Oscillations multimodales dans les équations différentielles stochastiques

Nils Berglund

MAPMO, Université d'Orléans CNRS, UMR 6628 et Fédération Denis Poisson www.univ-orleans.fr/mapmo/membres/berglund

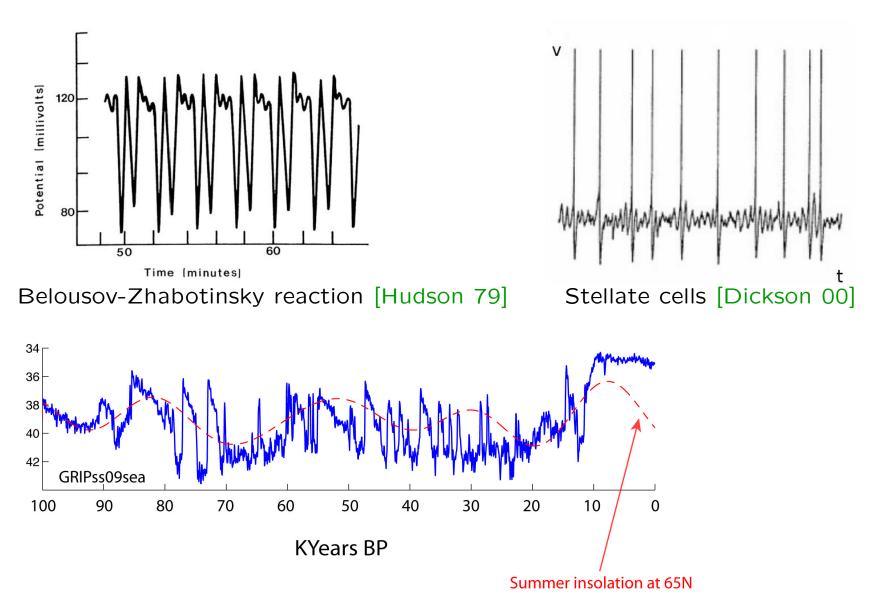
Collaborateurs:

Stéphane Cordier, Damien Landon, Simona Mancini, MAPMO, Orléans Barbara Gentz, University of Bielefeld Christian Kuehn, Max Planck Institute, Dresden

Projet ANR MANDy, Mathematical Analysis of Neuronal Dynamics

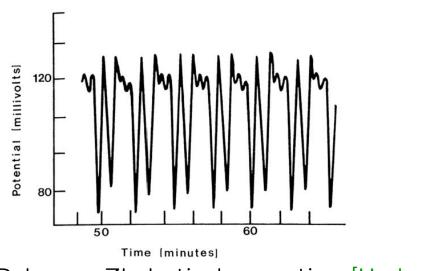
GdT Modélisation, LPMA, Paris, 5 mai 2011

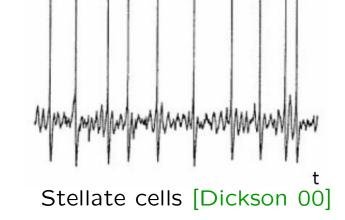
Oscillations in natural systems



Mean temperature based on ice core measurements [Johnson et al 01]

Oscillations in natural systems





Belousov-Zhabotinsky reaction [Hudson 79]

basa assillations avist

Deterministic models reproducing these oscillations exist
 and have been abundantly studied

They often involve singular perturbation theory

Noise may also induce oscillations not present in deterministic case

Example: Van der Pol oscillator $x'' + \varepsilon^{-1/2}(x^2 - 1)x' + x = 0$

$$x'' + \varepsilon^{-1/2}(x^2 - 1)x' + x = 0$$

$$\dot{x} = y + x - \frac{1}{3}x^{3} \qquad \qquad t \mapsto \varepsilon t \qquad \qquad \varepsilon \dot{x} = y + x - \frac{1}{3}x^{3}$$

$$\dot{y} = -\varepsilon x \qquad \qquad \dot{y} = -x$$

Example: Van der Pol oscillator

$$x'' + \varepsilon^{-1/2}(x^2 - 1)x' + x = 0$$

$$\dot{x} = y + x - \frac{1}{3}x^{3} \qquad \longleftrightarrow \qquad \dot{\varepsilon}\dot{x} = y + x - \frac{1}{3}x^{3} \\
\dot{y} = -\varepsilon x \qquad \qquad \dot{y} = -x$$

$$\downarrow \qquad \varepsilon \to 0 \qquad \qquad \downarrow \qquad \varepsilon \to 0$$

$$\dot{x} = y + x - \frac{1}{3}x^{3} \qquad \longleftrightarrow \qquad y = -(x - \frac{1}{3}x^{3})$$

$$\dot{y} = 0 \qquad \qquad \dot{y} = -x$$

$$\Rightarrow \dot{x} = \frac{x}{1 - x^{2}}$$

Example: Van der Pol oscillator $x'' + \varepsilon^{-1/2}(x^2 - 1)x' + x = 0$

$$x'' + \varepsilon^{-1/2}(x^2 - 1)x' + x = 0$$

$$\dot{x} = y + x - \frac{1}{3}x^{3} \qquad \qquad t \mapsto \varepsilon t \qquad \qquad \varepsilon \dot{x} = y + x - \frac{1}{3}x^{3}$$

$$\dot{y} = -\varepsilon x \qquad \qquad \dot{y} = -x$$

$$\varepsilon \rightarrow 0$$

$$\varepsilon \rightarrow 0$$

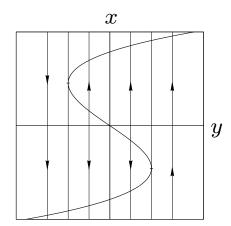
$$\dot{x} = y + x - \frac{1}{3}x^3 \qquad \iff \qquad y = -(x - \frac{1}{3}x^3)$$

