

Quantification of unique continuation for the wave equation with Dirichlet boundary condition

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1 Introduction and main result

The purpose of this note is to prove the following result:

Theorem .- *Let Ω be a bounded connected C^2 domain in $\mathbb{R}^n$, $n > 1$, and ω be a non-empty open subset of Ω . Let us consider the following wave equation in $\Omega \times \mathbb{R}$, with initial data and homogeneous Dirichlet boundary condition :*

$$\begin{cases} \partial_t^2 u - \Delta u = 0 & \text{in } \Omega \times \mathbb{R} \\ u = 0 & \text{on } \partial\Omega \times \mathbb{R} \\ u(\cdot, 0) = u_0, \partial_t u(\cdot, 0) = u_1 & \text{in } \Omega . \end{cases} \quad (1.1)$$

Then, there exist a constant $c > 0$ and a time $T > 0$ such that for all initial data $(u_0, u_1) \in H_0^1(\Omega) \times L^2(\Omega)$, $(u_0, u_1) \neq 0$, the solution u of (1.1) satisfies the following estimate :

$$\|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2 \leq e^{\frac{c \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2}{\|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2}} \int_0^T \int_{\omega} |u(x, t)|^2 dx dt . \quad (1.2)$$

Remark that estimate (1.2) is equivalent to

$$\|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2 \leq \frac{c}{\ln \left(2 + \frac{\|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}}{\|u\|_{L^2(\omega \times (0, T))}} \right)} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 , \quad (1.3)$$

or equivalently,

$$\|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2 \leq ce^{c/\varepsilon} \int_0^T \int_{\omega} |u(x, t)|^2 dx dt + \varepsilon \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 , \quad \forall \varepsilon > 0 , \quad (1.4)$$

where the value of the constant $c > 0$ may be changed from line (1.2) to line (1.4) but not its dependence.

Similar kind of estimates already appears in [R], [LR] in the framework of boundary control and stabilization for hyperbolic equations. Here, we use the ideas from [R].

2 Proof of Theorem

2.1 step 1: Energy estimates

First, recall the conservation identity of the energy for the solution u of

$$\begin{cases} \partial_t^2 u - \Delta u = 0 & \text{in } \Omega \times \mathbb{R} \\ u = 0 & \text{on } \partial\Omega \times \mathbb{R} \\ u(\cdot, 0) = u_0, \partial_t u(\cdot, 0) = u_1 & \text{in } \Omega, \end{cases} \quad (2.1)$$

which is, for all $t \in \mathbb{R}$

$$\|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 = \int_{\Omega} \left(|\partial_t u(x, t)|^2 + |\nabla u(x, t)|^2 \right) dx. \quad (2.2)$$

Furthermore, we have the following inequality: there exists a constant $c > 0$ such that for all $T \geq 1$, we get

$$T \|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2 \leq c \int_0^T \int_{\Omega} |u(x, t)|^2 dx. \quad (2.3)$$

Indeed, let $\Psi \in C_0^\infty((0, T))$ such that $0 \leq \Psi \leq 1$, $\Psi \equiv 1$ on $(\frac{T}{4}, \frac{3T}{4})$ and $|\Psi'| \leq \frac{8}{T}$. Multiplying (2.1) by Ψu and integrating over $\Omega \times (0, T)$ by parts yields:

$$\begin{aligned} \|\nabla u\|_{L^2(\Omega \times (\frac{T}{4}, \frac{3T}{4}))}^2 &\leq \int_0^T \int_{\Omega} \Psi(t) |\nabla u(x, t)|^2 dx dt \\ &= - \int_0^T \int_{\Omega} \partial_t^2 u(x, t) \Psi(t) u(x, t) dx dt \\ &= \int_0^T \int_{\Omega} \partial_t u(x, t) \Psi'(t) u(x, t) dx dt + \int_0^T \int_{\Omega} \Psi(t) |\partial_t u(x, t)|^2 dx dt \\ &\leq \int_0^T \int_{\Omega} |\partial_t u(x, t)|^2 dx dt + \|\Psi' \partial_t u\|_{L^2(\Omega \times (0, T))} \|u\|_{L^2(\Omega \times (0, T))} \\ &\leq \int_0^T \int_{\Omega} |\partial_t u(x, t)|^2 dx dt + \frac{8}{T} \|\partial_t u\|_{L^2(\Omega \times (0, T))} \|u\|_{L^2(\Omega \times (0, T))} \end{aligned} \quad (2.4)$$

