Métastabilité dans des équations de Ginzburg–Landau avec bruit blanc spatiotemporel

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Metastability in physics

Examples:

- Supercooled liquid
- Supersaturated gas
- Wrongly magnetised ferromagnet

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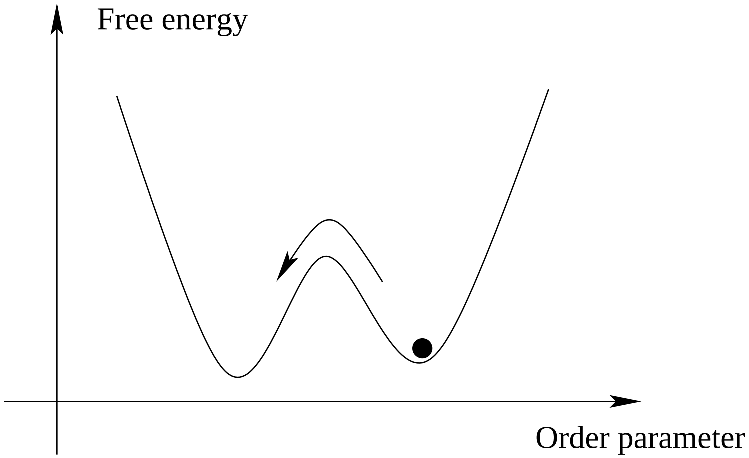
▷ Near first-order phase transition

▷ Nucleation implies crossing energy barrier



Metastability in stochastic lattice models

- ▷ Lattice: $\Lambda \subset\subset \mathbb{Z}^d$
- ▷ Configuration space: $\mathcal{X} = S^\Lambda$, S finite set (e.g. $\{-1, 1\}$)
- ▷ Hamiltonian: $H : \mathcal{X} \to \mathbb{R}$ (e.g. Ising or lattice gas)
- ▷ Gibbs measure: $\mu_\beta(x) = \mathrm{e}^{-\beta H(x)} / Z_\beta$
- ▷ Dynamics: Markov chain with invariant measure μ_β
  (e.g. Metropolis: Glauber or Kawasaki)

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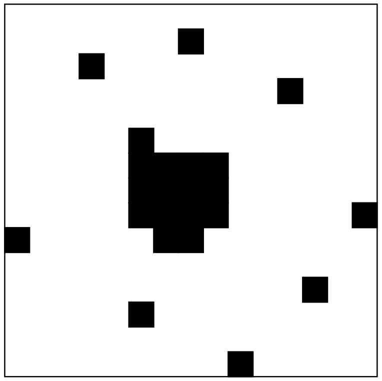
Results (for $\beta \gg 1$) on

- Transition time between $+$ and $-$ or empty and full configuration
- Transition path
- Shape of critical droplet



▷ Frank den Hollander, *Metastability under stochastic dynamics*, Stochastic Process. Appl. **114** (2004), 1–26.
▷ Enzo Olivieri and Maria Eulália Vares, *Large deviations and metastability*, Cambridge University Press, Cambridge, 2005.

Reversible diffusion

$$\mathrm{d}x_t = -\nabla V(x_t)\,\mathrm{d}t + \sqrt{2\varepsilon}\,\mathrm{d}W_t$$

- ▷ $V : \mathbb{R}^d \to \mathbb{R}$: potential, growing at infinity
- ▷ W_t: d-dim Brownian motion on $(\Omega, \mathcal{F}, \mathbb{P})$

Reversible w.r.t.
invariant measure:

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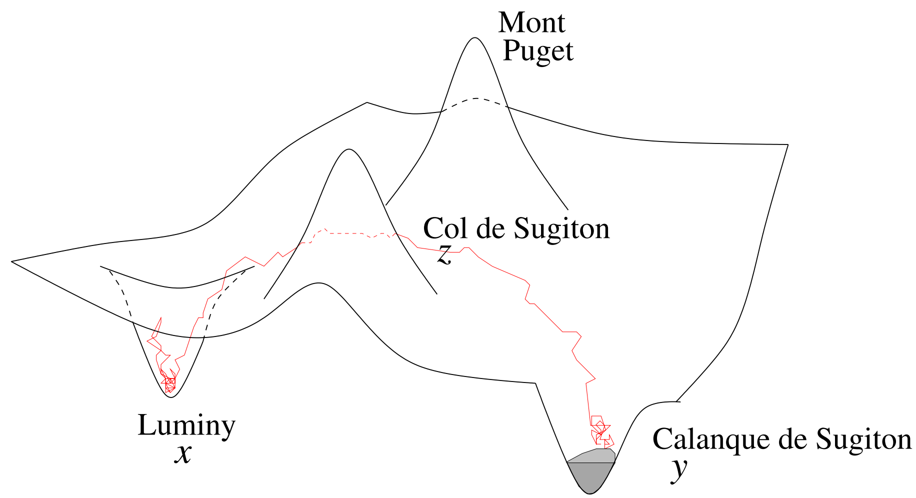

τ_y^x: first-hitting time of small ball $\mathcal{B}_\varepsilon(y)$, starting in x

"Eyring–Kramers law" (Eyring 1935, Kramers 1940)

- Dim 1: $\mathbb{E}[\tau_y^x] \simeq \dfrac{2\pi}{\sqrt{V''(x)|V''(z)|}}\,\mathrm{e}^{[V(z)-V(x)]/\varepsilon}$
- Dim$\geqslant 2$: $\mathbb{E}[\tau_y^x] \simeq \dfrac{2\pi}{|\lambda_1(z)|}\sqrt{\dfrac{|\det(\nabla^2 V(z))|}{\det(\nabla^2 V(x))}}\,\mathrm{e}^{[V(z)-V(x)]/\varepsilon}$

Towards a proof of Kramers' law

- Large deviations (Wentzell & Freidlin 1969):

$$\lim_{\varepsilon\to 0} \varepsilon \log(\mathbb{E}[\tau_y^x]) = V(z) - V(x)$$

- Analytic (Helffer, Sjöstrand 85, Miclo 95, Mathieu 95, Kolokoltsov 96,...): low-lying spectrum of generator
- Potential theory/variational (Bovier, Eckhoff, Gayrard, Klein 2004):

$$\mathbb{E}[\tau_y^x] = \frac{2\pi}{|\lambda_1(z)|}\sqrt{\frac{|\det(\nabla^2 V(z))|}{\det(\nabla^2 V(x))}}\, \mathrm{e}^{[V(z)-V(x)]/\varepsilon}\Big[1+\mathcal{O}(\varepsilon^{1/2}|\log\varepsilon|^{1/2})\Big]$$

  and similar asymptotics for eigenvalues of generator
- Witten complex (Helffer, Klein, Nier 2004): full asymptotic expansion of prefactor
- Distribution of τ_y^x (Day 1983, Bovier *et al* 2005):

$$\lim_{\varepsilon\to 0} \mathbb{P}\Big\{\tau_y^x > t\mathbb{E}[\tau_y^x]\Big\} = \mathrm{e}^{-t}$$

Ginzburg–Landau equation

$\partial_t u(x,t) = \partial_{xx} u(x,t) + u(x,t) - u(x,t)^3$ $+$ noise

$x \in [0, L]$, $u(x,t) \in \mathbb{R}$ represents e.g. magnetisation

- Periodic b.c.
- Neumann b.c. $\partial_x u(0,t) = \partial_x u(L,t) = 0$

Noise: weak, white in space and time

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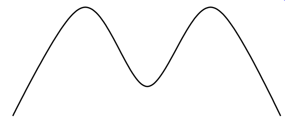
Deterministic system is gradient

$$\partial_{xx} u + u - u^3 = -\frac{\delta V}{\delta u}$$

$$V[u] = \int_0^L \left[\frac{1}{2} u'(x)^2 - \frac{1}{2} u(x)^2 + \frac{1}{4} u(x)^4\right] \mathrm{d}x \qquad \to \quad \text{min}$$

Stationary solutions

$$u''(x) = -u(x) + u(x)^3 = -\frac{\mathsf{d}}{\mathsf{d}x}\Bigg[$$



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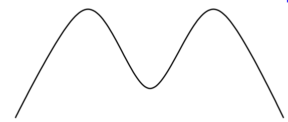

