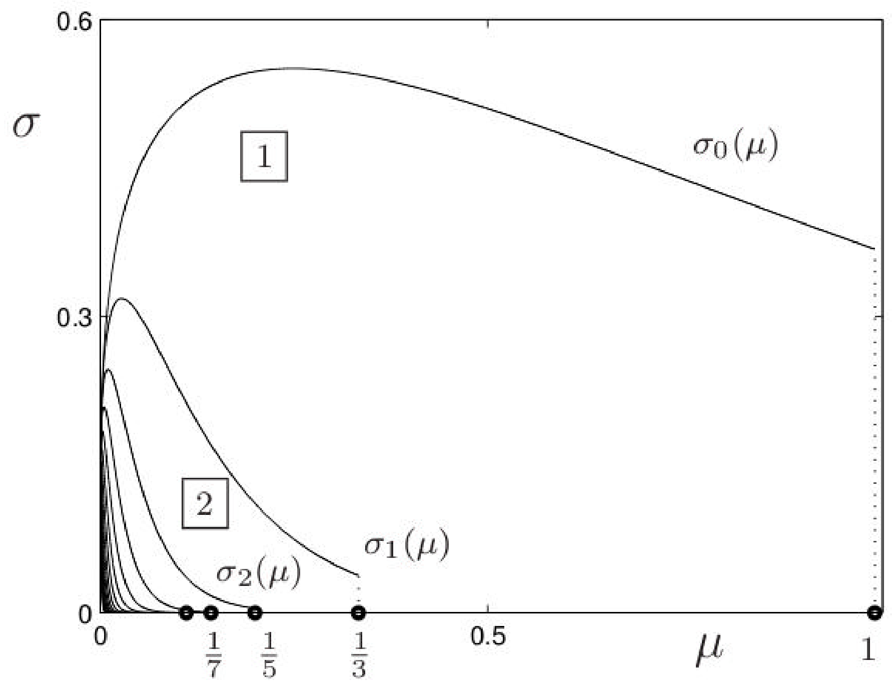
Canards with $\frac{2k+1}{4}$ oscillations become indistinguishable from noisy fluctuations for $\sigma > \sigma_k(\mu)$



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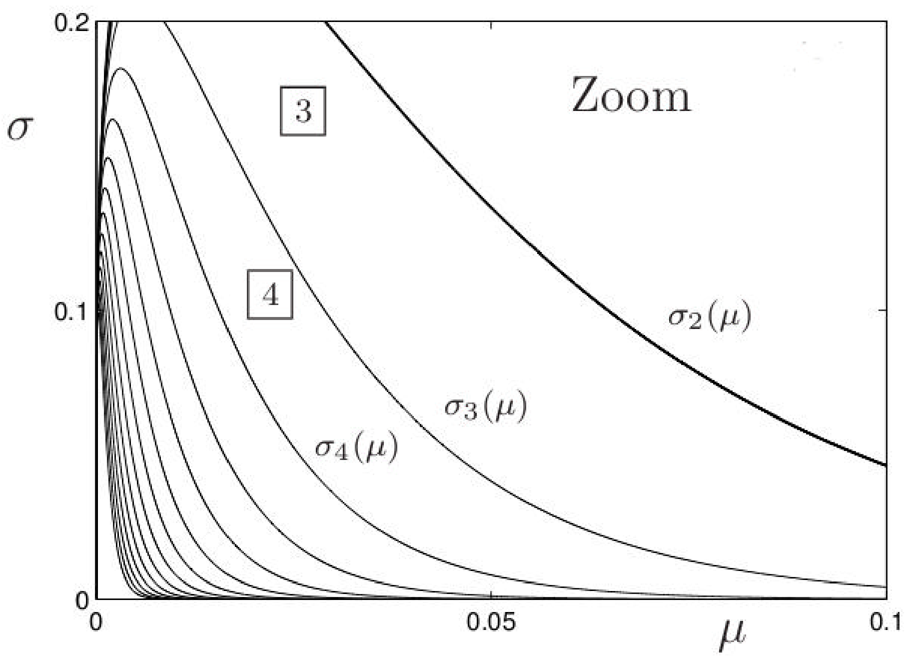