Adding $\|\partial_t u\|_{L^2(\Omega \times (\frac{T}{4}, \frac{3T}{4}))}^2$, using Poincaré inequality and conservation of energy (2.2), then there exists a constant $c > 0$ such that for all $T \geq 1$,

$$T \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \leq \sqrt{cT} \|\partial_t u\|_{L^2(\Omega \times (0, T))} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}, \quad (2.5)$$

that is

$$T \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \leq c \int_0^T \int_{\Omega} |\partial_t u(x, t)|^2 dx. \quad (2.6)$$

Now, we deduce (2.3) from (2.6). Indeed, let $v(x, t) = \int_0^t u(x, \ell) d\ell - (-\Delta^{-1}) u_1(x)$, then $\Delta v(x, t) = \int_0^t \Delta u(x, \ell) d\ell + u_1(x)$ and $\partial_t v(x, t) = u(x, t)$,

$$\begin{aligned} \partial_t^2 v(x, t) &= \partial_t u(x, t) \\ &= \partial_t u(x, 0) + \int_0^t \partial_t^2 u(x, \ell) d\ell \\ &= u_1(x) + \int_0^t \Delta u(x, \ell) d\ell \\ &= \Delta v(x, t). \end{aligned} \quad (2.7)$$

Also, $v(x, 0) = -(-\Delta^{-1}) u_1(x)$, $\partial_t v(x, 0) = u_0(x)$, and the homogeneous Dirichlet boundary condition is satisfied. Consequently,

$$\begin{aligned} \|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2 &= \|((- \Delta^{-1}) u_1, u_0)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \\ &= \|(v(x, 0), \partial_t v(x, 0))\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \leq \frac{c}{T} \int_0^T \int_{\Omega} |\partial_t v(x, t)|^2 dx dt \\ &\leq \frac{c}{T} \int_0^T \int_{\Omega} |u(x, t)|^2 dx dt, \end{aligned} \quad (2.8)$$

that is (2.3).

2.2 step 2: FBI transform

Let $F(z) = \frac{1}{2\pi} \int_{\mathbb{R}} e^{iz\tau} e^{-\tau^2} d\tau$, then $F(z) = \frac{\sqrt{\pi}}{2\pi} e^{\frac{1}{4}(|\operatorname{Im} z|^2 - |\operatorname{Re} z|^2)} e^{-\frac{i}{2}(\operatorname{Im} z \operatorname{Re} z)}$. Also, with $\lambda > 0$, let us consider

$$F_\lambda(z) = \lambda F(\lambda z) = \frac{1}{2\pi} \int_{\mathbb{R}} e^{iz\tau} e^{-\left(\frac{\tau}{\lambda}\right)^2} d\tau. \quad (2.9)$$

We have,

$$|F_\lambda(z)| = \frac{\sqrt{\pi}}{2\pi} \lambda e^{\frac{\lambda^2}{4}(|\operatorname{Im} z|^2 - |\operatorname{Re} z|^2)}, \quad (2.10)$$

and, denoting $\widehat{F}$ the Fourier transform of F ,

$$\widehat{F}_\lambda(\tau) = e^{-\left(\frac{\tau}{\lambda}\right)^2}. \quad (2.11)$$