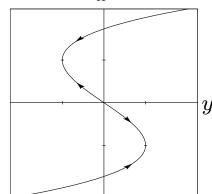
$$\dot{y} = 0 \qquad \qquad \dot{y} = -x$$

$$y = -\left(x - \frac{1}{3}x^3\right)$$

$$\dot{y} = -x$$

$$\Rightarrow \dot{x} = \frac{x}{1 - x^2}$$



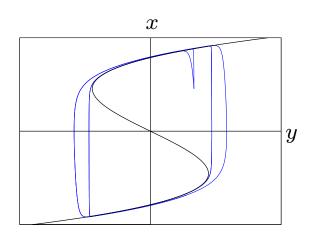


Example: Van der Pol oscillator

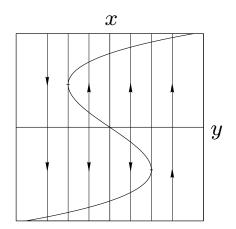
$$x'' + \varepsilon^{-1/2}(x^2 - 1)x' + x = 0$$

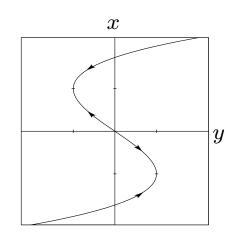
$$\dot{x} = y + x - \frac{1}{3}x^3$$

$$\dot{y} = -\varepsilon x$$



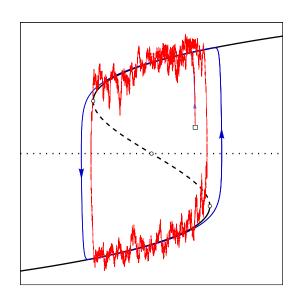
Relaxation oscillations





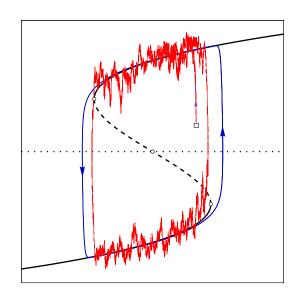
Effect of noise on the Van der Pol oscillator

$$dx_t = \left[y_t + x_t - \frac{x_t^3}{3} \right] dt + \sigma dW_t$$
$$dy_t = -\varepsilon x_t dt$$



Effect of noise on the Van der Pol oscillator

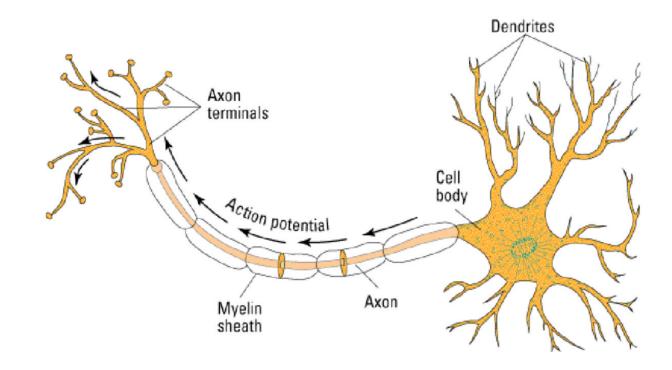
$$dx_t = \left[y_t + x_t - \frac{x_t^3}{3} \right] dt + \sigma dW_t$$
$$dy_t = -\varepsilon x_t dt$$



Theorem [B & Gentz 2006]

- \bullet $\sigma<\sqrt{\varepsilon}$: Cycles comparable to deterministic ones with probability $1-\mathcal{O}(\mathrm{e}^{-\varepsilon/\sigma^2})$
- $\sigma > \sqrt{\varepsilon}$: Cycles are smaller, by $\mathcal{O}(\sigma^{4/3})$, than deterministic cycles, with probability $1 \mathcal{O}(\mathrm{e}^{-\sigma^2/\varepsilon|\log\sigma|})$

Neuron



- > Single neuron communicates by generating action potential
- ▷ Excitable: small change in parameters yields spike generation
- ▶ May display Mixed-Mode Oscillations (MMOs) and Relaxation Oscillations

Hodgkin-Huxley model (1952)

$$C\dot{v} = -\sum_i \bar{g}_i \varphi_i^{\alpha_i} \chi_i^{\beta_i} (v - v_i^*)$$
 voltage
$$\tau_{\varphi,i}(v) \dot{\varphi}_i = -(\varphi_i - \varphi_i^*(v))$$
 activation
$$\tau_{\chi,i}(v) \dot{\chi}_i = -(\chi_i - \chi_i^*(v))$$
 inactivation

 $\triangleright i \in \{\text{Na}^+, \text{K}^+, \dots\}$ describes different types of ion channels $\triangleright \varphi_i^*(v), \chi_i^*(v)$ sigmoïdal functions, e.g. $\tanh(av + b)$

Hodgkin-Huxley model (1952)

$$C\dot{v} = -\sum_{i} \bar{g}_{i} \varphi_{i}^{\alpha_{i}} \chi_{i}^{\beta_{i}} (v - v_{i}^{*}) \qquad \text{voltage}$$

$$\tau_{\varphi,i}(v)\dot{\varphi}_{i} = -(\varphi_{i} - \varphi_{i}^{*}(v)) \qquad \text{activation}$$

$$\tau_{\chi,i}(v)\dot{\chi}_{i} = -(\chi_{i} - \chi_{i}^{*}(v)) \qquad \text{inactivation}$$