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- $u_{\pm}(x) \equiv \pm 1$: global minima of V, stable
- $u_0(x) \equiv 0$: unstable
- Neumann b.c: for $k = 1, 2, \ldots$, if $L > \pi k$,
  $u_{k,\pm}(x) = \pm\sqrt{\frac{2m}{m+1}}\,\mathrm{sn}\Big(\frac{kx}{\sqrt{m+1}} + \mathsf{K}(m), m\Big) \qquad 2k\sqrt{m+1}\,\mathsf{K}(m) = L$

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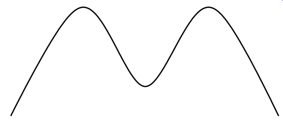
$$u''(x) = -u(x) + u(x)^3 = -\frac{\mathsf{d}}{\mathsf{d}x}\Big[$$  $$\Big]$$

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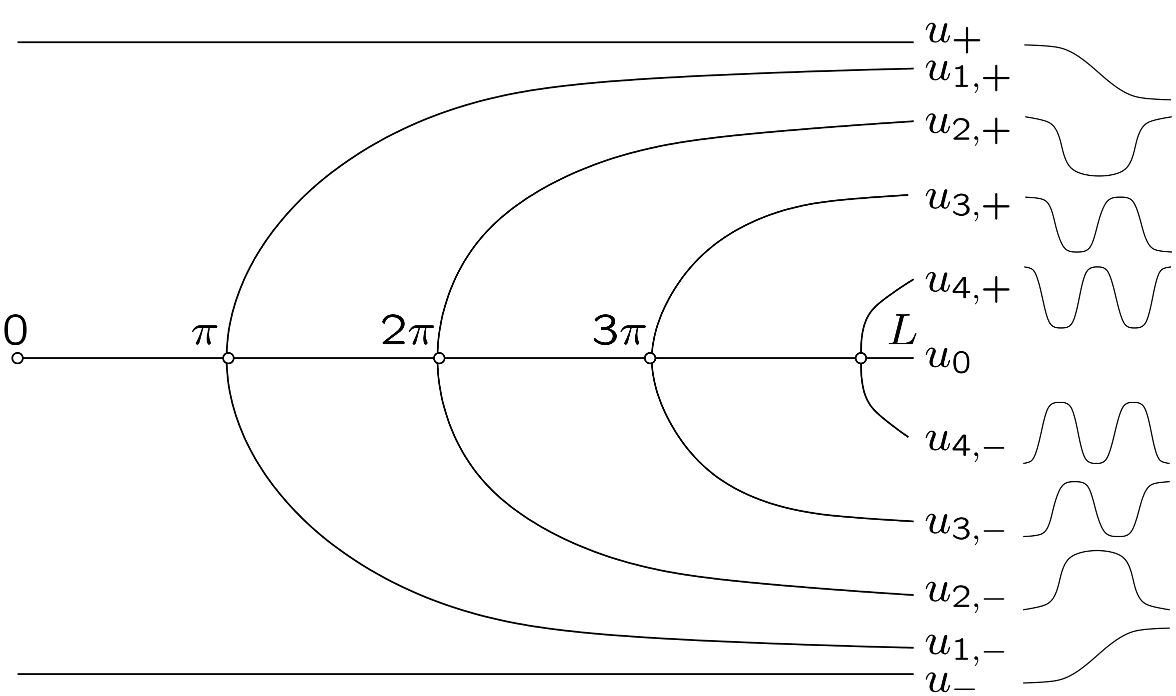

Stability of stationary solutions

Linearisation at $u(x)$: $\partial_t \varphi = A[u]\varphi$, $A[u] = \frac{\mathrm{d}^2}{\mathrm{d}x^2} + 1 - 3u(x)^2$