Let $s, \ell_o \in \mathbb{R}$, we introduce the following gaussian transformation which is a particular case of the Fourier-Bros-Iagolnitzer transformation in [LR] :

$$W_{\ell_o, \lambda}(x, s) = \int_{\mathbb{R}} F_\lambda(\ell_o + is - \ell) \Phi(\ell) u(x, \ell) d\ell \quad (2.12)$$

where $\Phi \in C_0^\infty(\mathbb{R})$. We remark that $\partial_s^2 F_\lambda(\ell_o + is - \ell) = -\partial_\ell^2 F_\lambda(\ell_o + is - \ell)$, so

$$\begin{aligned} \partial_s^2 W_{\ell_o, \lambda}(x, s) &= \int_{\mathbb{R}} -\partial_\ell^2 F_\lambda(\ell_o + is - \ell) \Phi(\ell) u(x, \ell) d\ell \\ &= \int_{\mathbb{R}} -F_\lambda(\ell_o + is - \ell) [\Phi''(\ell) u(x, \ell) + 2\Phi'(\ell) \partial_t u(x, \ell) + \Phi(\ell) \partial_t^2 u(x, \ell)] d\ell. \end{aligned} \quad (2.13)$$

As u is solution of the wave equation, $W_{\ell_o, \lambda}$ satisfies:

$$\begin{cases} \partial_s^2 W_{\ell_o, \lambda}(x, s) + \Delta W_{\ell_o, \lambda}(x, s) = \int_{\mathbb{R}} -F_\lambda(\ell_o + is - \ell) [\Phi''(\ell) u(x, \ell) + 2\Phi'(\ell) \partial_t u(x, \ell)] d\ell \\ W_{\ell_o, \lambda}(x, s) = 0 \quad \text{for } x \in \partial\Omega \\ W_{\ell_o, \lambda}(x, 0) = (F_\lambda * \Phi u(x, \cdot))(\ell_o) \quad \text{for } x \in \Omega. \end{cases} \quad (2.14)$$

In another hand, we also have for some $T > 0$,

$$\begin{aligned} \|\Phi u(x, \cdot)\|_{L^2\left(\left(\frac{T}{2}-1, \frac{T}{2}+1\right)\right)} &\leq \|\Phi u(x, \cdot) - F_\lambda * \Phi u(x, \cdot)\|_{L^2\left(\left(\frac{T}{2}-1, \frac{T}{2}+1\right)\right)} \\ &\quad + \|F_\lambda * \Phi u(x, \cdot)\|_{L^2\left(\left(\frac{T}{2}-1, \frac{T}{2}+1\right)\right)} \\ &\leq \|\Phi u(x, \cdot) - F_\lambda * \Phi u(x, \cdot)\|_{L^2(\mathbb{R})} \\ &\quad + \left(\int_{t \in \left(\frac{T}{2}-1, \frac{T}{2}+1\right)} |W_{t, \lambda}(x, 0)|^2 dt \right)^{1/2}. \end{aligned} \quad (2.15)$$

But using Parseval equality and (2.11),

$$\begin{aligned} \|\Phi u(x, \cdot) - F_\lambda * \Phi u(x, \cdot)\|_{L^2(\mathbb{R})} &= \frac{1}{\sqrt{2\pi}} \left\| \widehat{\Phi u(x, \cdot)} - \widehat{F_\lambda * \Phi u(x, \cdot)} \right\|_{L^2(\mathbb{R})} \\ &= \frac{1}{\sqrt{2\pi}} \left(\int_{\mathbb{R}} \left| \left(1 - e^{-\left(\frac{\tau}{\lambda}\right)^2}\right) \widehat{\Phi u(x, \cdot)}(\tau) \right|^2 d\tau \right)^{1/2} \\ &\leq \frac{1}{\sqrt{2\pi}} \left(\int_{\mathbb{R}} \left| \left(\sqrt{2} \frac{\tau}{\lambda}\right) \widehat{\Phi u(x, \cdot)}(\tau) \right|^2 d\tau \right)^{1/2} \\ &\leq \frac{\sqrt{2}}{\sqrt{2\pi}} \frac{1}{\lambda} \left(\int_{\mathbb{R}} \left| \partial_t (\widehat{\Phi u(x, \cdot)}) (\tau) \right|^2 d\tau \right)^{1/2} \\ &\leq \sqrt{2} \frac{1}{\lambda} \|\partial_t (\Phi u(x, \cdot))\|_{L^2(\mathbb{R})}, \end{aligned} \quad (2.16)$$

so,

$$\int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |\Phi(t)u(x, t)|^2 dt \leq 4 \frac{1}{\lambda^2} \|\partial_t (\Phi u(x, \cdot))\|_{L^2(\mathbb{R})}^2 + 2 \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |W_{t, \lambda}(x, 0)|^2 dt \quad (2.17)$$

