

Metastability
in irreversible diffusion processes
and stochastic resonance

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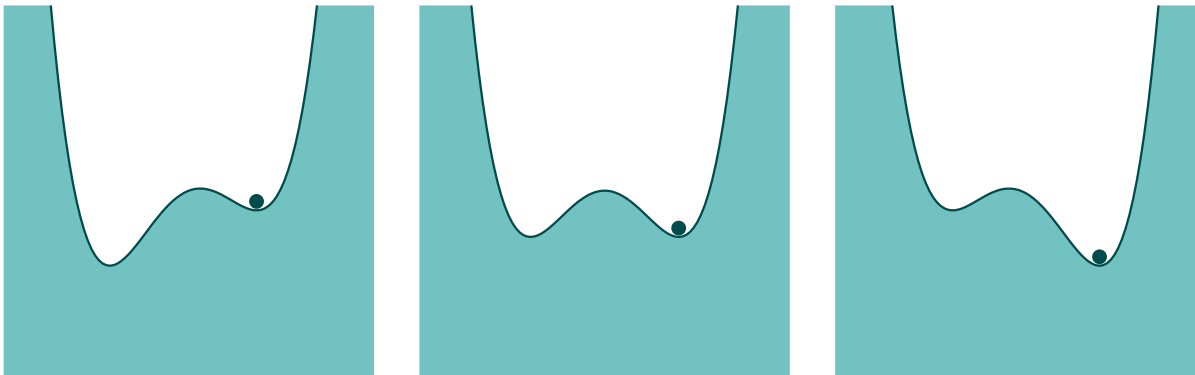
Joint work with [Barbara Gentz](#), WIAS, Berlin

Max-Planck-Institut Leipzig, May 2005

Stochastic resonance: typical example

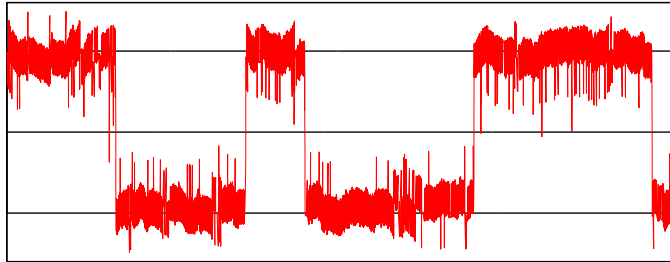
$$\begin{aligned} dx_t &= \underbrace{\left[x_t - x_t^3 + A \cos \varepsilon t \right]}_{= -\frac{\partial}{\partial x} V(x_t, t)} dt + \sigma dW_t \end{aligned}$$