 $\triangleright i \in \{\text{Na}^+, \text{K}^+, \dots\}$ describes different types of ion channels $\triangleright \varphi_i^*(v), \chi_i^*(v)$ sigmoïdal functions, e.g. $\tanh(av + b)$

For $C/\bar{g}_i \ll \tau_{x,i}$: slow-fast systems of the form

$$\varepsilon \dot{v} = f(v, w)
\dot{w}_i = g_i(v, w)$$

Fitzhugh-Nagumo model (1962)

$$\varepsilon \dot{x} = x - x^3 + y$$
$$\dot{y} = \alpha - \beta x - \gamma y$$

Fitzhugh-Nagumo model (1962)

$$\varepsilon \dot{x} = x - x^3 + y$$

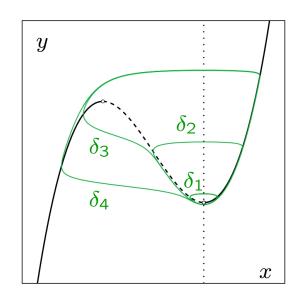
$$\dot{y} = \alpha - \beta x - \gamma y$$

$$= \frac{1}{\sqrt{3}} + \delta - x$$

The canard (french duck) phenomenon

[J.-L. Callot, F. Diener, M. Diener (1978), E. Benoît (1981), ...]

$$\varepsilon = 0.05$$
 $\alpha = \frac{1}{\sqrt{3}} + \delta$
 $\beta = 1$
 $\gamma = 0$
 $\delta_1 = -0.003$
 $\delta_2 = -0.003765458$
 $\delta_3 = -0.003765459$
 $\delta_4 = -0.005$



Fitzhugh-Nagumo model (1962)

$$\varepsilon \dot{x} = x - x^3 + y$$

$$\dot{y} = \alpha - \beta x - \gamma y$$

$$= \frac{1}{\sqrt{3}} + \delta - x$$

The canard (french duck) phenomenon

[J.-L. Callot, F. Diener, M. Diener (1978), E. Benoît (1981), ...]

$$\varepsilon = 0.05$$
 $\alpha = \frac{1}{\sqrt{3}} + \delta$
 $\beta = 1$
 $\gamma = 0$
 $\delta_1 = -0.003$
 $\delta_2 = -0.003765458$
 $\delta_3 = -0.003765459$
 $\delta_4 = -0.005$

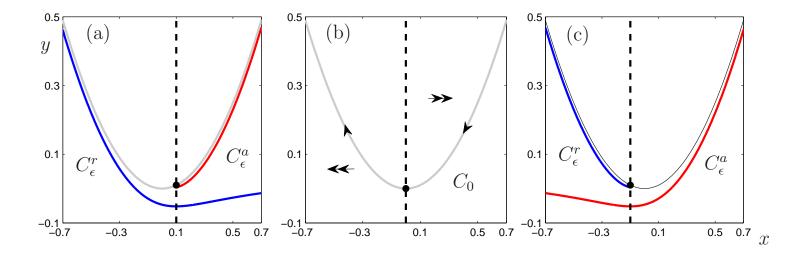


The canard (french duck) phenomenon

Normal form near fold point

$$\varepsilon \dot{x} = y - x^{2}$$

$$\dot{y} = \delta - x$$
(+ higher-order terms)



Folded node singularity

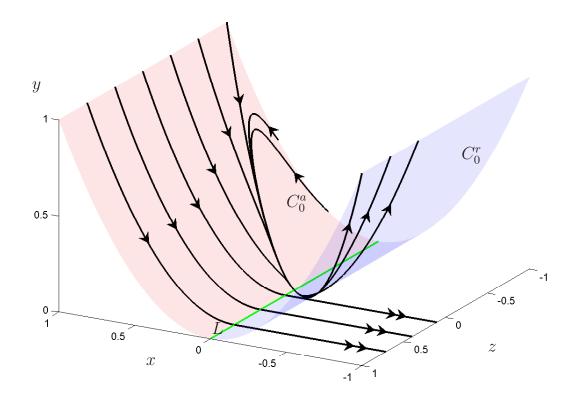
Normal form [Benoît, Lobry '82, Szmolyan, Wechselberger '01]:

$$\begin{split} \epsilon \dot{x} &= y - x^2 \\ \dot{y} &= -(\mu + 1)x - z \\ \dot{z} &= \frac{\mu}{2} \end{split} \tag{+ higher-order terms)}$$

Folded node singularity

Normal form [Benoît, Lobry '82, Szmolyan, Wechselberger '01]:

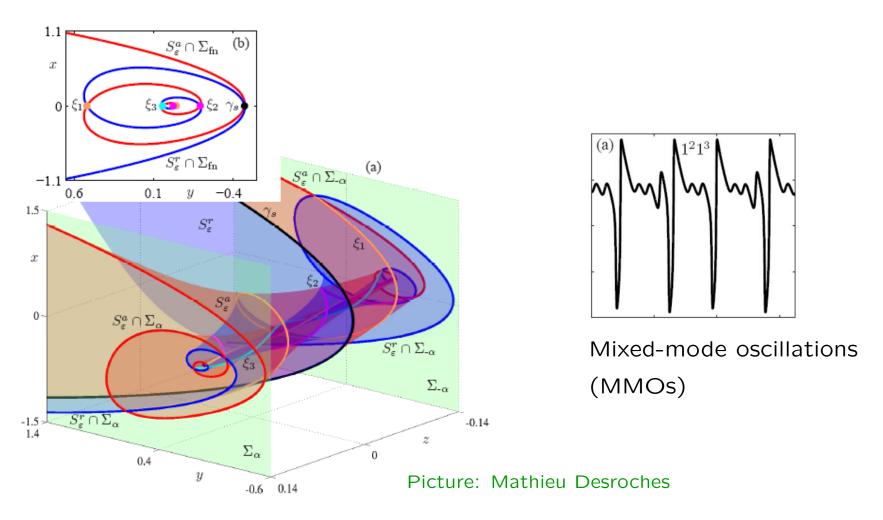
$$\begin{split} \epsilon \dot{x} &= y - x^2 \\ \dot{y} &= -(\mu + 1)x - z \\ \dot{z} &= \frac{\mu}{2} \end{split} \tag{+ higher-order terms)}$$



Folded node singularity

Theorem [Benoît, Lobry '82, Szmolyan, Wechselberger '01]:

For $2k+1<\mu^{-1}<2k+3$, the system admits k canard solutions The j^{th} canard makes (2j+1)/2 oscillations

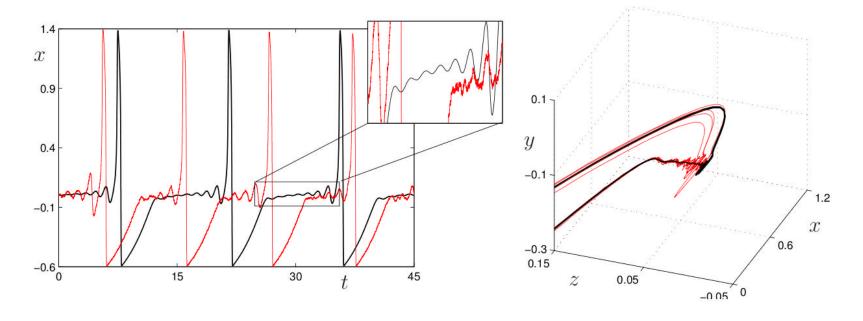


Effect of noise

$$dx_t = \frac{1}{\varepsilon} (y_t - x_t^2) dt + \frac{\sigma}{\sqrt{\varepsilon}} dW_t^{(1)}$$

$$dy_t = [-(\mu + 1)x_t - z_t] dt + \sigma dW_t^{(2)}$$

$$dz_t = \frac{\mu}{2} dt$$



- Noise smears out small amplitude oscillations
- Early transitions modify the mixed-mode pattern

Linearized stochastic equation around a canard $(x_t^{\text{det}}, y_t^{\text{det}}, z_t^{\text{det}})$

$$d\zeta_t = A(t)\zeta_t dt + \sigma dW_t \qquad A(t) = \begin{pmatrix} -2x_t^{\text{det}} & 1\\ -(1+\mu) & 0 \end{pmatrix}$$

$$\zeta_t = U(t)\zeta_0 + \sigma \int_0^t U(t,s) \, dW_s$$
 $(U(t,s) : principal solution of $\dot{U} = AU)$$

Gaussian process with covariance matrix

$$Cov(\zeta_t) = \sigma^2 V(t)$$
 $V(t) = U(t)V(0)U(t)^{-1} + \int_0^t U(t,s)U(t,s)^T ds$

Linearized stochastic equation around a canard $(x_t^{\text{det}}, y_t^{\text{det}}, z_t^{\text{det}})$

$$d\zeta_t = A(t)\zeta_t dt + \sigma dW_t \qquad A(t) = \begin{pmatrix} -2x_t^{\text{det}} & 1\\ -(1+\mu) & 0 \end{pmatrix}$$

$$\zeta_t = U(t)\zeta_0 + \sigma \int_0^t U(t,s) \, dW_s$$
 $(U(t,s) : principal solution of $\dot{U} = AU)$$

Gaussian process with covariance matrix

$$Cov(\zeta_t) = \sigma^2 V(t)$$
 $V(t) = U(t)V(0)U(t)^{-1} + \int_0^t U(t,s)U(t,s)^T ds$

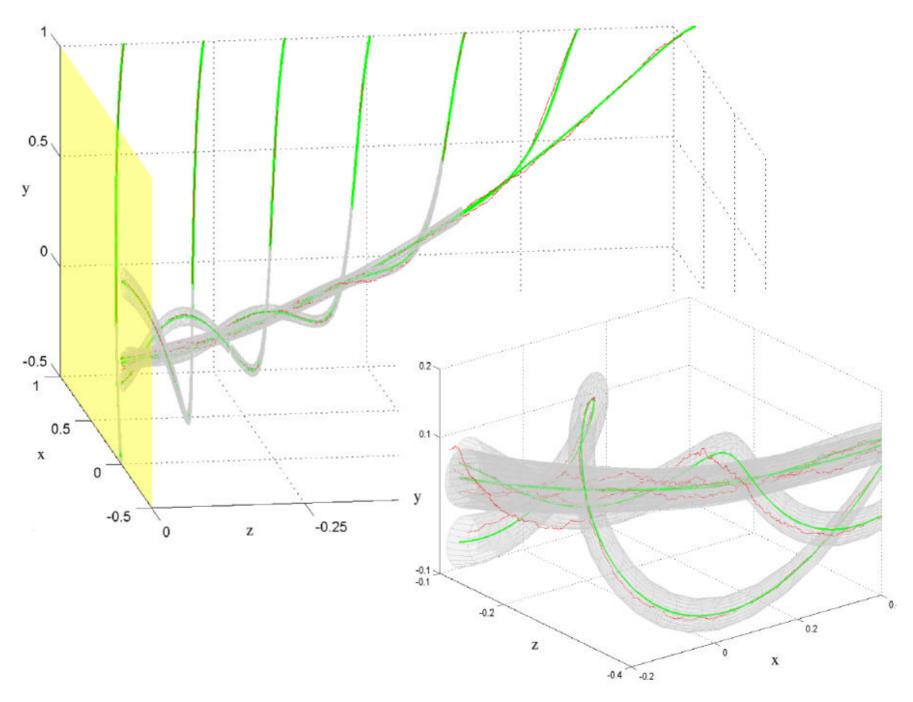
Covariance tube :