Now, recall that from the Cauchy theorem we have:

Proposition .- *Let f be a holomorphic function in a domain $D \subset \mathbb{C}$. Let $a, b > 0$, $z \in \mathbb{C}$. We suppose that*

$$D_o = \{(x, y) \in \mathbb{R}^2 \simeq \mathbb{C} \mid |x - \operatorname{Re}(z)| \leq a, |y - \operatorname{Im}(z)| \leq b\} \subset D,$$

then

$$f(z) = \frac{1}{\pi ab} \int \int_{\left| \frac{x - \operatorname{Re}(z)}{a} \right|^2 + \left| \frac{y - \operatorname{Im}(z)}{b} \right|^2 \leq 1} f(x + iy) dx dy.$$

Choosing $z = t \in (\frac{T}{2} - 1, \frac{T}{2} + 1) \subset \mathbb{R}$ and $x + iy = \ell_o + is$, we deduce that

$$\begin{aligned} |W_{t, \lambda}(x, 0)| &\leq \frac{1}{\pi ab} \int_{|\ell_o - t| \leq a} \int_{|s| \leq b} |W_{\ell_o + is, \lambda}(x, 0)| d\ell_o ds \\ &\leq \frac{1}{\pi ab} \int_{|\ell_o - t| \leq a} \int_{|s| \leq b} |W_{\ell_o, \lambda}(x, s)| ds d\ell_o \\ &\leq \frac{1}{\pi ab} \left(\int_{|\ell_o - t| \leq a} \int_{|s| \leq b} |W_{\ell_o, \lambda}(x, s)|^2 ds d\ell_o \right)^{1/2} \left(\int_{|\ell_o - t| \leq a} \int_{|s| \leq b} ds d\ell_o \right)^{1/2} \\ &\leq \frac{2}{\pi \sqrt{ab}} \left(\int_{|\ell_o - t| \leq a} \int_{|s| \leq b} |W_{\ell_o, \lambda}(x, s)|^2 ds d\ell_o \right)^{1/2}, \end{aligned} \quad (2.18)$$

and with $a = b = 1$,

$$\begin{aligned} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |W_{t, \lambda}(x, 0)|^2 dt &\leq \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} \left(\int_{|\ell_o - t| \leq 1} \int_{|s| \leq 1} |W_{\ell_o, \lambda}(x, s)|^2 ds d\ell_o \right) dt \\ &\leq \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} \int_{\ell_o \in (\frac{T}{2}-2, \frac{T}{2}+2)} \int_{|s| \leq 1} |W_{\ell_o, \lambda}(x, s)|^2 ds d\ell_o dt \\ &\leq 2 \int_{\ell_o \in (\frac{T}{2}-2, \frac{T}{2}+2)} \int_{|s| \leq 1} |W_{\ell_o, \lambda}(x, s)|^2 ds d\ell_o. \end{aligned} \quad (2.19)$$

Consequently, from (2.17), (2.19) and integrating over Ω , we get

$$\begin{aligned} \int_{\Omega} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |\Phi(t)u(x, t)|^2 dt dx &\leq 4 \frac{1}{\lambda^2} \int_{\Omega} \int_{\mathbb{R}} |\Phi'(t)u(x, t) + \Phi(t)\partial_t u(x, t)|^2 dt dx \\ &\quad + 4 \int_{\ell_o \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{\Omega} \int_{|s| \leq 1} |W_{\ell_o, \lambda}(x, s)|^2 ds dx \right) d\ell_o. \end{aligned} \quad (2.20)$$