Potential: $V(x, t) = \frac{1}{4}x^4 - \frac{1}{2}x^2 - Ax \cos \varepsilon t.$

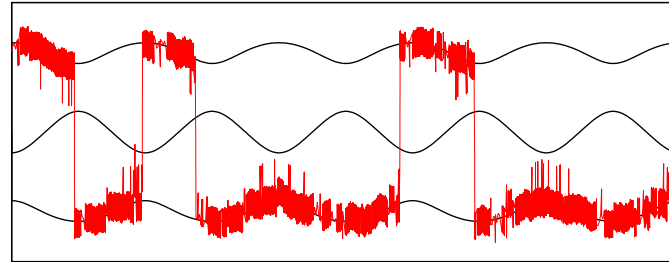


Sample paths $\{x_t\}_t$

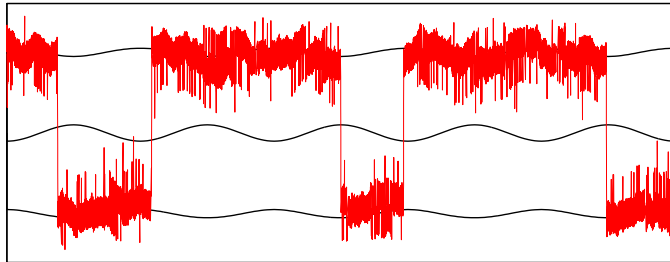
$\varepsilon = 0.001$



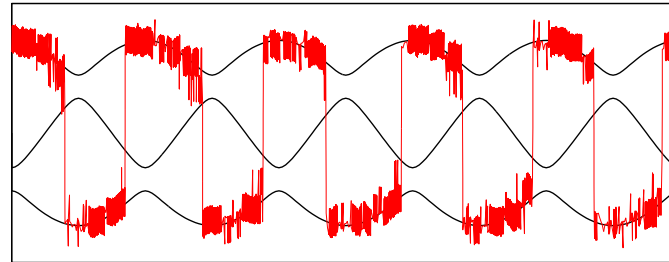
$A = 0, \sigma = 0.3$



$A = 0.24, \sigma = 0.2$



$A = 0.1, \sigma = 0.27$



$A = 0.35, \sigma = 0.2$

Qualitative measures:

- Power spectrum, signal-to-noise ratio.
- Residence-time distribution: law of interwell transition times.

Theory of large deviations

[Freidlin, Wentzell]

$$dx_t = f(x_t) dt + \sigma dW_t$$

$$\text{Rate function: } I_{[0,T]}(\varphi) = \begin{cases} \frac{1}{2} \int_0^T \|\dot{\varphi}_t - f(\varphi_t)\|^2 dt & \text{for } \varphi \in H_1 \\ +\infty & \text{otherwise} \end{cases}$$

Probability of staying close to $\varphi \sim e^{-I(\varphi)/\sigma^2}$.

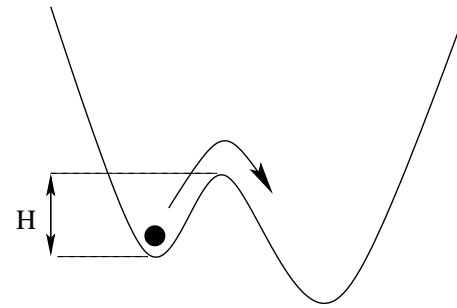
Large-deviation principle: for Γ set of paths,

$$\mathbb{P}\{(x_t)_{0 \leq t \leq T} \in \Gamma\} \sim e^{-\inf_{\Gamma} I/\sigma^2} \text{ as } \sigma \rightarrow 0.$$

Meaning:

$$\begin{aligned} -\inf_{\Gamma^\circ} I &\leq \liminf_{\sigma \rightarrow 0} \sigma^2 \log \mathbb{P}\{(x_t)_t \in \Gamma\} \\ &\leq \limsup_{\sigma \rightarrow 0} \sigma^2 \log \mathbb{P}\{(x_t)_t \in \Gamma\} \leq -\inf_{\Gamma} I \end{aligned}$$

Gradient case (reversible)



$$dx_t = -\nabla V(x_t) dt + \sigma dW_t$$

- $I(\varphi)$ minimized on paths with $\dot{\varphi}_t = -f(\varphi_t)$.
- Cost of leaving potential well
 $\inf\{I(\varphi) : \varphi_0 = \text{bottom}, \exists t : \varphi_t \notin \text{well}\} = 2H$.
- Expected time to leave well: $\mathbb{E}(\tau) \sim e^{2H/\sigma^2}$
[Eyring, Kramers, Freidlin, Wentzell, ...]
- Law of $\tau/\mathbb{E}(\tau) \rightarrow \text{Exp}(1)$ as $\sigma \rightarrow 0$ [Day]
- Subexponential behaviour known, related to small eigenvalues of generator of diffusion
[Bovier, Eckhoff, Gaynard, Klein], [Helffer, Klein, Nier]

Stochastic resonance: quasistatic regime

$$dx_t = \left[x_t - x_t^3 + A \cos \varepsilon t \right] dt + \sigma dW_t \quad (\text{for simplicity})$$

$$\text{Take } \varepsilon = \frac{2\pi}{T(\sigma)}, \quad T(\sigma) \sim e^{\lambda/\sigma^2}$$

x_t crosses barrier whenever $T(\sigma) \gg e^{2H(t)/\sigma^2}$

$H(t)$ = instantaneous barrier height

$\phi(t, \lambda)$: deterministic function, switches to deeper well whenever $2H(t) < \lambda$. (\Rightarrow possibility of **hysteresis**.)