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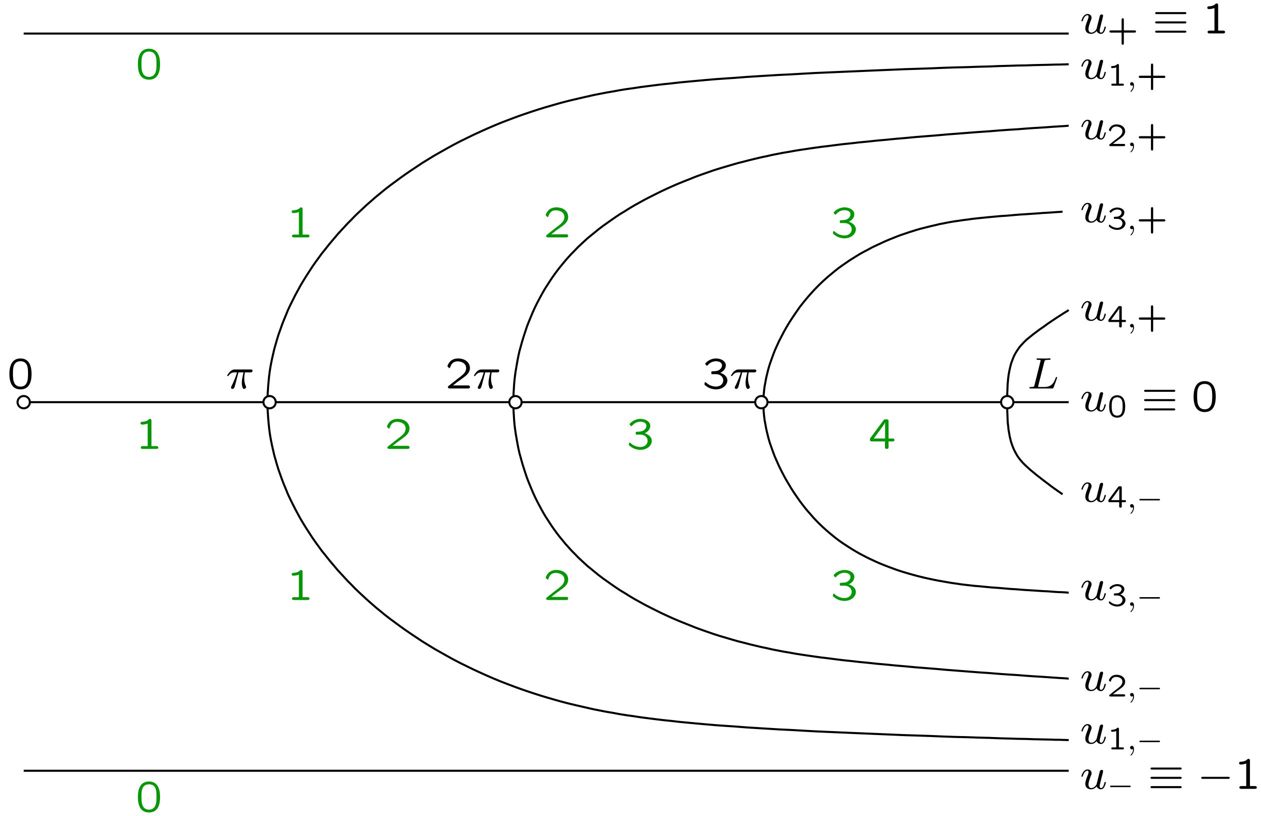
- $u_\pm(x) \equiv \pm 1$: eigenvalues $-(2 + (\pi k/L)^2)$, $k \in \mathbb{N}$
- $u_0(x) \equiv 0$: eigenvalues $1 - (\pi k/L)^2$, $k \in \mathbb{N}$

Number of positive eigenvalues:



Ginzburg–Landau equation with noise

$$\dot{u}_t(x) = \Delta u_t(x) + f(u_t(x)) \quad + \sqrt{2\varepsilon}\,\ddot{W}_{tx} \qquad\qquad (\Delta \equiv \partial_{xx}, f(u) = u - u^3)$$

$\ddot{W}_{tx}$: space–time white noise (formal derivative of Brownian sheet)

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Construction of mild solution:

1. $\dot{u}_t = \Delta u_t \qquad \Rightarrow \qquad u_t = \mathrm{e}^{\Delta t}\, u_0$

   where $\mathrm{e}^{\Delta t} \cos(k\pi x/L) = \mathrm{e}^{-(k\pi/L)^2 t} \cos(k\pi x/L)$