$$\mathcal{B}(h) = \left\{ \langle (x, y) - (x_t^{\text{det}}, y_t^{\text{det}}), V(t)^{-1} [(x, y) - (x_t^{\text{det}}, y_t^{\text{det}})] \rangle < h^2 \right\}$$

Theorem [B, Gentz, Kuehn 2010]

Probability of leaving covariance tube before time t (with $z_t \leq 0$):

$$\mathbb{P}\big\{\tau_{\mathcal{B}(h)} < t\big\} \leqslant C(t) \,\mathrm{e}^{-\kappa h^2/2\sigma^2}$$



Theorem [B, Gentz, Kuehn 2010]

Probability of leaving covariance tube before time t (with $z_t \leq 0$):

$$\mathbb{P}\left\{\tau_{\mathcal{B}(h)} < t\right\} \leqslant C(t) \,\mathrm{e}^{-\kappa h^2/2\sigma^2}$$

Sketch of proof:

- \triangleright (Sub)martingale : $\{M_t\}_{t\geqslant 0}$, $\mathbb{E}\{M_t|M_s\}=(\geqslant)M_s$ for $t\geqslant s\geqslant 0$
- ho Doob's submartingale inequality : $\mathbb{P}\Big\{\sup_{0\leqslant t\leqslant T}M_t\geqslant L\Big\}\leqslant rac{1}{L}\mathbb{E}[M_T]$

Theorem [B, Gentz, Kuehn 2010]

Probability of leaving covariance tube before time t (with $z_t \leq 0$):

$$\mathbb{P}\left\{\tau_{\mathcal{B}(h)} < t\right\} \leqslant C(t) \,\mathrm{e}^{-\kappa h^2/2\sigma^2}$$

Sketch of proof:

- \triangleright (Sub)martingale : $\{M_t\}_{t\geqslant 0}$, $\mathbb{E}\{M_t|M_s\}=(\geqslant)M_s$ for $t\geqslant s\geqslant 0$
- ho Doob's submartingale inequality : $\mathbb{P}\Big\{\sup_{0\leqslant t\leqslant T}M_t\geqslant L\Big\}\leqslant rac{1}{L}\mathbb{E}[M_T]$
- ho Linear equation : $\zeta_t = \sigma \int_0^t U(t,s) \, \mathrm{d}W_s$ is no martingale but can be approximated by martingale on small time intervals
- $\triangleright \exp\{\gamma\langle \zeta_t, V(t)^{-1}\zeta_t\rangle\}$ approximated by submartingale
- \triangleright Doob's inequality yields bound on probability of leaving $\mathcal{B}(h)$ during small time intervals. Then sum over all time intervals

Theorem [B, Gentz, Kuehn 2010]

Probability of leaving covariance tube before time t (with $z_t \leq 0$):

$$\mathbb{P}\left\{\tau_{\mathcal{B}(h)} < t\right\} \leqslant C(t) \,\mathrm{e}^{-\kappa h^2/2\sigma^2}$$

Sketch of proof:

- \triangleright (Sub)martingale : $\{M_t\}_{t\geqslant 0}$, $\mathbb{E}\{M_t|M_s\}=(\geqslant)M_s$ for $t\geqslant s\geqslant 0$
- ho Doob's submartingale inequality : $\mathbb{P}\Big\{\sup_{0\leqslant t\leqslant T}M_t\geqslant L\Big\}\leqslant rac{1}{L}\mathbb{E}[M_T]$
- ho Linear equation : $\zeta_t = \sigma \int_0^t U(t,s) \, \mathrm{d}W_s$ is no martingale but can be approximated by martingale on small time intervals
- $\triangleright \exp{\{\gamma\langle\zeta_t,V(t)^{-1}\zeta_t\rangle\}}$ approximated by submartingale
- \triangleright Doob's inequality yields bound on probability of leaving $\mathcal{B}(h)$ during small time intervals. Then sum over all time intervals
- \triangleright Nonlinear equation : $d\zeta_t = A(t)\zeta_t dt + b(\zeta_t, t) dt + \sigma dW_t$

$$\zeta_t = \sigma \int_0^t U(t,s) dW_s + \int_0^t U(t,s)b(\zeta_s,s) ds$$

Second integral can be treated as small perturbation for $t \leqslant \tau_{\mathcal{B}(h)}$