Theorem: [Freidlin, 2000]

For $S, p, \delta > 0$,

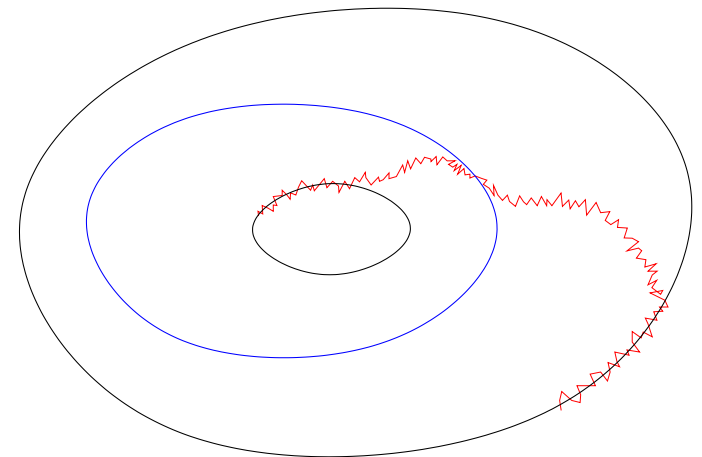
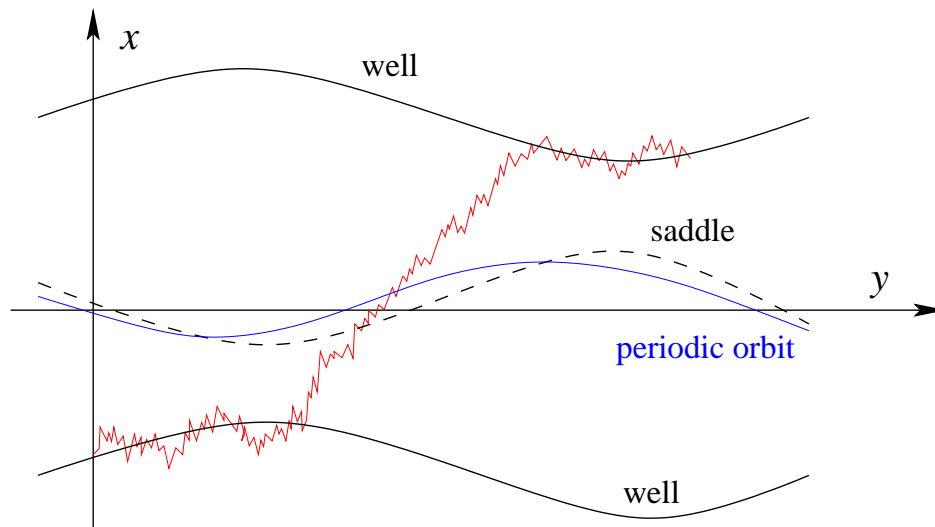
$$\lim_{\sigma \rightarrow 0} \mathbb{P} \left\{ \int_0^S |x_{sT(\sigma)} - \phi(sT(\sigma), \lambda)|^p ds > \delta \right\} = 0.$$

What about non-quasistatic regime?

$$\begin{aligned} dx_t &= [x_t - x_t^3 + A \cos \Omega y_t] dt + \sigma dW_t \\ dy_t &= dt \end{aligned}$$

Irreversible (degenerate) diffusion process.

Interwell transition \rightarrow crossing unstable periodic orbit tracking the saddle. **Distribution of transition locations?**



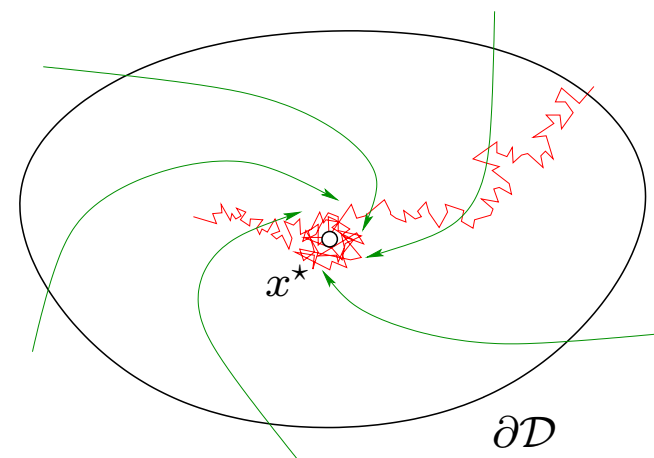
Exit problem

$$dx_t = f(x_t) dt + \sigma dW_t$$

$x_0 \in \mathcal{D}$, open, bdd set, $\partial\mathcal{D}$ smooth.