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2. $\dot{u}_t = \Delta u_t + \sqrt{2\varepsilon}\, \ddot{W}_{tx}$

   $$\Rightarrow \qquad u_t = \mathrm{e}^{\Delta t}\, u_0 + \sqrt{2\varepsilon}\, \underbrace{\int_0^t \mathrm{e}^{\Delta(t-s)}\, \dot{W}_x(\mathrm{d}s)}_{=: w_t(x)}$$

   $$w_t(x) = \sum_k \int_0^t \mathrm{e}^{-(k\pi/L)^2(t-s)}\, \mathrm{d}W_s^{(k)} \cos(k\pi x/L)$$

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   $\Rightarrow \qquad$ Existence and a.s. uniqueness (Faris & Jona-Lasinio 1982)

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Fourier variables: $u_t(x) = \sum_{k=-\infty}^{\infty} z_k(t)\,\mathrm{e}^{\mathrm{i}\,\pi k x/L}$

$$\Rightarrow \qquad \mathrm{d}z_k = -\lambda_k z_k\,\mathrm{d}t - \sum_{k_1+k_2+k_3=k} z_{k_1} z_{k_2} z_{k_3}\,\mathrm{d}t \; + \sqrt{2\varepsilon}\,\mathrm{d}W_t^{(k)}$$

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Energy functional:

$$\frac{1}{L} V[u] = \frac{1}{2} \sum_{k=-\infty}^{\infty} \lambda_k |z_k|^2 + \frac{1}{4} \sum_{k_1+k_2+k_3+k_4=0} z_{k_1} z_{k_2} z_{k_3} z_{k_4}$$

The question

How long does the system take to get from $u_-(x) \equiv -1$ to (a neighbourhood of) $u_+(x) \equiv 1$?

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Large deviations (Faris & Jona-Lasinio 1982)

- $\triangleright$ $L \leqslant \pi$: $\Delta W = V[u_0] - V[u_-] = \frac{L}{4}$
- $\triangleright$ $L > \pi$: $\Delta W = V[u_{1,\pm}] - V[u_-] = \frac{1}{3\sqrt{1+m}}\left[8\,\mathsf{E}(m) - \frac{(1-m)(3m+5)}{1+m}\,\mathsf{K}(m)\right]$

Formal computation for Ginzburg–Landau (R.S. Maier, D. Stein, 01)

Case $L < \pi$:

Kramers: $\Gamma_0 \simeq \dfrac{|\lambda_1(u_0)|}{2\pi}\sqrt{\dfrac{\det(\nabla^2 V[u_-])}{|\det(\nabla^2 V[u_0])|}}$

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Case $L > \pi$: Spectral determinant computed by Gelfand's method

Problems:

1. What happens when $L \to \pi_-$? (bifurcation)
2. Is the formal computation correct in infinite dimension?

Potential theory

Consider first Brownian motion $W_t^x = x + W_t$

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$$\begin{aligned} \Delta w_A(x) &= 1 \qquad && x \in A^c \\ w_A(x) &= 0 \qquad && x \in A \end{aligned}$$

$G_{A^c}(x,y)$ Green's function $\Rightarrow w_A(x) = \int_{A^c} G_{A^c}(x,y)\,\mathrm{d}y$

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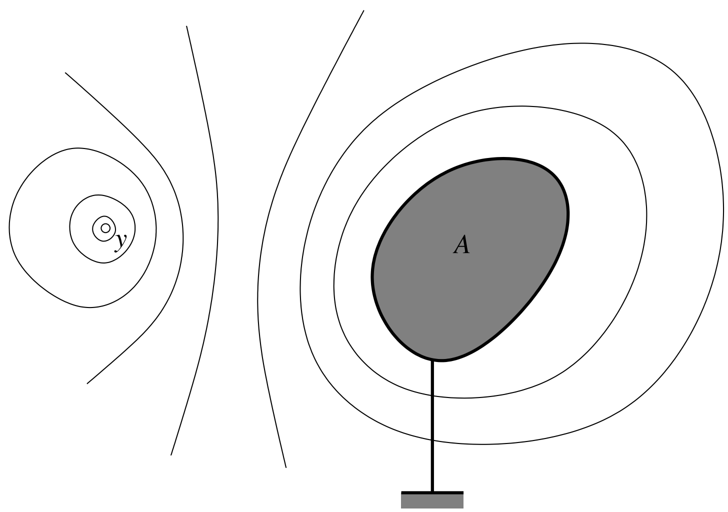