Now recall the following quantification result for unique continuation of elliptic equation with Dirichlet boundary condition:

Proposition .- *Let Ω be a bounded connected C^2 domain in $\mathbb{R}^n$, $n > 1$, and ω be a non-empty open subset of Ω . Let $f = f(x, s) \in L^2(\Omega \times (-2, 2))$. Then there exists $C > 0$ such that for all $w = w(x, s) \in H^2(\Omega \times (-2, 2))$ solution of*

$$\begin{cases} \partial_s^2 w + \Delta_x w = f & \text{in } \Omega \times (-2, 2) \\ w = 0 & \text{on } \partial\Omega \times (-2, 2) \end{cases}$$

for all $\varepsilon > 0$, we have :

$$\begin{aligned} \int_{|s| \leq 1} \int_{\Omega} |w(x, s)|^2 dx ds &\leq C e^{C/\varepsilon} \left(\int_{|s| \leq 2} \int_{\omega} |w(x, s)|^2 dx ds + \int_{|s| \leq 2} \int_{\Omega} |f(x, s)|^2 dx ds \right) \\ &\quad + e^{-4/\varepsilon} \int_{|s| \leq 2} \int_{\Omega} |w(x, s)|^2 dx ds. \end{aligned}$$

Applying to $W_{\ell_0, \lambda}$, we deduce that for all $\varepsilon > 0$,

$$\begin{aligned}
& \int_{|s| \leq 1} \int_{\Omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \\
& \leq e^{-4/\varepsilon} \int_{|s| \leq 2} \int_{\Omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \\
& \quad + C e^{C/\varepsilon} \int_{|s| \leq 2} \int_{\omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \\
& \quad + C e^{C/\varepsilon} \int_{|s| \leq 2} \int_{\Omega} \left| \int_{\mathbb{R}} -F_{\lambda}(\ell_0 + is - \ell) [\Phi''(\ell) u(x, \ell) + 2\Phi'(\ell) \partial_t u(x, \ell)] d\ell \right|^2 dx ds.
\end{aligned} \tag{2.21}$$

Consequently, from (2.20) and (2.21), there exists a constant $C > 0$, such that for all $\varepsilon > 0$,

$$\begin{aligned}
& \int_{\Omega} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |\Phi(t) u(x, t)|^2 dt dx \\
& \leq 4 \frac{1}{\lambda^2} \int_{\Omega} \int_{\mathbb{R}} |\Phi'(t) u(x, t) + \Phi(t) \partial_t u(x, t)|^2 dt dx \\
& \quad + 4e^{-4/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \right) d\ell_0 \\
& \quad + 4C e^{C/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \right) d\ell_0 \\
& \quad + 4C e^{C/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_{\mathbb{R}} -F_{\lambda}(\ell_0 + is - \ell) [\Phi''(\ell) u(x, \ell) + 2\Phi'(\ell) \partial_t u(x, \ell)] d\ell \right|^2 dx ds \right) d\ell_0.
\end{aligned} \tag{2.22}$$

Let define $\Phi \in C_0^\infty(\mathbb{R})$ more precisely now: let $T > 16$, we choose $\Phi \in C_0^\infty((0, T))$, $0 \leq \Phi \leq 1$, $\Phi \equiv 1$ on $(\frac{T}{4}, \frac{3T}{4})$. Furthermore, let $K = [0, \frac{T}{4}] \cup [\frac{3T}{4}, T]$ such that $\text{supp}(\Phi') = K$ and $\text{supp}(\Phi'') \subset K$. Let $K_0 = [\frac{3T}{8}, \frac{5T}{8}]$ such that $\text{dist}(K, K_0) = \frac{T}{8}$. We will choose $\ell_0 \in (\frac{T}{2} - 2, \frac{T}{2} + 2) \subset K_0$.

Till the end of the proof, C_T will denote a generic positive constant independent of ε and λ but dependent on (Ω, T) , whose value may change from line to line.