First-exit time: $\tau = \inf\{t > 0: x_t \notin \mathcal{D}\}$.

First-exit location: $x_\tau \in \partial\mathcal{D}$.



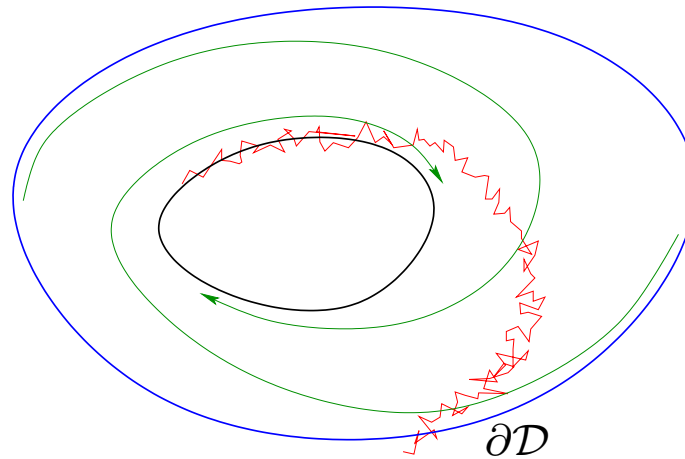
Case 1: $\bar{\mathcal{D}} \subset$ basin of attraction of asympt. stable equil. point x^*

Quasipotential

$$V(x^*, y) = \inf_{t > 0} \{I(\varphi) : \varphi_0 = x^*, \varphi_t = y\}.$$
$$\bar{V} = \inf_{y \in \partial\mathcal{D}} V(x^*, y).$$

- $\mathbb{E}(\tau) \sim e^{\bar{V}/\sigma^2}$
- $\lim_{\sigma \rightarrow 0} \mathbb{P}\left\{e^{(\bar{V}-\delta)/\sigma^2} \leq \tau \leq e^{(\bar{V}+\delta)/\sigma^2}\right\} = 1.$
- If $y \mapsto V(x^*, y)$ has non-degenerate global minimum in $x_1 \in \partial\mathcal{D}$,
 $\lim_{\sigma \rightarrow 0} \mathbb{P}\left\{\|x_\tau - x_1\| > \delta\right\} = 0.$ [Freidlin, Wentzell]

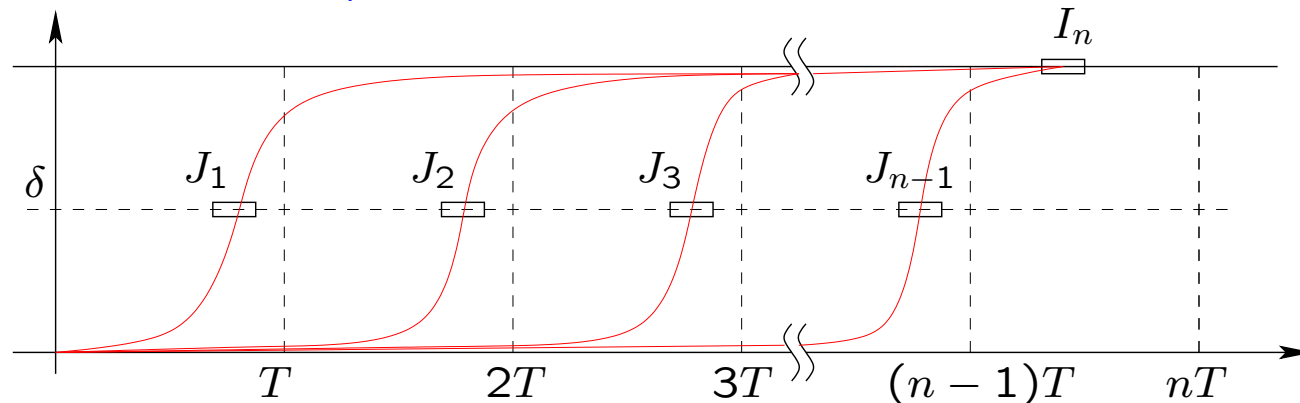
Exit problem



Case 2: $\partial\mathcal{D}$ is unstable periodic orbit (characteristic boundary)