Potential theory

Fact 2: $h_{A,B}(x) = \mathbb{P}[\tau_A^x < \tau_B^x]$ satisfies

$$\begin{aligned} \Delta h_{A,B}(x) &= 0 \qquad && x \in (A \cup B)^c \\ h_{A,B}(x) &= 1 && x \in A \\ h_{A,B}(x) &= 0 && x \in B \end{aligned}$$

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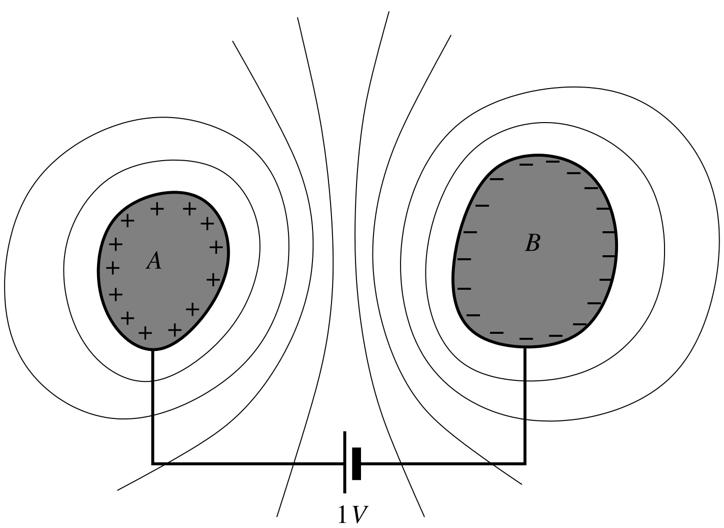
$\rho_{A,B}$: "surface charge density" on ∂A



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Capacity: $\mathsf{cap}_A(B) = \int_{\partial A} \rho_{A,B}(\mathsf{d}y)$

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Key observation: let $C = \mathcal{B}_\varepsilon(x)$, then (using $G(y,z) = G(z,y)$)

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$$\Rightarrow \qquad \mathbb{E}[\tau_A^x] = w_A(x) \simeq \frac{\int_{A^c} h_{\mathcal{B}_\varepsilon(x),A}(y)\,\text{d}y}{\text{cap}_{\mathcal{B}_\varepsilon(x)}(A)}$$

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Variational representation: Dirichlet form

$$\mathrm{cap}_A(B) = \int_{(A\cup B)^c} \|\nabla h_{A,B}(x)\|^2\,\mathrm{d}x = \inf_{h\in\mathcal{H}_{A,B}} \int_{(A\cup B)^c} \|\nabla h(x)\|^2\,\mathrm{d}x$$

($\mathcal{H}_{A,B}$: set of sufficiently smooth functions satisfying b.c.)

Potential theory

General case: $\mathrm{d}x_t = -\nabla V(x_t)\,\mathrm{d}t + \sqrt{2\varepsilon}\,\mathrm{d}W_t$

Generator: $\Delta \mapsto \varepsilon\Delta - \nabla V \cdot \nabla$

Then

$$\mathbb{E}[\tau_A^x] = w_A(x) \simeq \frac{\int_{A^c} h_{\mathcal{B}_\varepsilon(x),A}(y)\,\mathrm{e}^{-V(y)/\varepsilon}\,\mathrm{d}y}{\mathrm{cap}_{\mathcal{B}_\varepsilon(x)}(A)}$$

where

$$\mathrm{cap}_A(B) = \varepsilon \inf_{h\in\mathcal{H}_{A,B}} \int_{(A\cup B)^c} \|\nabla h(x)\|^2\,\mathrm{e}^{-V(x)/\varepsilon}\,\mathrm{d}x$$

Potential theory

General case: $\mathrm{d}x_t = -\nabla V(x_t)\,\mathrm{d}t + \sqrt{2\varepsilon}\,\mathrm{d}W_t$