One shows that for z = 0

- \triangleright The distance between the $k^{\rm th}$ and $k+1^{\rm st}$ canard has order ${\rm e}^{-(2k+1)^2\mu}$
- \triangleright The section of $\mathcal{B}(h)$ is close to circular with radius $\mu^{-1/4}h$

One shows that for z=0

- \triangleright The distance between the k^{th} and $k+1^{\text{st}}$ canard has order $\mathrm{e}^{-(2k+1)^2\mu}$
- \triangleright The section of $\mathcal{B}(h)$ is close to circular with radius $\mu^{-1/4}h$

Sketch of proof:

- Dynamic diagonalization of equation linearized around central ("weak") canard
- $\triangleright V(t) = \sigma^{-2} \operatorname{Cov}(\zeta_t)$ satisfies fast-slow equation

$$\mu \frac{\mathrm{d}V}{\mathrm{d}z} = A(z)V + VA(z)^T + 1$$

which can be studied by singular perturbation theory.

Note: Hopf bifurcation at z = 0!

One shows that for z=0

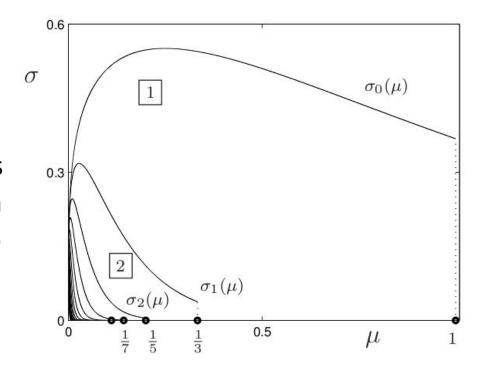
- \triangleright The distance between the k^{th} and $k+1^{\text{st}}$ canard has order $\mathrm{e}^{-(2k+1)^2\mu}$
- \triangleright The section of $\mathcal{B}(h)$ is close to circular with radius $\mu^{-1/4}h$

Corollary

Let

$$\sigma_k(\mu) = \mu^{1/4} e^{-(2k+1)^2 \mu}$$

Canards with $\frac{2k+1}{4}$ oscillations become indistinguishable from noisy fluctuations for $\sigma > \sigma_k(\mu)$



One shows that for z=0

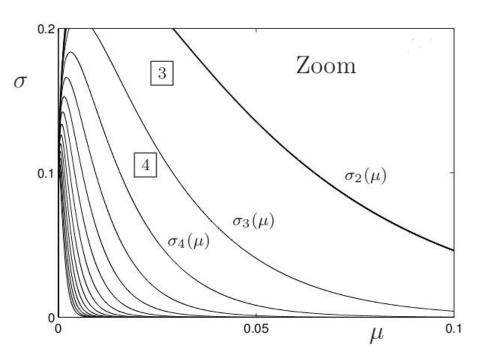
- \triangleright The distance between the k^{th} and $k+1^{\text{st}}$ canard has order $\mathrm{e}^{-(2k+1)^2\mu}$
- \triangleright The section of $\mathcal{B}(h)$ is close to circular with radius $\mu^{-1/4}h$

Corollary

Let

$$\sigma_k(\mu) = \mu^{1/4} e^{-(2k+1)^2 \mu}$$

Canards with $\frac{2k+1}{4}$ oscillations become indistinguishable from noisy fluctuations for $\sigma > \sigma_k(\mu)$



Early transitions

Let \mathcal{D} be neighbourhood of size \sqrt{z} of a canard for z > 0 (unstable)

Theorem [B, Gentz, Kuehn 2010]

 $\exists \kappa, C, \gamma_1, \gamma_2 > 0$ such that for $\sigma |\log \sigma|^{\gamma_1} \leqslant \mu^{3/4}$ probability of leaving \mathcal{D} after $z_t = z$ satisfies

$$\mathbb{P}\left\{z_{\tau_{\mathcal{D}}} > z\right\} \leqslant C |\log \sigma|^{\gamma_2} e^{-\kappa(z^2 - \mu)/(\mu |\log \sigma|)}$$

Small for $z\gg \sqrt{\mu|\log\sigma|/\kappa}$

Early transitions

Let \mathcal{D} be neighbourhood of size \sqrt{z} of a canard for z > 0 (unstable)

Theorem [B, Gentz, Kuehn 2010]

 $\exists \kappa, C, \gamma_1, \gamma_2 > 0$ such that for $\sigma |\log \sigma|^{\gamma_1} \leqslant \mu^{3/4}$ probability of leaving \mathcal{D} after $z_t = z$ satisfies

$$\mathbb{P}\left\{z_{\tau_{\mathcal{D}}} > z\right\} \leqslant C |\log \sigma|^{\gamma_2} e^{-\kappa(z^2 - \mu)/(\mu |\log \sigma|)}$$

Small for $z\gg\sqrt{\mu|\log\sigma|/\kappa}$

Sketch of proof:

- \triangleright Escape from neighbourhood of size $\sigma |\log \sigma|/\sqrt{z}$: compare with linearized equation on small time intervals + Markov property
- \triangleright Escape from annulus $\sigma |\log \sigma|/\sqrt{z} \leqslant ||\zeta|| \leqslant \sqrt{z}$: use polar coordinates and averaging
- ▶ To combine the two regimes : use Laplace transforms