- Quasipotential $V(x^*, y) = \bar{V}$ is constant on $\partial\mathcal{D}$.
 \Rightarrow **no concentration of exit location?**
- $\mathbb{E}(\tau) \sim e^{\bar{V}/\sigma^2}$ still holds. [Day]
- As $\sigma \rightarrow 0$, density of x_τ is translated along $\partial\mathcal{D}$ proportionally to $|\log \sigma|$: **cycling**. [Day]

How to compute law of exit location



If exit occurs in $I_n = [y, y + \Delta] \subset [nT, (n+1)T]$:
Rate function has n minimizers of comparable value.

$$\mathbb{P}^{0,0}\{y_\tau \in I_n\} \simeq \sum_{\ell=1}^n \underbrace{\mathbb{P}^{\delta, J_\ell}\{y_\tau \in I_n\}}_{Q_{n-\ell}(y)} \underbrace{\mathbb{P}^{0,0}\{y_{\tau'} \in J_\ell\}}_{P_\ell}$$

$$P_\ell \simeq \text{const} e^{-\ell\varepsilon} \exp\left\{-\frac{\bar{V}_1}{\sigma^2} \left(1 - e^{-2\ell\lambda T}\right)\right\}, \quad \varepsilon = T e^{-\bar{V}}/\sigma^2$$

$$Q_k(y) \simeq C(y) e^{-2k\lambda T} \exp\left\{-\frac{\bar{V}_2}{\sigma^2} \left(1 - c(y) e^{-2k\lambda T}\right)\right\}$$

$\lambda =$ Lyapunov exponent of unstable orbit.

Theorem:

$\partial\mathcal{D}$: unstable periodic orbit, Lyapunov exponent λ .

$\theta(y)$: “natural” parametrization of boundary, $\theta(y+T) = \theta(y) + \lambda T$.

$$\forall \Delta > \sigma^{1/2},$$

$$\mathbb{P}^{0, \theta^{-1}(\theta_0)} \{ \theta(y_{T\mathcal{D}}) \in [\theta_1, \theta_1 + \Delta] \} = \int_{\theta_1}^{\theta_1 + \Delta} p(\theta | \theta_0) d\theta [1 + \mathcal{O}(\sigma^{1/2})]$$

$$p(\theta | \theta_0) = \frac{1}{N} f_{\text{trans}}(\theta, \theta_0) \frac{e^{-(\theta - \theta_0)/\lambda T_{\mathcal{K}}}}{\lambda T_{\mathcal{K}}} P_{\lambda T}(\theta - |\log \sigma|)$$

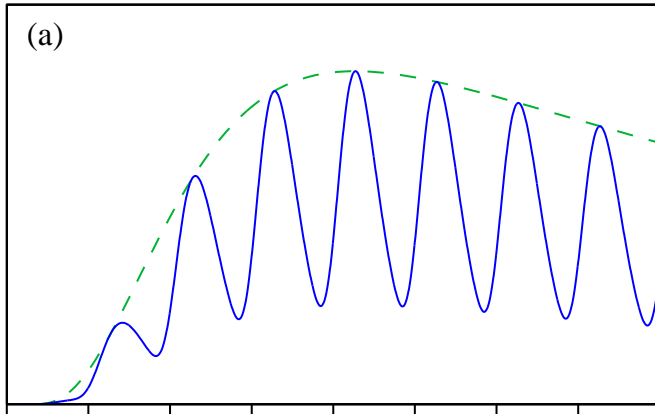
- $f_{\text{trans}}(\theta, \theta_0)$ grows from 0 to 1 for $\theta - \theta_0$ of order $|\log \sigma|$.