Early transitions

Let $\mathcal{D}$ be neighbourhood of size $\sqrt{z}$ of a canard for $z > 0$ (unstable)

Theorem 4: [B, Gentz, Kuehn 2010]
$\exists \kappa, C, \gamma_1, \gamma_2 > 0$ such that for $\sigma|\log\sigma|^{\gamma_1} \leqslant \mu^{3/4}$ probability of leaving $\mathcal{D}$ after $z_t = z$ satisfies

$$\mathbb{P}\Big\{z_{\tau_{\mathcal{D}}} > z\Big\} \leqslant C|\log\sigma|^{\gamma_2}\,\mathrm{e}^{-\kappa(z^2-\mu)/(\mu|\log\sigma|)}$$

Small for $z \gg \sqrt{\mu|\log\sigma|/\kappa}$

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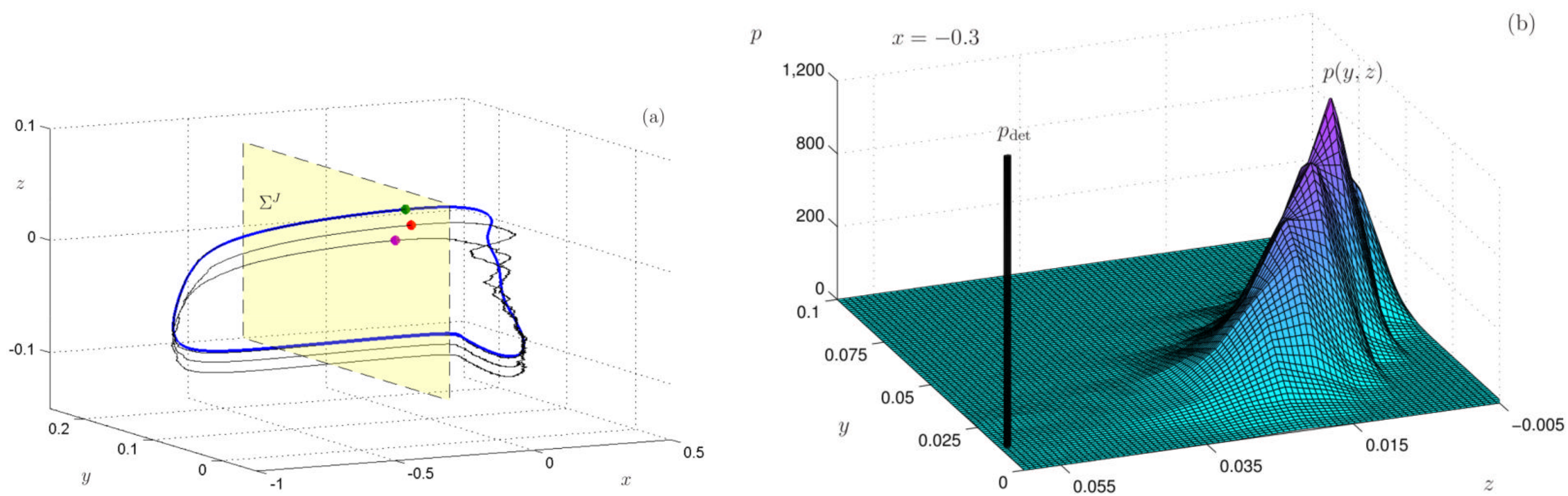