Generator: $\Delta \mapsto \varepsilon\Delta - \nabla V \cdot \nabla$

Then

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Rough a priori bounds on h show that if x potential minimum,

$$\int_{A^c} h_{\mathcal{B}_\varepsilon(x),A}(y)\,\mathrm{e}^{-V(y)/\varepsilon}\,\mathrm{d}y \simeq \frac{(2\pi\varepsilon)^{d/2}\,\mathrm{e}^{-V(x)/\varepsilon}}{\sqrt{\det(\nabla^2 V(x))}}$$

Estimation of capacity

Truncated energy functional: retain only modes with $k \leqslant d$

$$\frac{1}{L}V[u] = -\frac{1}{2}z_0^2 + u_1(z_1) + \frac{1}{2}\sum_{k=2}^{d} \lambda_k |z_k|^2 + \ldots$$

$$u_1(z_1) = \tfrac{1}{2}\lambda_1 z_1^2 + \tfrac{3}{8} z_1^4$$

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Theorem: For all $L < \pi$,

$$\operatorname{cap}_{\mathcal{B}_\varepsilon(u_-)}(\mathcal{B}_\varepsilon(u_+)) = \varepsilon\frac{\displaystyle\int_{-\infty}^{\infty} \mathrm{e}^{-u_1(z_1)/\varepsilon}\,\mathrm{d}z_1}{\sqrt{2\pi\varepsilon}}\prod_{j=2}^{d}\sqrt{\frac{2\pi\varepsilon}{\lambda_j}}\Big[1+R(\varepsilon)\Big]$$

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Corollary:
Expected first-passage time converges in the limit $d \to \infty$
(Liu, 2003)

Sketch of proof

Upper bound:

$$\mathrm{cap} = \inf_h \Phi(h) \leqslant \Phi(h_+) \qquad \Phi(h) = \varepsilon \int \|\nabla h(z)\|^2 \, \mathrm{e}^{-V(z)/\varepsilon} \, \mathrm{d}z$$

Let $\delta = \sqrt{c\varepsilon|\log \varepsilon|}$, choose

$$h_+(z) = \begin{cases} 1 & \text{for } z_0 < -\delta \\ f(z_0) & \text{for } -\delta < z_0 < \delta \\ 0 & \text{for } z_0 > \delta \end{cases}$$

where $\varepsilon f''(z_0) + \partial_{z_0} V(z_0, 0) f'(z_0) = 0$ with b.c. $f(\pm\delta) = 0, 1$

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Lower bound:
Bound Dirichlet Φ form below by restricting domain, taking only 1st component of gradient and use for b.c. a priori bound on h

Computation of prefactor

$L < \pi$:

$$\Gamma_0 = \frac{1}{2^{3/4}\pi}\sqrt{\frac{\sinh(\sqrt{2}L)}{\sin L}}$$

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where $\lambda_1 = -1 + (\pi/L)^2$ and

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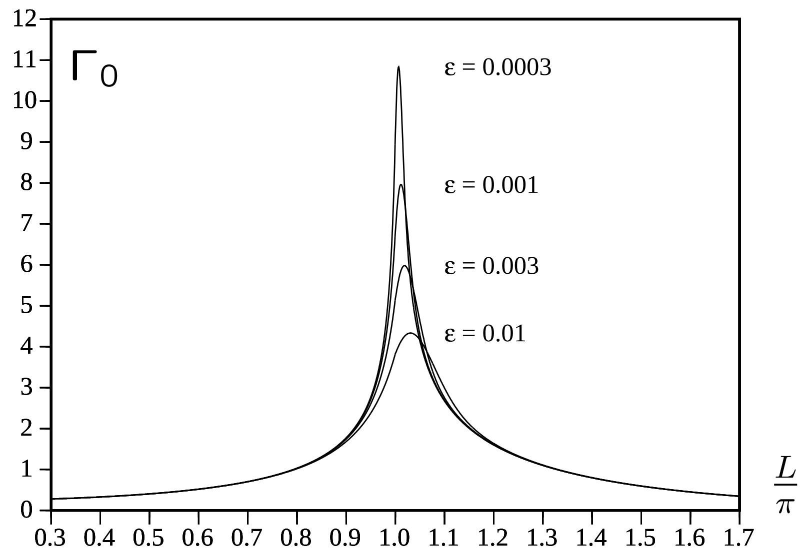
Similar expression for $L > \pi$

(product of eigenvalues computed using path-integral techniques, cf. Maier and Stein)



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