- $T_{\mathcal{K}} = T_{\mathcal{K}}(\sigma) = \frac{C}{\sigma} e^{\bar{V}/\sigma^2}$, Kramers time.

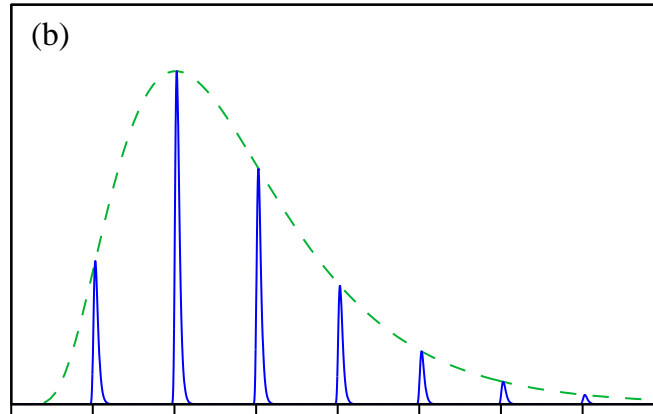
- $P_{\lambda T}(x) = \sum_{k=-\infty}^{\infty} A(x - k\lambda T)$ $A(z) = \frac{1}{2} e^{-2z} \exp\left\{-\frac{1}{2} e^{-2z}\right\}$.

First-passage-time distributions

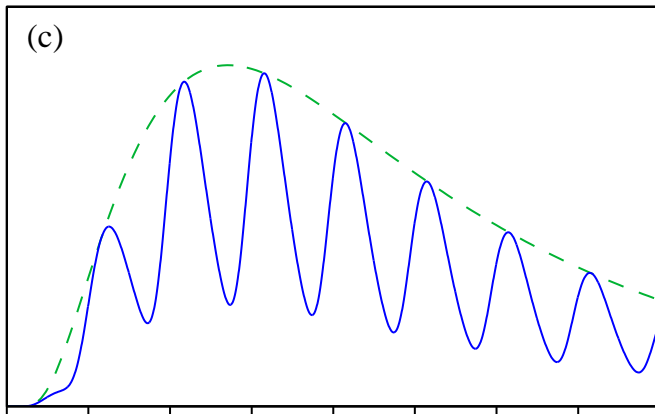
$$\bar{V} = 0.5, \lambda = 1$$



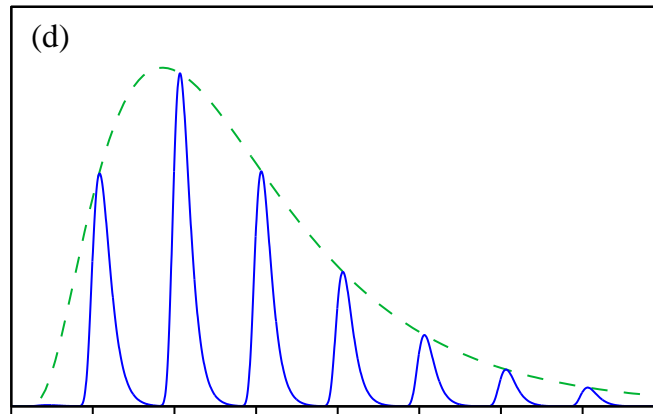
$\sigma = 0.4, T = 2$



$\sigma = 0.4, T = 20$



$\sigma = 0.5, T = 2$

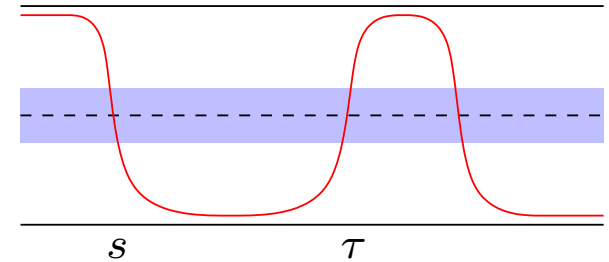


$\sigma = 0.5, T = 5$

Application to residence-time distribution

x_t crosses $x^{\text{per}}(t)$ at time s .

τ : time of first crossing back after s .