Early transitions

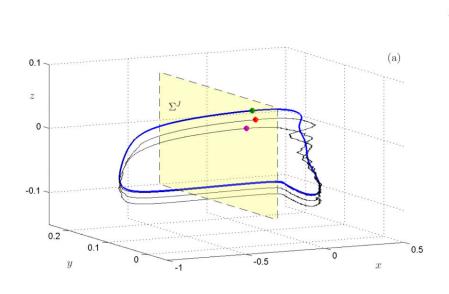
Let \mathcal{D} be neighbourhood of size \sqrt{z} of a canard for z>0 (unstable)

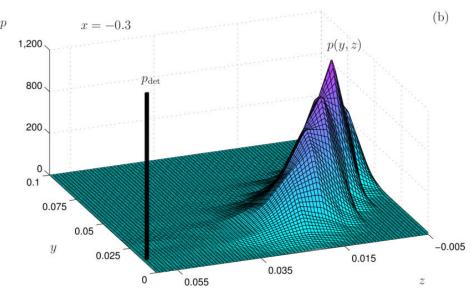
Theorem [B, Gentz, Kuehn 2010]

 $\exists \kappa, C, \gamma_1, \gamma_2 > 0$ such that for $\sigma |\log \sigma|^{\gamma_1} \leqslant \mu^{3/4}$ probability of leaving \mathcal{D} after $z_t = z$ satisfies

$$\mathbb{P}\left\{z_{\tau_{\mathcal{D}}} > z\right\} \leqslant C |\log \sigma|^{\gamma_2} e^{-\kappa(z^2 - \mu)/(\mu |\log \sigma|)}$$

Small for $z\gg \sqrt{\mu|\log\sigma|/\kappa}$



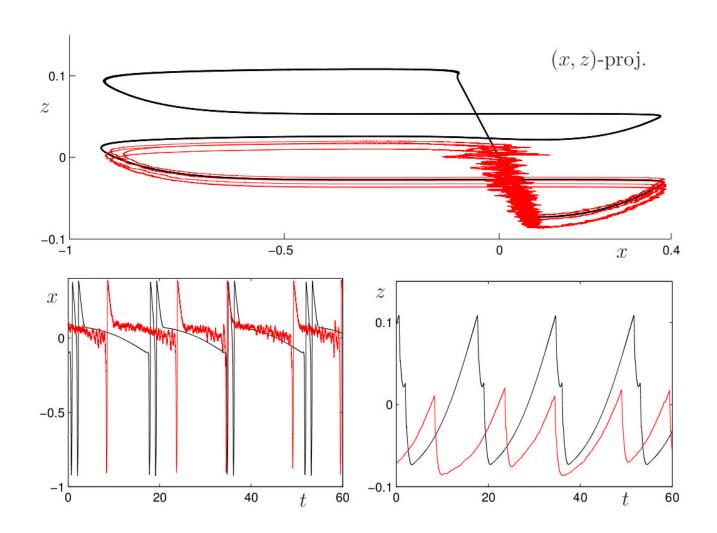


Further work

- ▶ Better understanding of distribution of noise-induced transitions
- ▷ Effect on mixed-mode pattern in conjunction with global return mechanism

Further work

- ▶ Better understanding of distribution of noise-induced transitions
- ▶ Effect on mixed-mode pattern in conjunction with global return mechanism



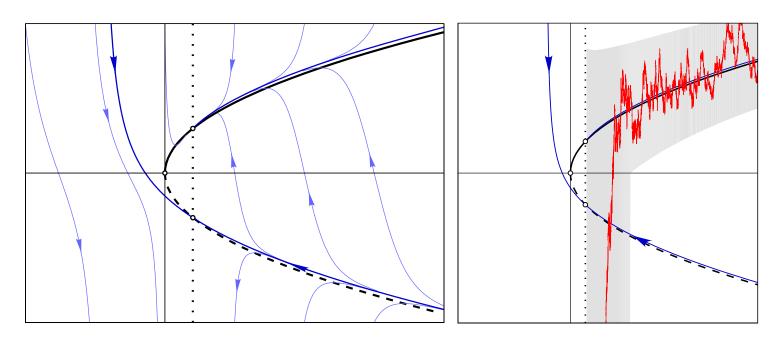
Noise-induced MMOs [D. Landon, PhD thesis, in progress]

FitzHugh-Nagumo, normal form near bifurcation point:

$$dx_t = (y_t - x_t^2) dt + \sigma dW_t$$
$$dy_t = \varepsilon(\delta - x_t) dt$$

 $\triangleright \delta > \sqrt{\varepsilon}$: equilibrium (δ, δ^2) is a node, effectively 1D problem

- $\bullet \ \sigma \ll \delta^{3/2}$: rare spikes, approx. exponential interspike times
- $\sigma \gg \delta^{3/2}$: repeated spikes



Noise-induced MMOs [D. Landon, PhD thesis, in progress]

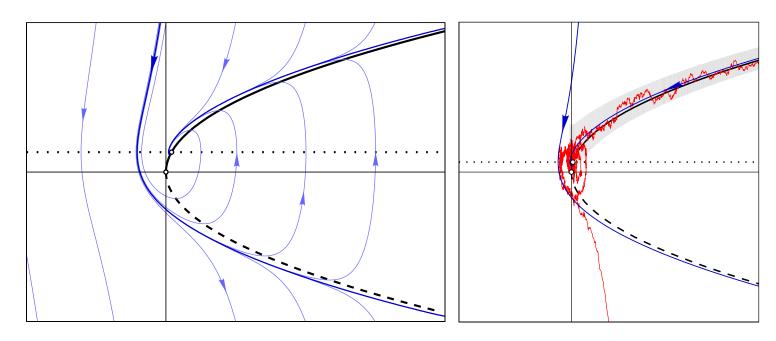
FitzHugh-Nagumo, normal form near bifurcation point:

$$dx_t = (y_t - x_t^2) dt + \sigma dW_t$$
$$dy_t = \varepsilon(\delta - x_t) dt$$