Further work

▷ Better understanding of distribution of noise-induced transitions

▷ Effect on mixed-mode pattern in conjunction with global return mechanism

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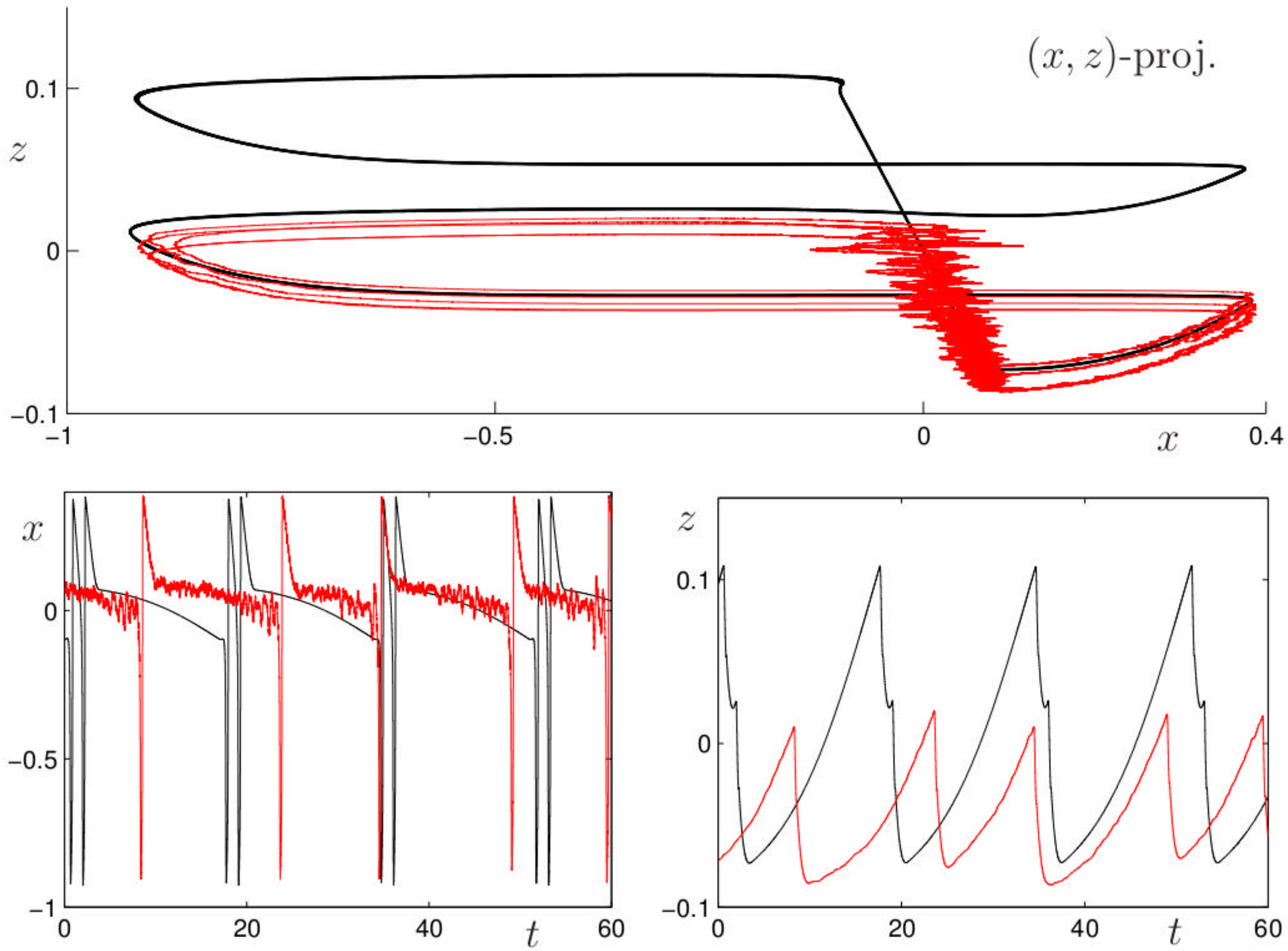
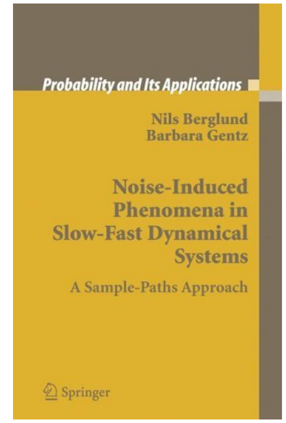
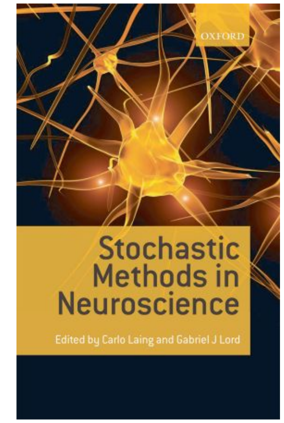