The first term of the second member of (2.22) becomes, using (2.2),

$$\frac{1}{\lambda^2} \int_{\Omega} \int_{\mathbb{R}} |\Phi'(t) u(x, t) + \Phi(t) \partial_t u(x, t)|^2 dt dx \leq \frac{C_T}{\lambda^2} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \tag{2.23}$$

The second term of the second member of (2.22) becomes, using (2.10),

$$\begin{aligned}
& e^{-4/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \right) d\ell_0 \\
& = e^{-4/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_{\mathbb{R}} F_{\lambda}(\ell_0 + is - \ell) \Phi(\ell) u(x, \ell) d\ell \right|^2 dx ds \right) d\ell_0 \\
& \leq e^{-4/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_0^T \frac{\sqrt{\pi}}{2\pi} \lambda e^{\frac{\lambda^2}{4}(|s|^2 - |\ell_0 - \ell|^2)} |u(x, \ell)| d\ell \right|^2 dx ds \right) d\ell_0 \\
& \leq \left(\frac{\sqrt{\pi}}{2\pi} \lambda e^{\lambda^2} \right)^2 e^{-4/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_0^T |u(x, \ell)| d\ell \right|^2 dx ds \right) d\ell_0 \\
& \leq \left(\frac{\sqrt{\pi}}{2\pi} \lambda e^{\lambda^2} \right)^2 e^{-4/\varepsilon} 4 \left(4T \int_{\Omega} \int_0^T |u(x, \ell)|^2 d\ell dx \right) \\
& \leq C_T \lambda^2 e^{2\lambda^2} e^{-4/\varepsilon} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2
\end{aligned} \tag{2.24}$$

The third term of the second member of (2.22) becomes, using (2.10),

$$\begin{aligned}
& e^{C/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\omega} |W_{\ell_0, \lambda}(x, s)|^2 dx ds \right) d\ell_0 \\
& = e^{C/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\omega} \left| \int_{\mathbb{R}} F_{\lambda}(\ell_0 + is - \ell) \Phi(\ell) u(x, \ell) d\ell \right|^2 dx ds \right) d\ell_0 \\
& \leq e^{C/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\omega} \left| \int_0^T \frac{\sqrt{\pi}}{2\pi} \lambda e^{\frac{\lambda^2}{4}(|s|^2 - |\ell_0 - \ell|^2)} |\Phi(\ell) u(x, \ell)| d\ell \right|^2 dx ds \right) d\ell_0 \\
& \leq \left(\frac{\sqrt{\pi}}{2\pi} \lambda e^{\lambda^2} \right)^2 e^{C/\varepsilon} \int_{\ell_0 \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\omega} \left| \int_0^T |u(x, \ell)| d\ell \right|^2 dx ds \right) d\ell_0 \\
& \leq 4\lambda^2 T e^{2\lambda^2} e^{C/\varepsilon} \int_{\omega} \int_0^T |u(x, t)|^2 dt dx
\end{aligned} \tag{2.25}$$

The fourth term of the second member of (2.22) becomes, using (2.10) and the choice of Φ ,