- First-passage-time density:

$$p(t|s) = \frac{\partial}{\partial t} \mathbb{P}^{s, x^{\text{per}}(s)} \{ \tau < t \}.$$

- Asymptotic transition-phase density:

$$\psi(t) = \int_{-\infty}^t p(t|s) \psi(s - T/2) ds = \psi(t + T).$$

- Residence-time distribution:

$$q(t) = \int_0^T p(s + t|s) \psi(s - T/2) ds.$$

Computation of residence-time distribution

Without forcing ($A = 0$):

$p(t|s) \sim$ exponential, $\psi(t)$ uniform $\Rightarrow q(t) \sim$ exponential.

With forcing ($A \gg \sigma^2$): time change $t \mapsto \theta(t)/\lambda$

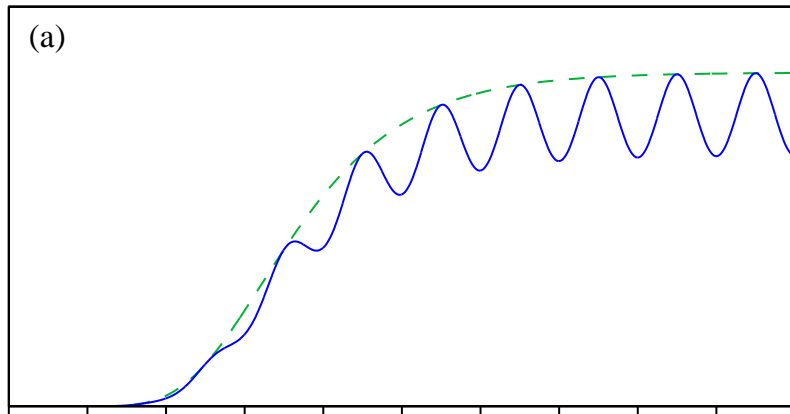
$$p(t|s) \simeq \frac{1}{N} f_{\text{trans}}(t, s) \frac{e^{-(t-s)/T_{\mathbf{K}}}}{T_{\mathbf{K}}} P_{\lambda T}(\lambda t - |\log \sigma|)$$

$$\psi(s) \simeq \frac{1}{T} P_{\lambda T}(\lambda t - |\log \sigma|) [1 + \mathcal{O}(T/T_{\mathbf{K}})]$$

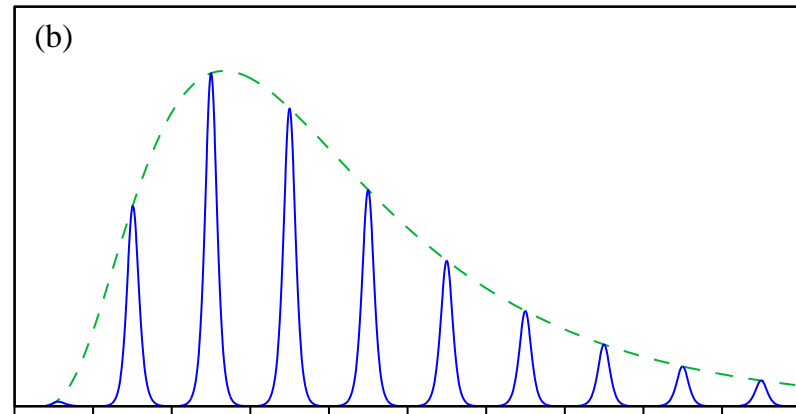
$$q(t) \simeq \tilde{f}_{\text{trans}}(t) \frac{e^{-t/T_{\mathbf{K}}}}{T_{\mathbf{K}}} \frac{\lambda T}{2} \sum_{k=-\infty}^{\infty} \frac{1}{\cosh^2(\lambda(t + T/2 - kT))}$$

Residence-time distributions

$$\bar{V} = 0.5, \lambda = 1$$



$$\sigma = 0.2, T = 2$$



$$\sigma = 0.4, T = 10$$

References

- N. B., B. Gentz, *On the noise-induced passage through an unstable periodic orbit I: Two-level model*, J. Statist. Phys. **114**, 1577–1618 (2004)
- _____, *Universality of first-passage and residence-time distributions in non-adiabatic stochastic resonance*, Europhys. Letters **70**, 1–7 (2005)
- In preparation . . .