Further reading

N.B. and Barbara Gentz, *Noise-induced phenomena in slow-fast dynamical systems, A sample-paths approach*, Springer, Probability and its Applications (2006)

N.B. and Barbara Gentz, *Stochastic dynamic bifurcations and excitability*, in C. Laing and G. Lord, (Eds.), *Stochastic methods in Neuroscience*, p. 65-93, Oxford University Press (2009)





N.B., Barbara Gentz and Christian Kuehn, *Hunting French Ducks in a Noisy Environment*, `arXiv:1011.3193`, submitted (2010)

N.B. and Damien Landon, *Mixed-mode oscillations and interspike interval statistics in the stochastic FitzHugh–Nagumo model*, `arXiv:1105.1278`, submitted (2011)

N.B., *Kramers' law: Validity, derivations and generalisations*, `arXiv:1106.5799`, submitted (2011)

Appendix

Covariance tubes

Theorem 3: [B, Gentz, Kuehn 2010]

Probability of leaving covariance tube before time t (with $z_t \leqslant 0$) :

$$\mathbb{P}\Big\{\tau_{\mathcal{B}(h)} < t\Big\} \leqslant C(t)\, \mathrm{e}^{-\kappa h^2/2\sigma^2}$$

Sketch of proof :

▷ (Sub)martingale : $\{M_t\}_{t\geqslant 0}$, $\mathbb{E}\{M_t|M_s\} = (\geqslant) M_s$ for $t \geqslant s \geqslant 0$

▷ Doob's submartingale inequality : $\mathbb{P}\Big\{\sup_{0\leqslant t\leqslant T} M_t \geqslant L\Big\} \leqslant \frac{1}{L}\mathbb{E}[M_T]$

Covariance tubes

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- ▷ Linear equation : $\zeta_t = \sigma \displaystyle\int_0^t U(t,s)\,\mathrm{d}W_s$ is no martingale but can be approximated by martingale on small time intervals
- ▷ $\exp\{\gamma\langle\zeta_t, V(t)^{-1}\zeta_t\rangle\}$ approximated by submartingale
- ▷ Doob's inequality yields bound on probability of leaving $\mathcal{B}(h)$ during small time intervals. Then sum over all time intervals

Covariance tubes

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- ▷ Doob's inequality yields bound on probability of leaving $\mathcal{B}(h)$ during small time intervals. Then sum over all time intervals
- ▷ Nonlinear equation : $\mathrm{d}\zeta_t = A(t)\zeta_t\,\mathrm{d}t + b(\zeta_t,t)\,\mathrm{d}t + \sigma\,\mathrm{d}W_t$

$$\zeta_t = \sigma\int_0^t U(t,s)\,\mathrm{d}W_s + \int_0^t U(t,s)b(\zeta_s,s)\,\mathrm{d}s$$

Second integral can be treated as small perturbation for $t \leqslant \tau_{\mathcal{B}(h)}$

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Small for $z \gg \sqrt{\mu|\log\sigma|/\kappa}$

Sketch of proof :

- ▷ Escape from neighbourhood of size $\sigma|\log\sigma|/\sqrt{z}$:
  compare with linearized equation on small time intervals + Markov property
- ▷ Escape from annulus $\sigma|\log\sigma|/\sqrt{z} \leqslant \|\zeta\| \leqslant \sqrt{z}$:
  use polar coordinates and averaging
- ▷ To combine the two regimes : use Laplace transforms