 $\triangleright \delta > \sqrt{\varepsilon}$: equilibrium (δ, δ^2) is a node, effectively 1D problem

- $\bullet \ \sigma \ll \delta^{3/2}$: rare spikes, approx. exponential interspike times
- $\sigma \gg \delta^{3/2}$: repeated spikes

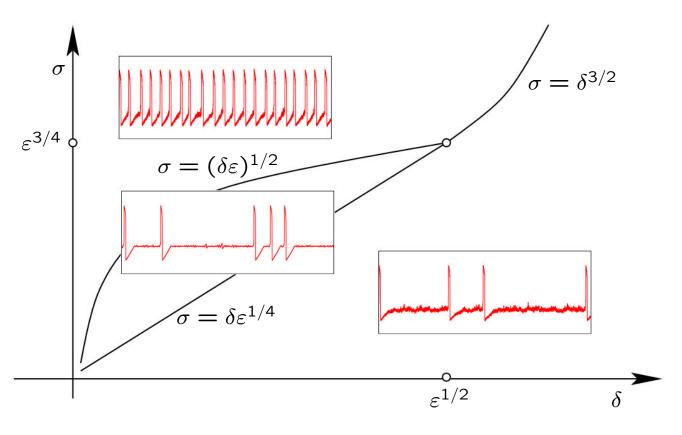
 $\triangleright \delta < \sqrt{\varepsilon}$: equilibrium (δ, δ^2) is a focus. Two-dimensional problem



Noise-induced MMOs

[D. Landon, PhD thesis, in progress]

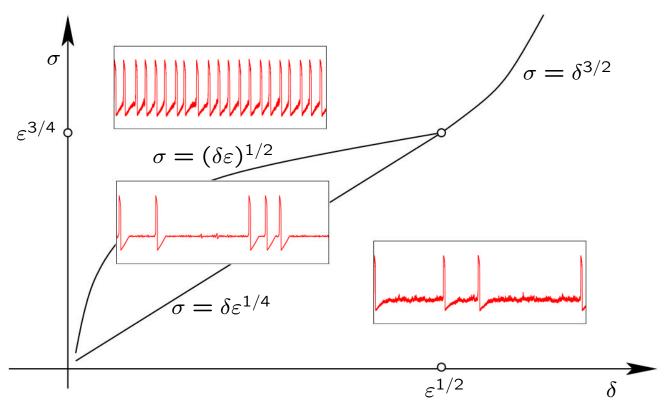
Conjectured bifurcation diagram [Muratov and Vanden Eijnden (2007)]:



Noise-induced MMOs

[D. Landon, PhD thesis, in progress]

Conjectured bifurcation diagram [Muratov and Vanden Eijnden (2007)]:



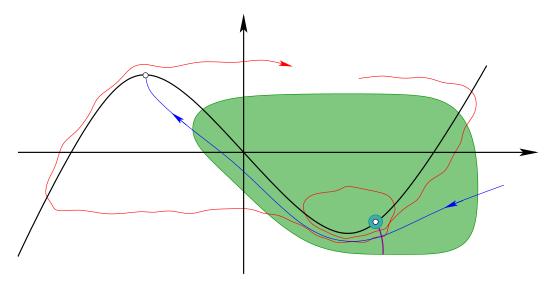
Work in progress:

- ▷ Prove bifurcation diagram is correct
- ▷ Characterize interspike time statistics and spike train statistics
- ▷ Characterize distribution of mixed-mode patterns

Noise-induced MMOs

[D. Landon, PhD thesis, in progress]

Definition of random number of SAOs N:



N = survival time of substochastic Markov chain

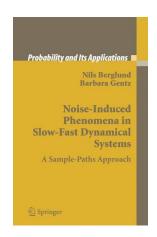
Theorem (2011):

- $\lim_{n\to\infty} \mathbb{P}\{N=n+1|N>n\}=1-\lambda_0,\ \lambda_0=\text{principal ev}$
- Weak noise: $\sigma_1^2 + \sigma_2^2 \leqslant (\varepsilon^{1/4}\delta)^2 \Rightarrow 1 \lambda_0 \leqslant e^{-\kappa(\varepsilon^{1/4}\delta)^2/(\sigma_1^2 + \sigma_2^2)}$
- Increasing noise:

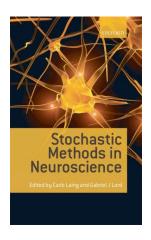
$$1 - \lambda_0 \simeq \Phi \left(- \frac{(\pi \varepsilon)^{1/4} (\delta - \sigma_1^2 / \varepsilon)}{\sqrt{\sigma_1^2 + \sigma_2^2}} \right)$$

References

N.B. and Barbara Gentz, *Noise-induced phe-nomena in slow-fast dynamical systems, A sample-paths approach*, Springer, Probability and its Applications (2006)



N.B. and Barbara Gentz, Stochastic dynamic bifurcations and excitability, in C. Laing and G. Lord, (Eds.), Stochastic methods in Neuroscience, p. 65-93, Oxford University Press (2009)



N.B., Barbara Gentz and Christian Kuehn, Hunting French Ducks in a Noisy Environment, hal-00535928, submitted (2010)