$$\begin{aligned}
& e^{C/\varepsilon} \int_{\ell_o \in (\frac{T}{2}-2, \frac{T}{2}+2)} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_{\mathbb{R}} -F_{\lambda}(\ell_o + is - \ell) [\Phi''(\ell)u(x, \ell) + 2\Phi'(\ell)\partial_t u(x, \ell)] d\ell \right|^2 dx ds \right) d\ell_o \\
& \leq e^{C/\varepsilon} \int_{K_o} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_K \frac{\sqrt{\pi}}{2\pi} \lambda e^{\frac{\lambda^2}{4}(|s|^2 - |\ell_o - \ell|^2)} |\Phi''(\ell)u(x, \ell) + 2\Phi'(\ell)\partial_t u(x, \ell)| d\ell \right|^2 dx ds \right) d\ell_o \\
& \leq \left(\frac{\sqrt{\pi}}{2\pi} \lambda e^{\lambda^2} \right)^2 e^{C/\varepsilon} C_T \int_{K_o} \left(\int_{|s| \leq 2} \int_{\Omega} \left| \int_K e^{-\frac{\lambda^2}{4}|\ell_o - \ell|^2} (|u(x, \ell)| + |\partial_t u(x, \ell)|) d\ell \right|^2 dx ds \right) d\ell_o \\
& \leq C_T \left(\frac{\sqrt{\pi}}{2\pi} \lambda e^{\lambda^2} e^{-\frac{\lambda^2}{4} \text{dist}(K, K_o)^2} \right)^2 e^{C/\varepsilon} \int_{\Omega} \left| \int_K (|u(x, \ell)| + |\partial_t u(x, \ell)|) d\ell \right|^2 dx \\
& \leq C_T \lambda^2 e^{2\lambda^2} e^{-\frac{\lambda^2}{2} \text{dist}(K, K_o)^2} e^{C/\varepsilon} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \\
& \leq C_T \lambda^2 e^{\frac{\lambda^2}{2} \left(4 - \frac{T^2}{8^2}\right)} e^{C/\varepsilon} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2
\end{aligned} \tag{2.26}$$

We finally obtain from (2.23), (2.24), (2.25), (2.26) and (2.22), that

$$\begin{aligned}
\int_{\Omega} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |\Phi(t)u(x, t)|^2 dt dx & \leq \frac{C_T}{\lambda^2} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \\
& + C_T \lambda^2 e^{2\lambda^2} e^{-4/\varepsilon} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \\
& + 4c\lambda^2 T e^{2\lambda^2} e^{C/\varepsilon} \int_{\omega} \int_0^T |u(x, t)|^2 dt dx \\
& + C_T \lambda^2 e^{\frac{\lambda^2}{2} \left(4 - \frac{T^2}{8^2}\right)} e^{C/\varepsilon} \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 .
\end{aligned} \tag{2.27}$$

We begin to choose $\lambda^2 = \frac{1}{\varepsilon}$ in order that

$$\begin{aligned}
\int_{\Omega} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |\Phi(t)u(x, t)|^2 dt dx & \leq \varepsilon C_T \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \\
& + e^{-1/\varepsilon} C_T \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \\
& + e^{(3+C)/\varepsilon} C_T \int_{\omega} \int_0^T |u(x, t)|^2 dt dx \\
& + C_T \exp \left(\left(3 - \frac{1}{2} \frac{T^2}{8^2} + C \right) \frac{1}{\varepsilon} \right) \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 .
\end{aligned} \tag{2.28}$$

We finally choose $T > 16$ large enough such that $\left(3 - \frac{1}{2} \frac{T^2}{8^2} + C \right) \leq -1$ that is $8^2(8 + 2C) \leq T^2$, to deduce that

$$\begin{aligned}
\int_{\Omega} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |u(x, t)|^2 dt dx & \leq \int_{\Omega} \int_{t \in (\frac{T}{2}-1, \frac{T}{2}+1)} |\Phi(t)u(x, t)|^2 dt dx \\
& \leq \varepsilon C_T \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 + e^{(3+C)/\varepsilon} C_T \int_{\omega} \int_0^T |u(x, t)|^2 dt dx .
\end{aligned} \tag{2.29}$$

Now we conclude from (2.3), that there exist a constant $c > 0$ and a time $T > 0$ large enough such that for all $\varepsilon \in (0, 1)$ small enough, we have

$$\|(u_0, u_1)\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2 \leq e^{c/\varepsilon} \int_{\omega} \int_0^T |u(x, t)|^2 dt dx + \varepsilon \|(u_0, u_1)\|_{H_0^1(\Omega) \times L^2(\Omega)}^2 \tag{2.30}$$

References

- [LR] G. Lebeau and L. Robbiano, Stabilisation de l'équation des ondes par le bord, Duke Math. J. 86 (1997) 465-491.
- [R] L. Robbiano, Fonction de coût et contrôle des solutions des équations hyperboliques, Asymptotic Analysis 10 (1995) 95-115.